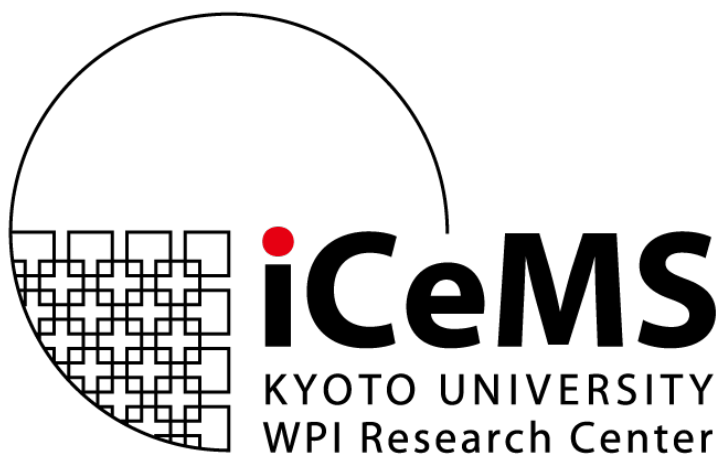


# Applications of data science and machine learning to organic materials

Daniel Packwood



# Self-introduction

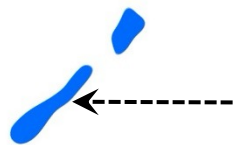
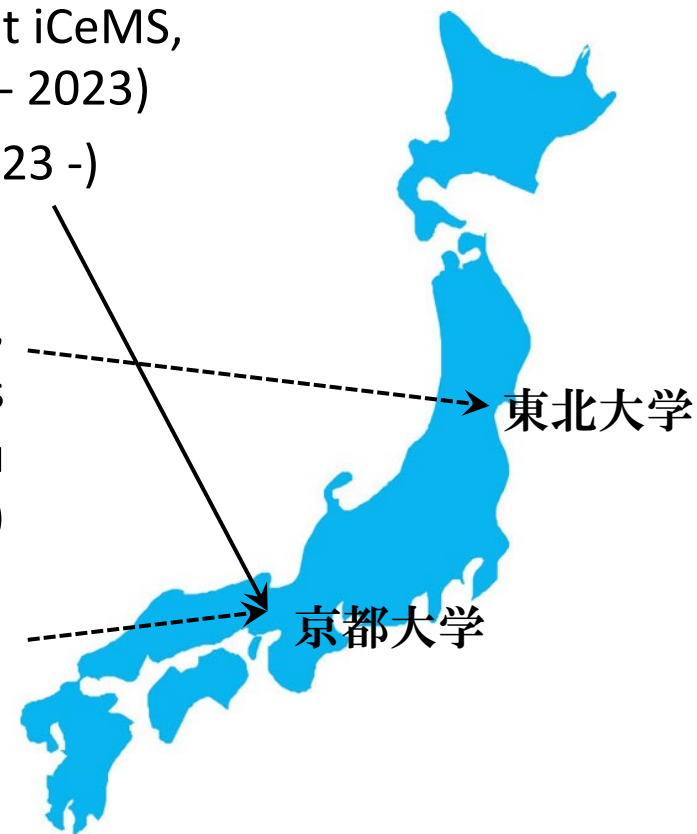


Senior lecturer and PI at iCeMS,  
Kyoto University (2016 - 2023)  
Associate Professor (2023 -)



Assistant Professor,  
Mathematical Sciences  
Unit, AIMR, Tohoku  
University (2012 – 2016)

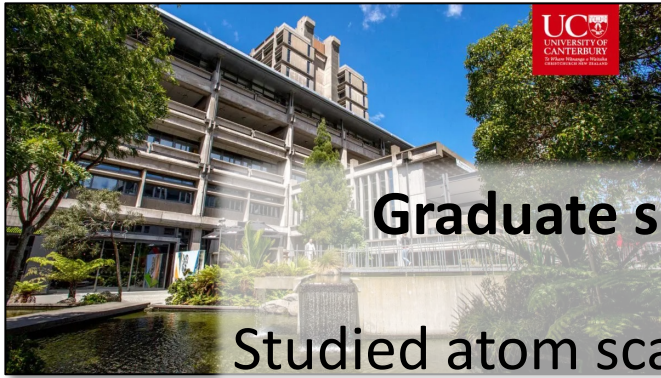
JSPS Postdoc, Quantum Chemistry Lab, Dept.  
of Chemistry, Graduate School of Science,  
Kyoto University (2010-2012)



University of Canterbury (PhD 2010)  
Major: Chemistry, Minor: Statistics



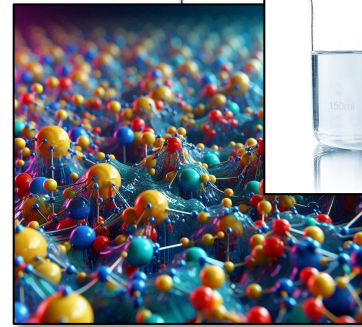
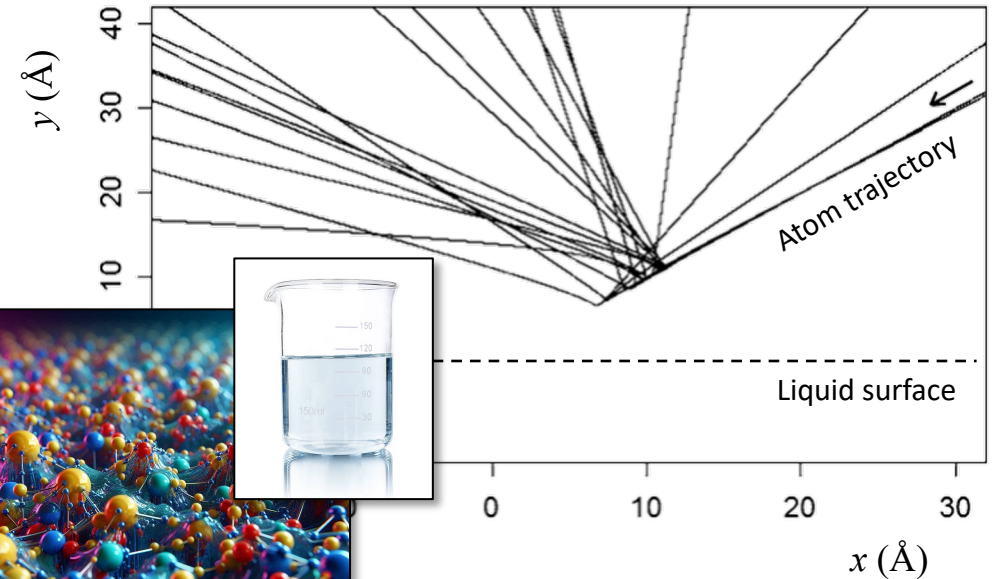




# Why did I come to Japan?

## Graduate school research (late 2000s)

Studied atom scattering from liquid surfaces using **non-equilibrium statistical mechanics** and **stochastic differential equations**.



Two names frequently appeared during my study:



**Kiyoshi Ito**

- Kyoto University mathematician
- **Stochastic differential equations pioneer**



**Ryogo Kubo**

- Tokyo University physicist
- **Non-equilibrium statistical mechanics pioneer**

**Realization: my field was pioneered in Japan!**

# Why did I come to Japan?



HOPE MEETINGS  
with Nobel Laureates

## ■ 第2回HOPEミーティング開催概要

日時：2009年9月27日（日）～10月1日（木）

会場：ザ・プリンス箱根

テーマ：Art in Science

対象分野：化学及び関連分野（物理学、生物学等）

主催：（独）日本学術振興会

<https://www.jsps.go.jp/hope/gaiyou2.html>



- Gathering of around 50 students from the Asia-Pacific region. Activities with Japanese graduate students and lectures from Japanese Nobel laureates.
- High level of research from the students impressed me, left me with a strong impression.

**→ Go to Japan for postdoctoral research!**



# Career in Japan

Tanimura laboratory, Department of Chemistry, Grad. School of Science, Kyoto University

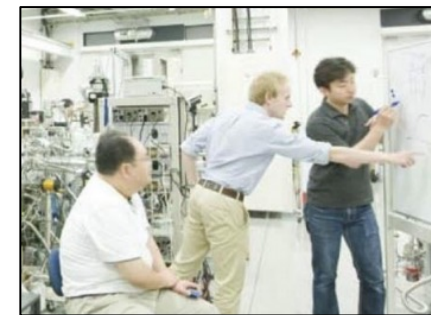


2010

Postdoc

2012

Mathematical Sciences Unit,  
Advanced Institute for Materials  
Research, Tohoku University

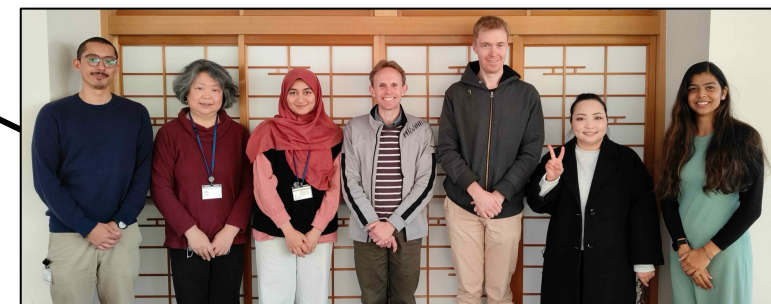


Assistant Professor

Institute for Integrated Cell-Material Sciences  
(iCeMS), Kyoto University

2016

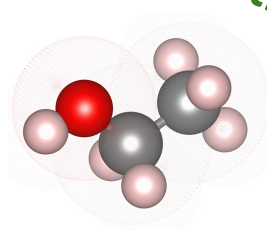
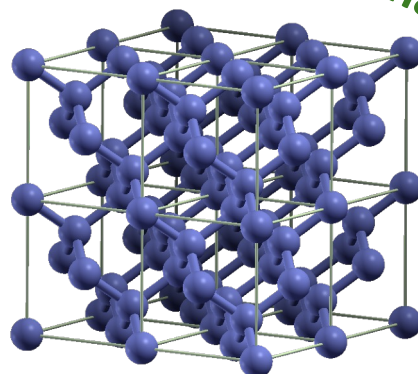
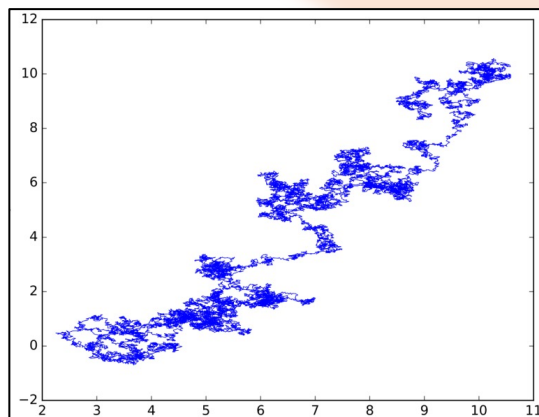
Lecturer



Associate Professor

Phase 1: Stochastic processes

Phase 2: Computational chemistry



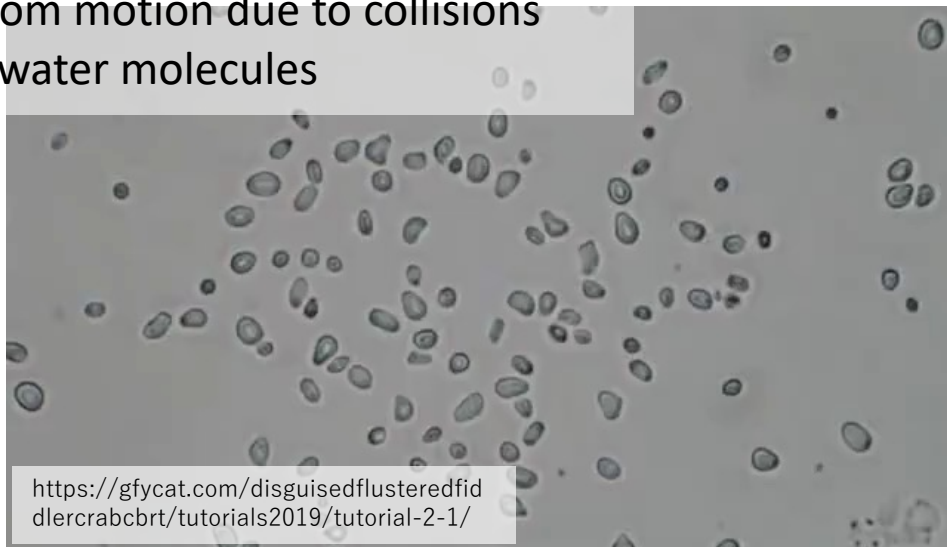
## *My previously research...*

I studied **stochastic processes** (random walks) in physics and chemistry.

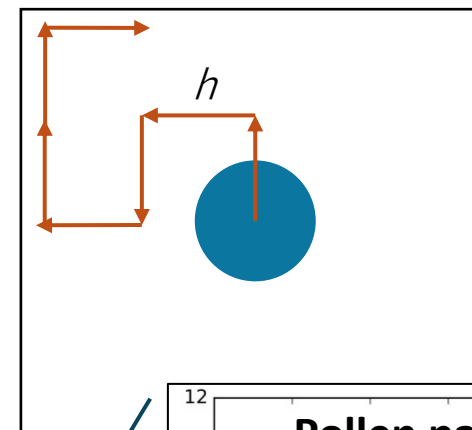
The most famous type of random walk in these fields is **Brownian motion**:

### **Pollen particles on water**

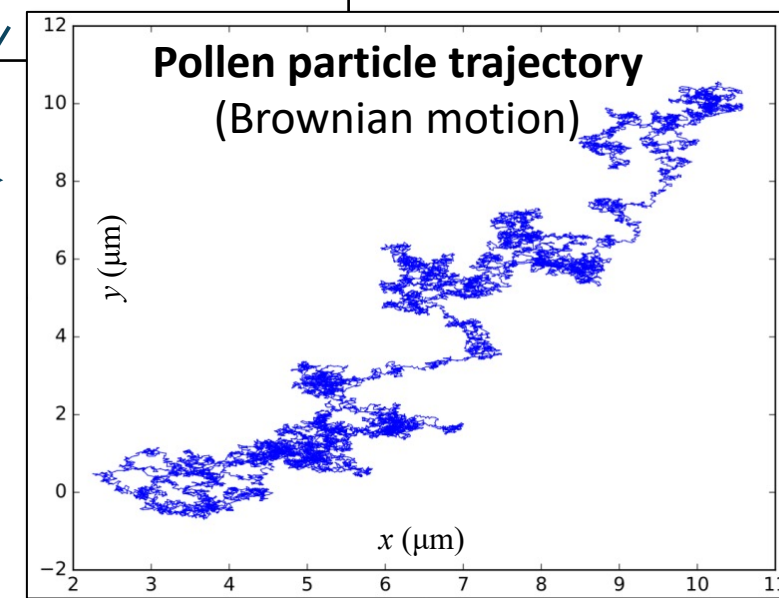
Random motion due to collisions  
with water molecules



Model motion as  
a random walk

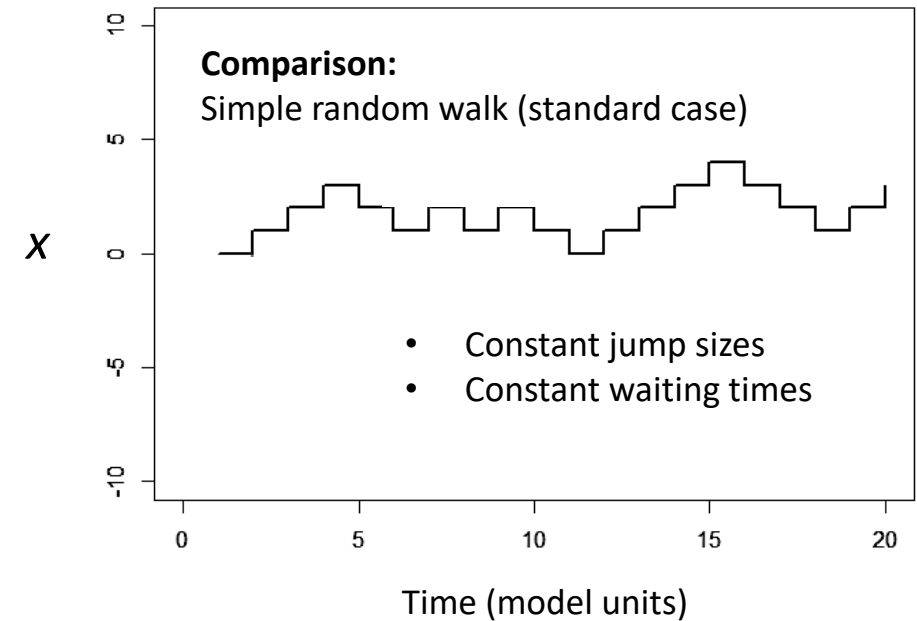
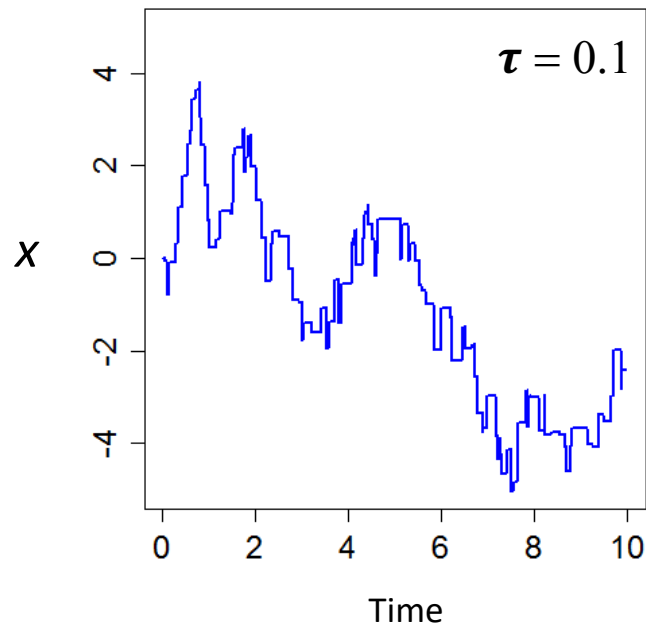
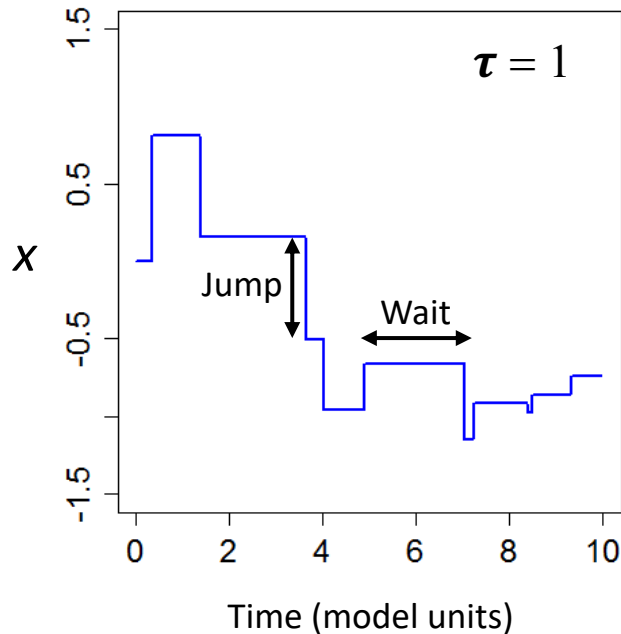


Take limit  
 $h \rightarrow 0$



# Postdoc research (2010 – 2012)

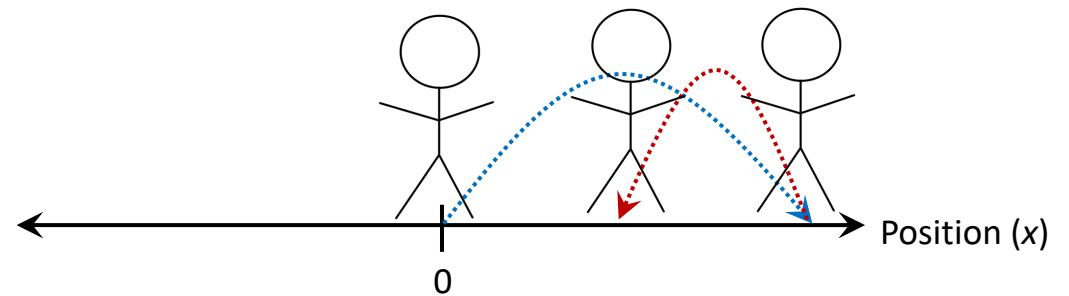
## Continuous-time random walk



Consider a person hopping in one dimension.  
Let  $x$  denote their position.

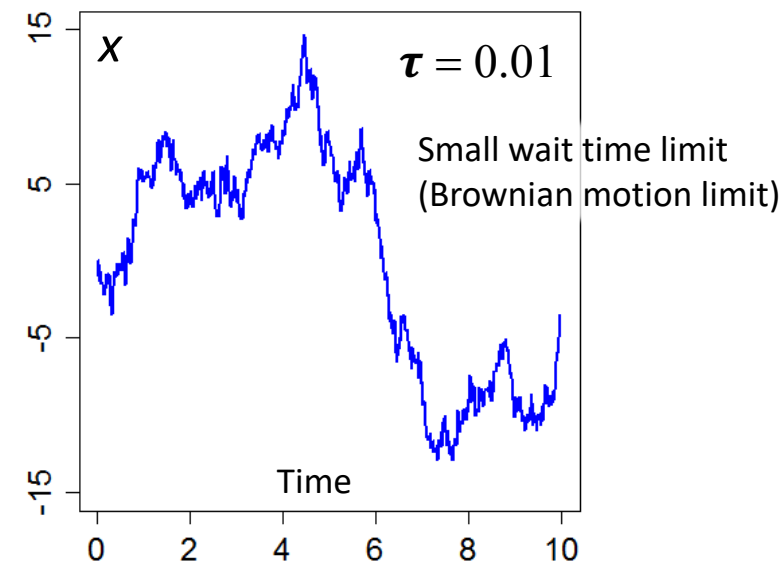
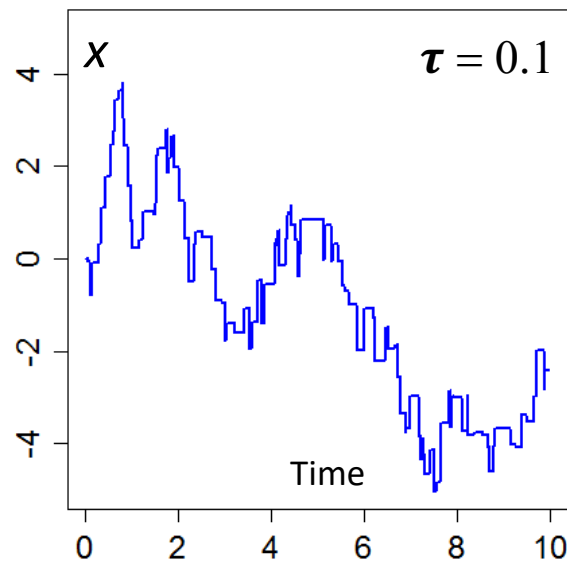
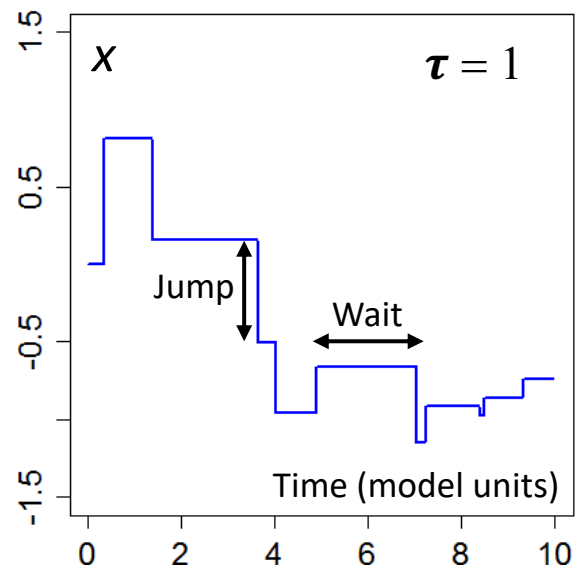
In the **continuous-time random walk**,

- the waiting times between jumps are random (average waiting time =  $\tau$ ), and
- the jump size is random





## Typical results



1. Derivation of conditions for convergence to Brownian motion as wait times go to zero.

(the maximum jump size  $M$  must become  $(3\tau)^{1/2}$ )

*J. Phys. A.: Math. Theor.* **43**, 2010, 464001

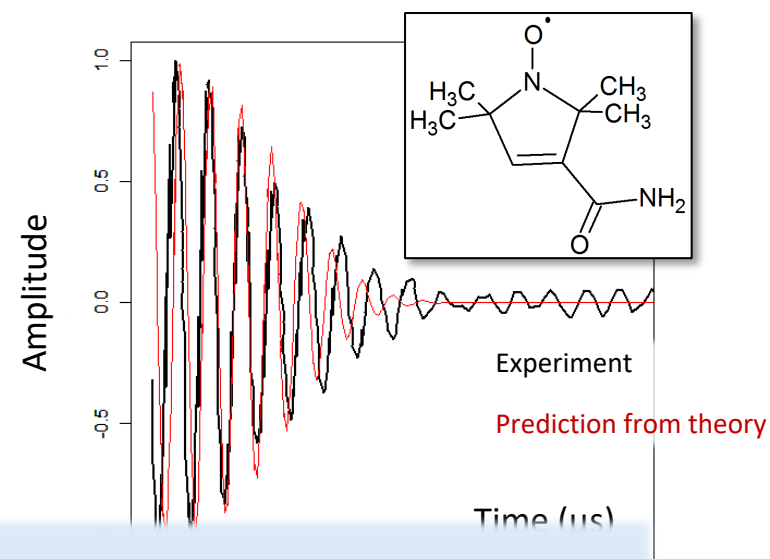
*arXiv*: 1105.6283 (2011)

*Phys. Rev. E.* **84**, 2011, 61111

2. Derivation of the linear response function (important in spectroscopy)

*Phys. Rev. E.* **86**, 2012, 11130

$$F(t) = e^{-t/\tau} e^{i\omega_0 t} \exp \left( \frac{1}{\tau M} \int_0^t \frac{\sin r/\tau}{r} dr \right)$$



Fun, but a very narrow topic...

# Career in Japan

Tanimura laboratory, Department of Chemistry, Grad. School of Science, Kyoto University

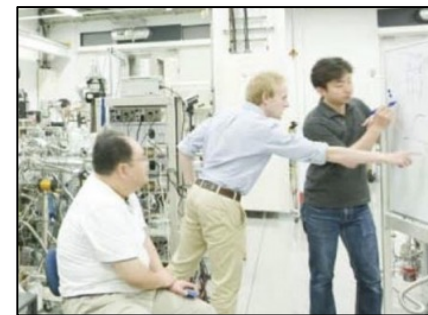


2010

Postdoc

2012

Mathematical Sciences Unit, Advanced Institute for Materials Research, Tohoku University

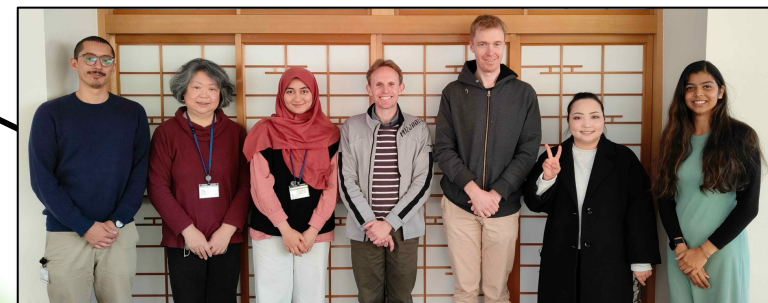


Assistant Professor

Institute for Integrated Cell-Material Sciences (iCeMS), Kyoto University

2016

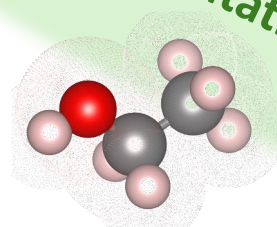
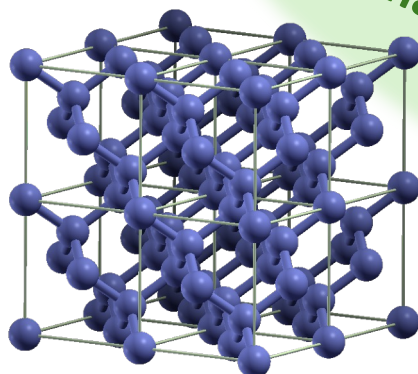
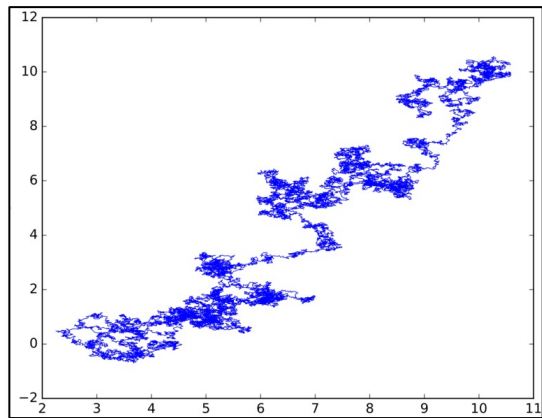
Lecturer



Associate Professor

Phase 1: Stochastic processes

Phase 2: Computational chemistry

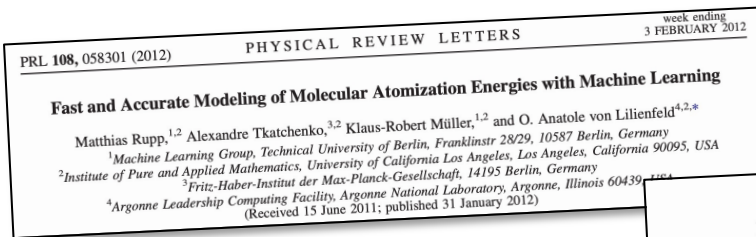




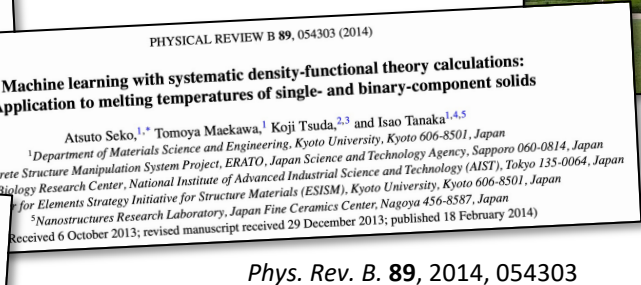
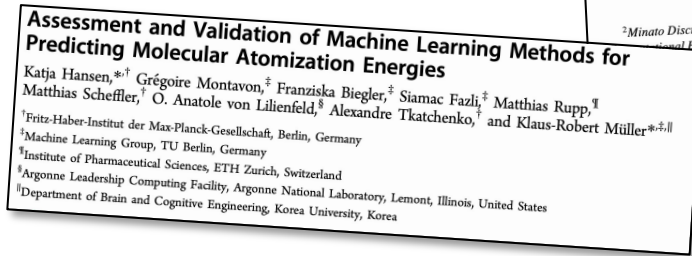
# Chance to shift to computational research:

## Emergence of data science in computational materials science

Around 2012 – 2014, several papers appeared showing how machine learning could be applied in computational materials science.



*Phys. Rev. Lett.* **108**, 2012, 058301



*Phys. Rev. B.* **89**, 2014, 054303

*J. Chem. Theory Comput.* **9**, 2013, 3404



### Materials for a Sustainable Energy Future

Long program at the Institute for Pure and Applied Mathematics, UCLA, Sep – Dec 2013

Moreover, in 2013 I got to spend time with some of the authors during a stay at University of California, Los Angeles (UCLA)

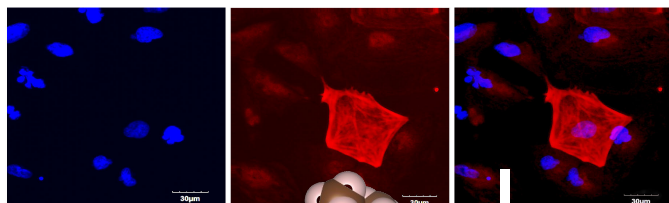
**Machine learning and data science gave mathematicians a way to enter the computational materials science field!**

... we just had to learn how to do density functional theory calculations.

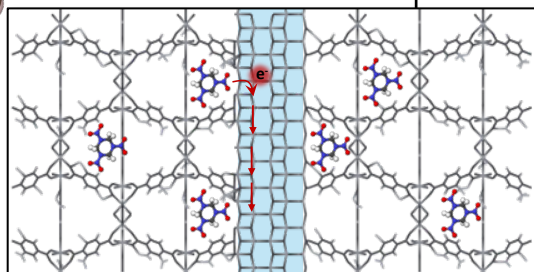
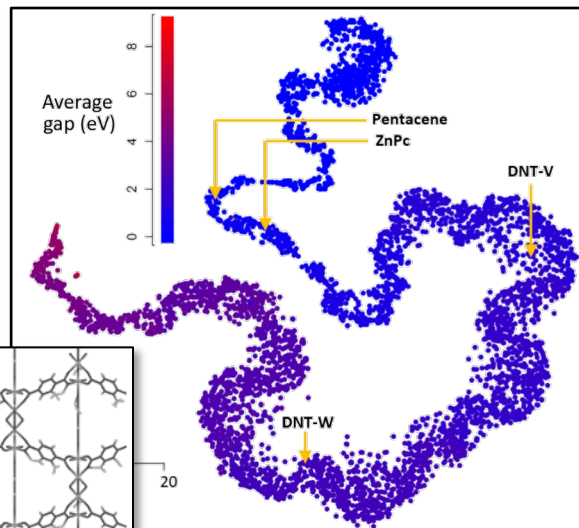
# Research in the Packwood group

organic materials x simulation x data science

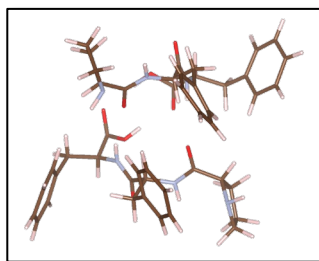
Computational design of bioactive compound for iPS cell control (*J. Chem. Inf. Model.* 2025)



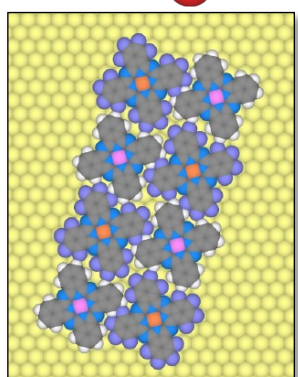
Visualization of organic semiconductor database (*Adv. Theory Simul.* 2023)



Disease detection potential in a MOF-semiconductor sensor (*Adv. Theory Simul.* 2025)



Game theory analysis of peptide dimer dynamics (in process)

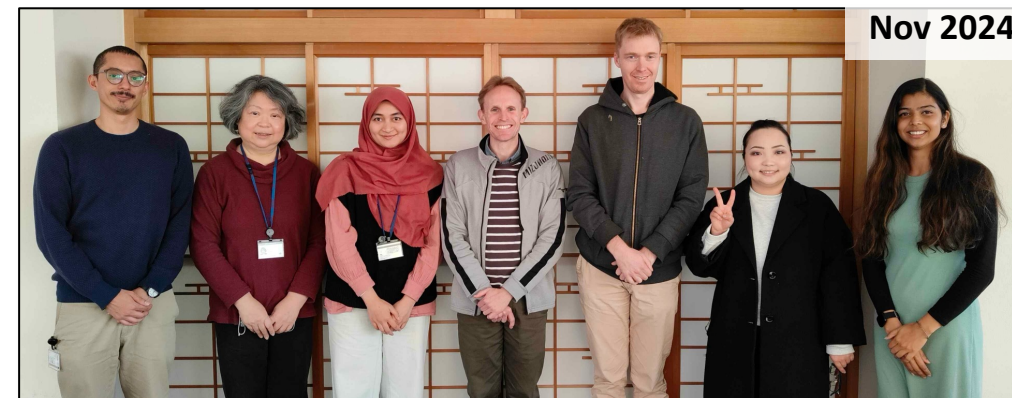


Simulation of metal-complex self-assembly (*Adv. Physics Res.* 2022)

Coordination polymers / metal-organic frameworks (2018 – now)

Organic semiconductor (2021 – now)

On-surface molecular self-assembly (2015 – 2022)



Nov 2024

Small molecule aggregation (2021 – now)

Bioactive molecule discovery (2021 – now)

Chemical biology

Materials science

2015

2025

## Lecture topics

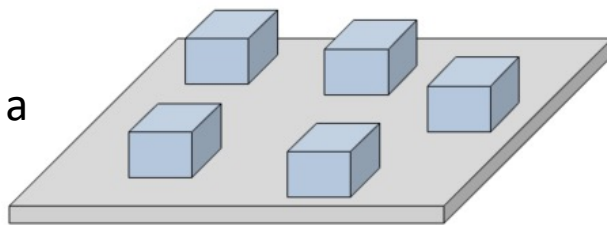
Simulation of on-surface molecular self-assembly

Machine learning for organic photovoltaic materials

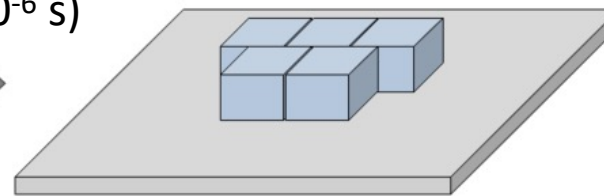


# On-Surface Molecular Self-Assembly

Molecules  
adsorbed to a  
surface



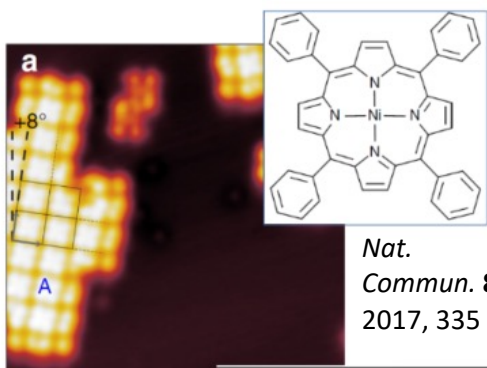
Time ( $< 10^{-6}$  s)



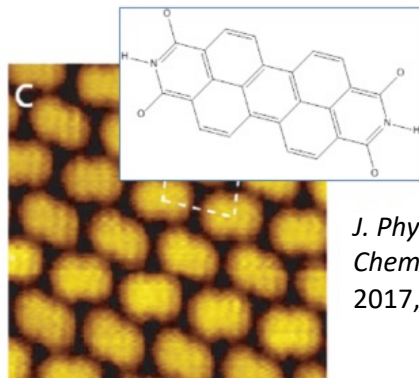
Molecular assembly  
(forms spontaneously)

## Many experimental reports...

(Scanning tunneling microscopy images)

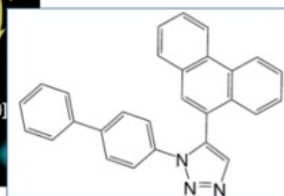
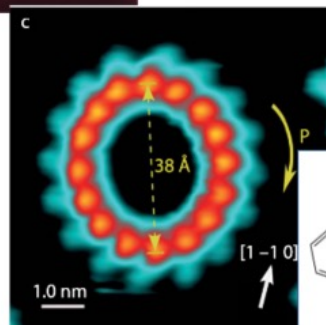


*Nat. Commun.* **8**,  
2017, 335

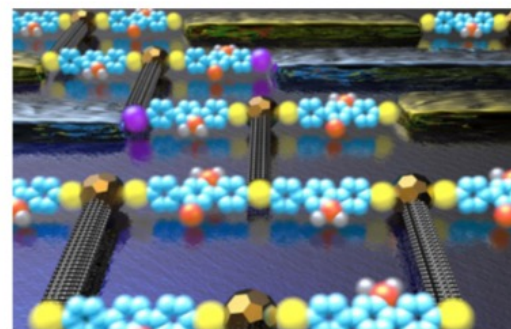


*J. Phys. Chem. C.* **121**,  
2017, 20986

*ACS Nano* **11**,  
2017, 8302



## ... and many potential applications



Molecular electronics

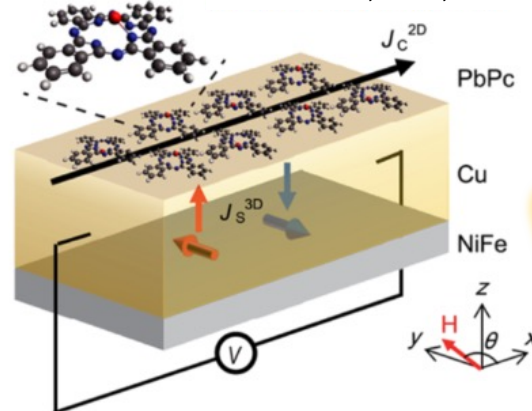
[http://asdn.net/asdn/electronics/molecular\\_electronics.php](http://asdn.net/asdn/electronics/molecular_electronics.php)

Organic electronics

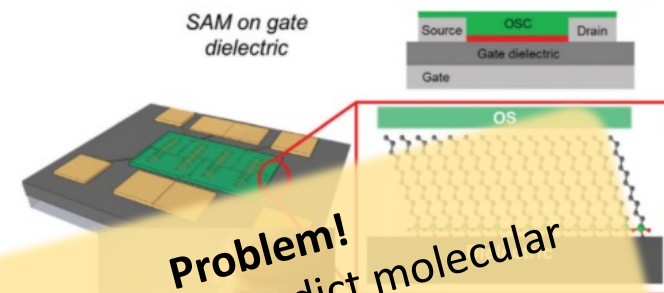
*Chem. Soc. Rev.* **46**, 2017, 40

Spintronics

*Nano. Lett.* **19**, 2019, 7119



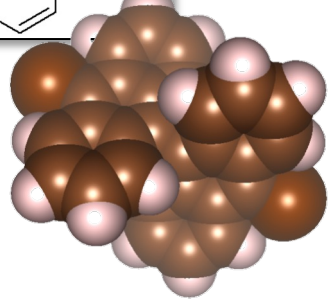
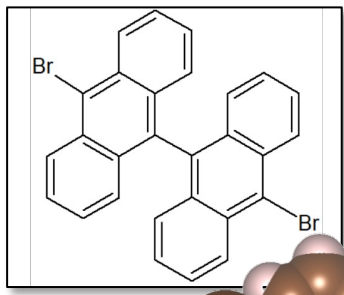
SAM on gate  
dielectric



**Problem!**  
Very hard to predict molecular  
assembly

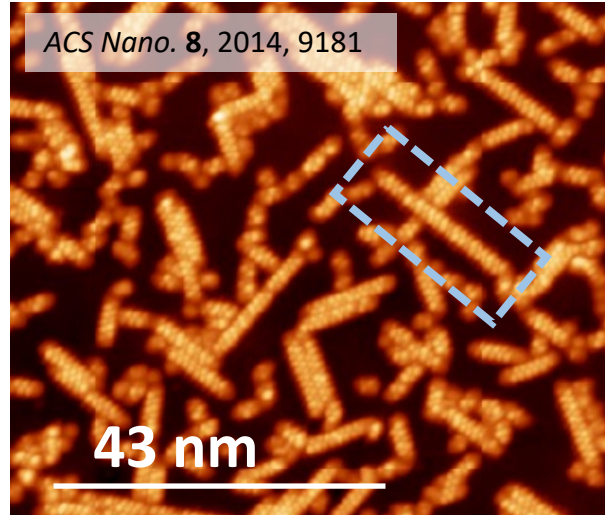
<https://www.scinotsci.com/electronics>

# How I came to this topic?



**Dibromo-bianthracene molecule**

Deposition on  
Cu(111)



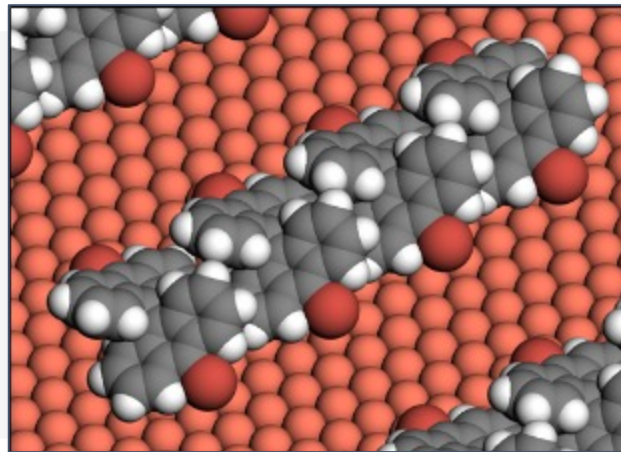
**Self-assembly of molecular chains**

About 10 years ago, colleagues at Tohoku University were studying molecular self-assembly using scanning tunneling microscopy (STM).

They wanted to know whether the chains could be predicted computationally.

Collaborating together, we created a computational method to predict chain formation.

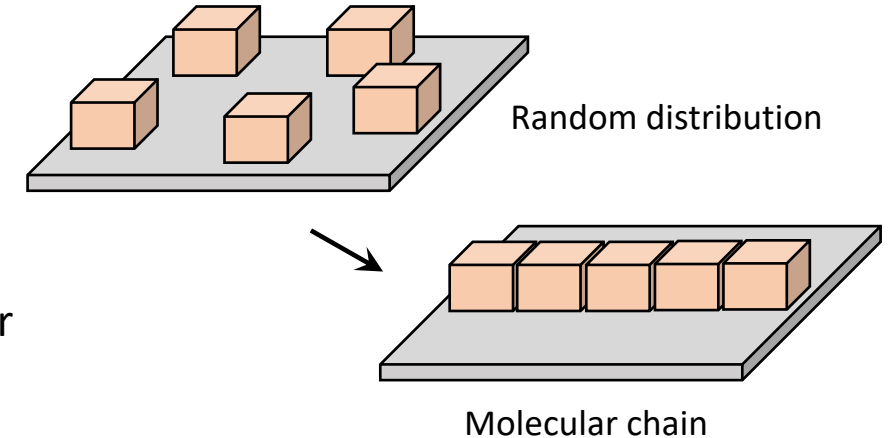
*Nat. Commun.* **8**, 2017, 14463



**Dr. Patrick Han**



**Prof. Taro Hitosugi**  
(now at Tokyo Univ)

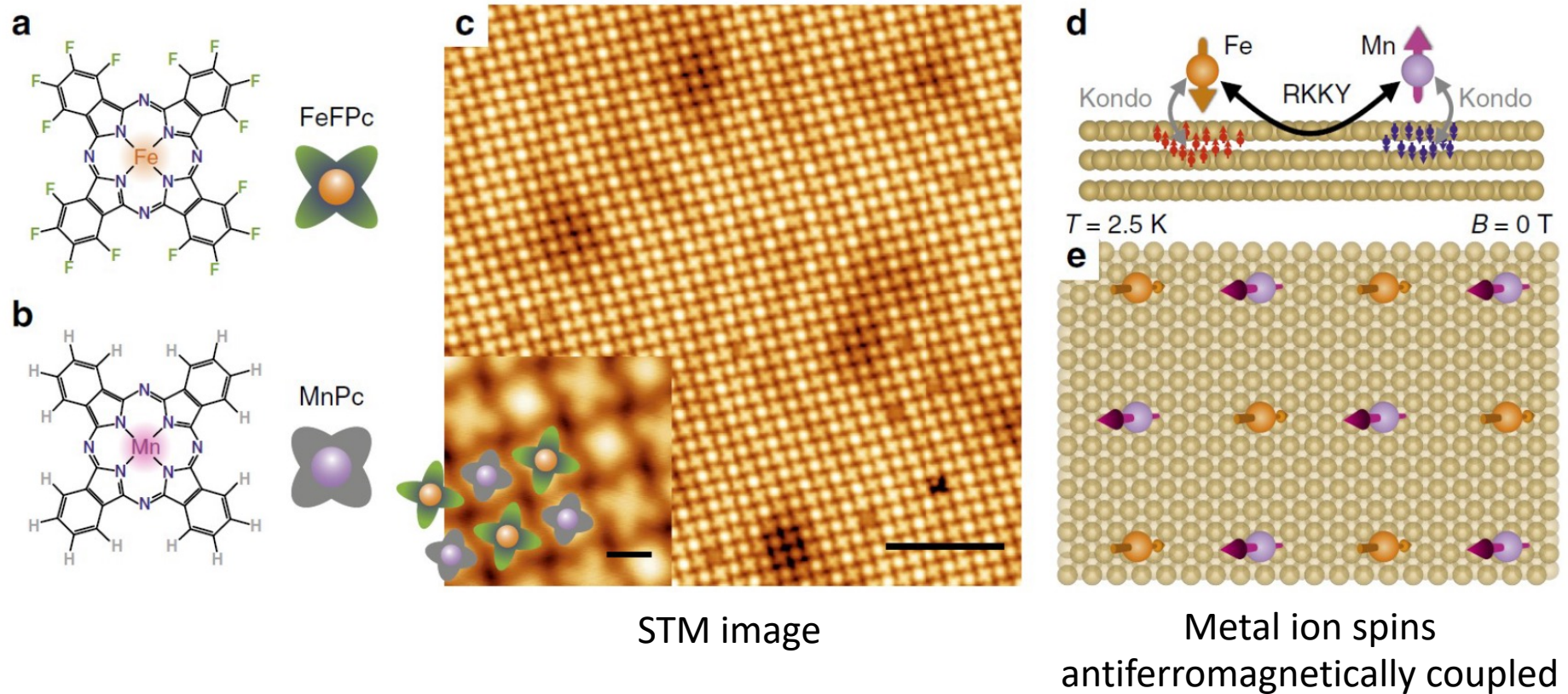


**Today, I will discuss an updated version of this method applied to a different type of system...**



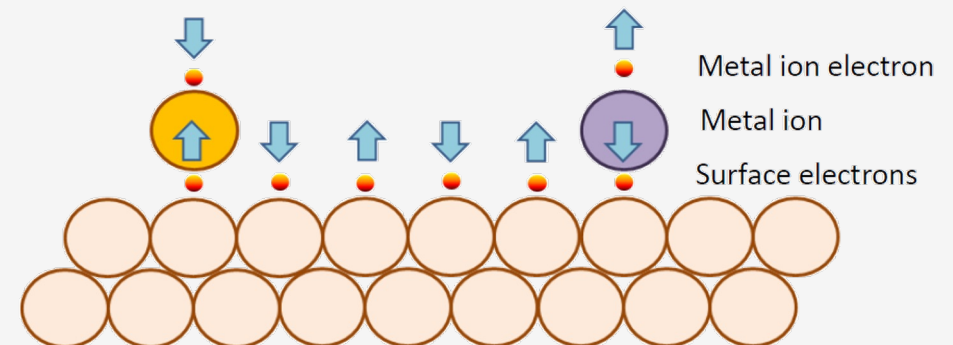
# Self-assembly of FeFPc and MnPc on Au(111)

Girovsky et al. Nat. Commun. 8, 2017, 15388

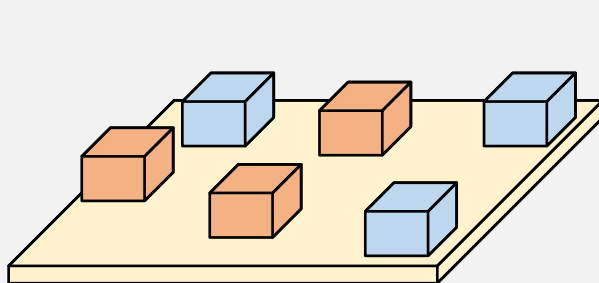


The antiferromagnetic coupling is due to the Ruderman-Kittel-Kasuya-Yoshida (RKKY) interaction

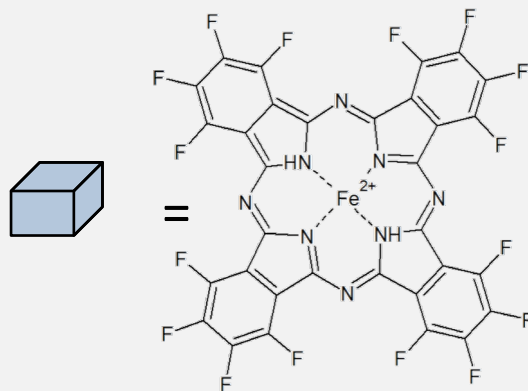
***Can we reproduce this result with computation?***



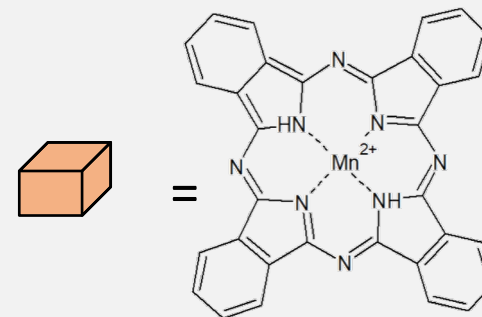
## Q1. How to build a simplified model for the system?



Metal complexes on gold surface (Au(111))

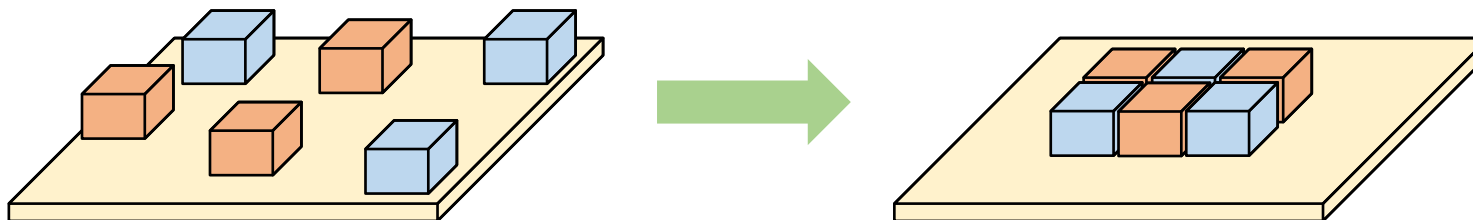


FeFPc



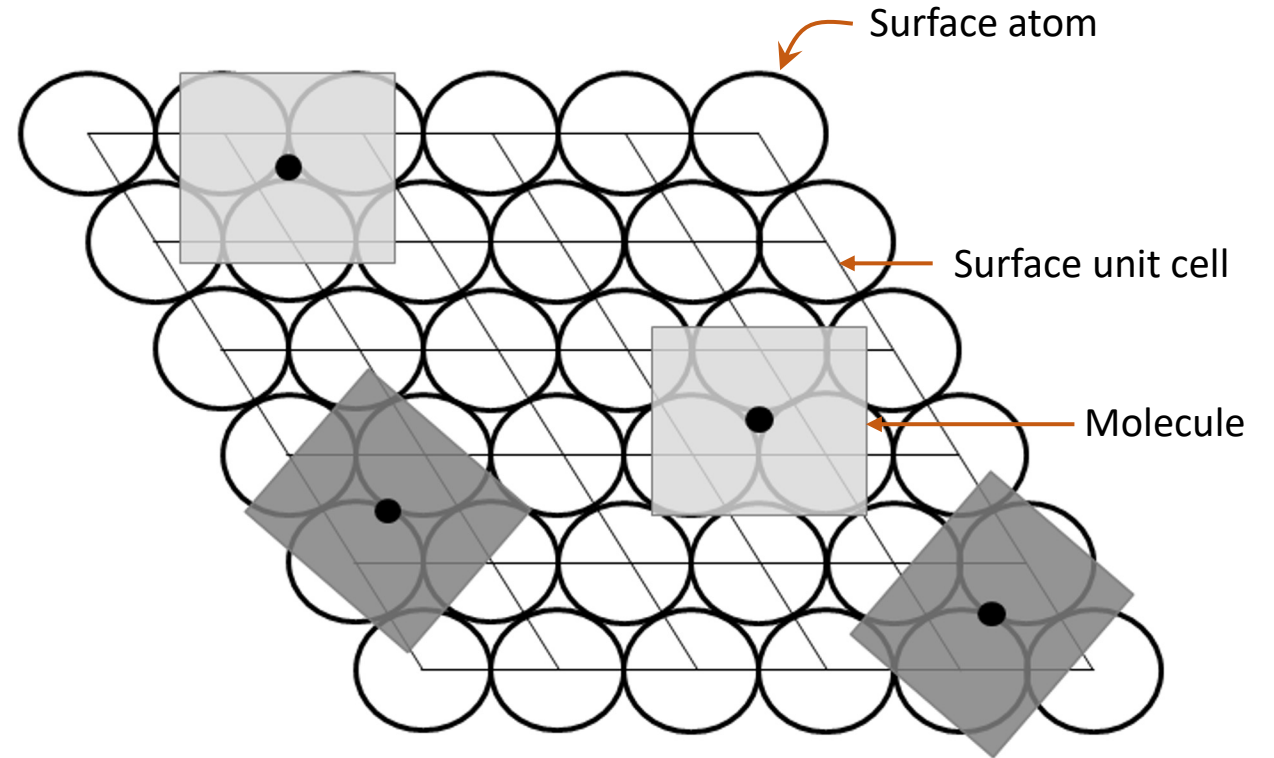
MnHPc

## Q2. How to obtain predictions about self-assembly from the model?



# Model assumptions

- Perfectly crystalline surface
- Two types of molecules on the surface;  $n_1$  of the first type,  $n_2$  of the second type.
- Finite number of adsorption sites (places where the molecules can sit)
- Finite number of molecule orientations.
- Rigid molecules. All molecules of the same type have the same conformation.
- Two-body energy function



$$E = \sum_i u_i + \sum_{i \neq j} v_{ij}$$

$u_i$  = molecule  $i$ -surface interaction energy

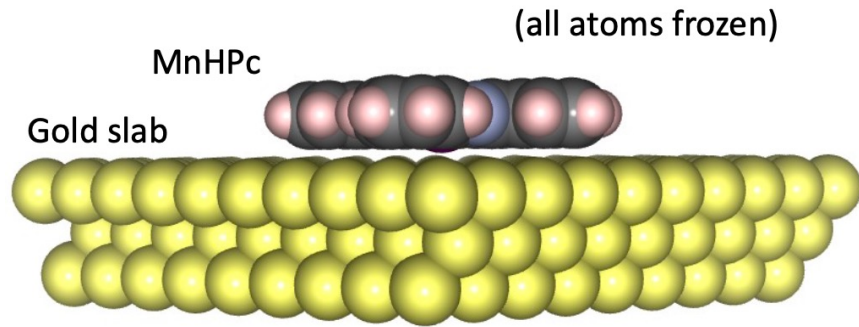
$v_{ij}$  = molecule  $i$ -molecule  $j$  interaction energy

*How do we assign the energy parameters ( $u_i$ ,  $v_{ij}$ )?*

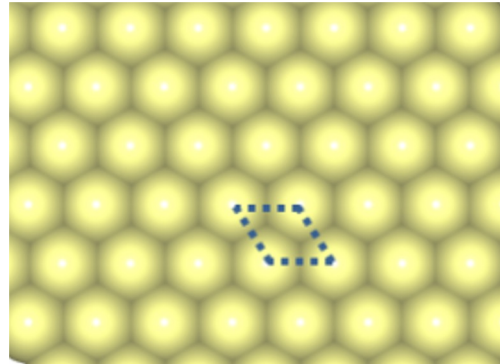


# Molecule-surface interaction energies – assign using DFT\*

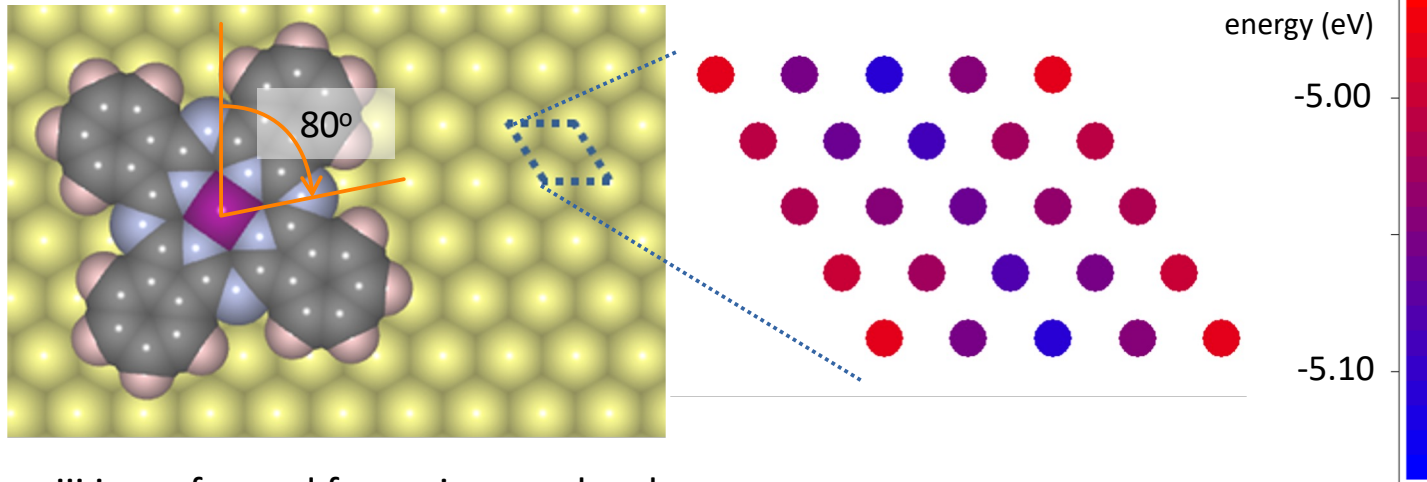
i. Obtain the optimal molecule-surface adsorption height



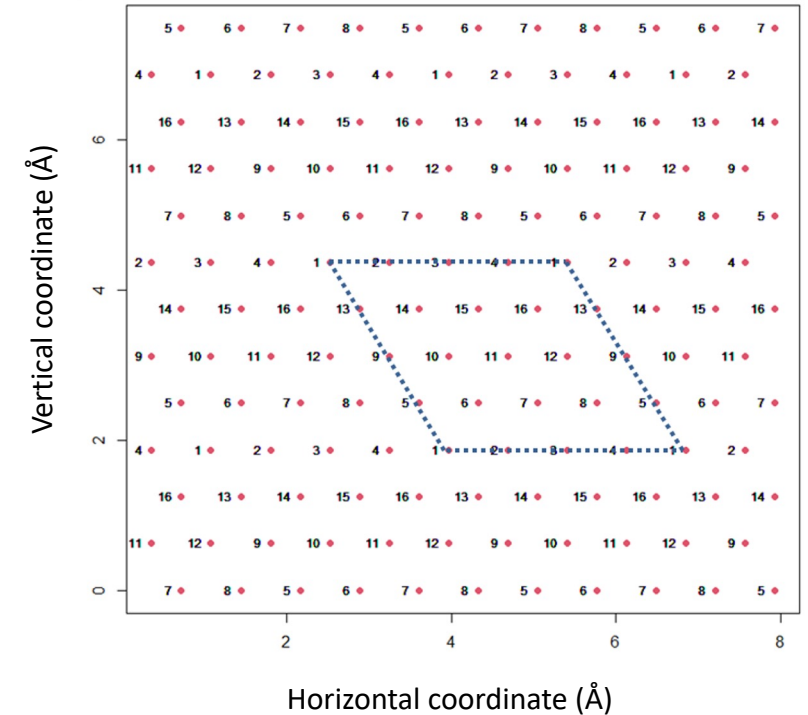
ii. Discretize surface (around 40,000 points)



iii. Compute interaction energy at each (symmetry unique) adsorption site



Step iii is performed for various molecule orientations (0°, 40°, 80°, ..., 320°)



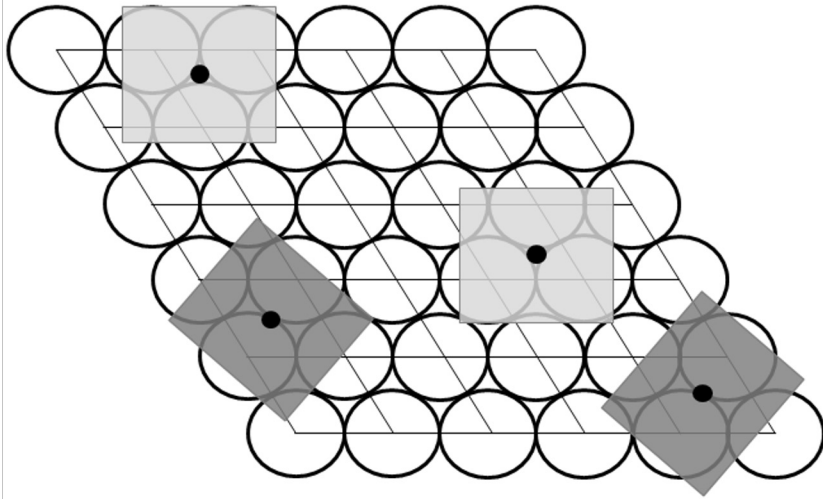
→ *Surface-molecule interaction parameters assigned!*

$$E = \sum_i u_i + \sum_{i \neq j} v_{ij}$$

\* DFT as implemented in FHI-aims, with the PBE exchange-correlation functional, TS vdW corrections, and “light” basis set defaults.

# Molecule-molecule interaction energies

$$E = \sum u_i + \sum \underset{\text{.....}}{v_{ij}}$$



## Problem!

For a model with 40,000 adsorption sites, around  $10^{11}$  unique molecule-molecule interactions are possible.

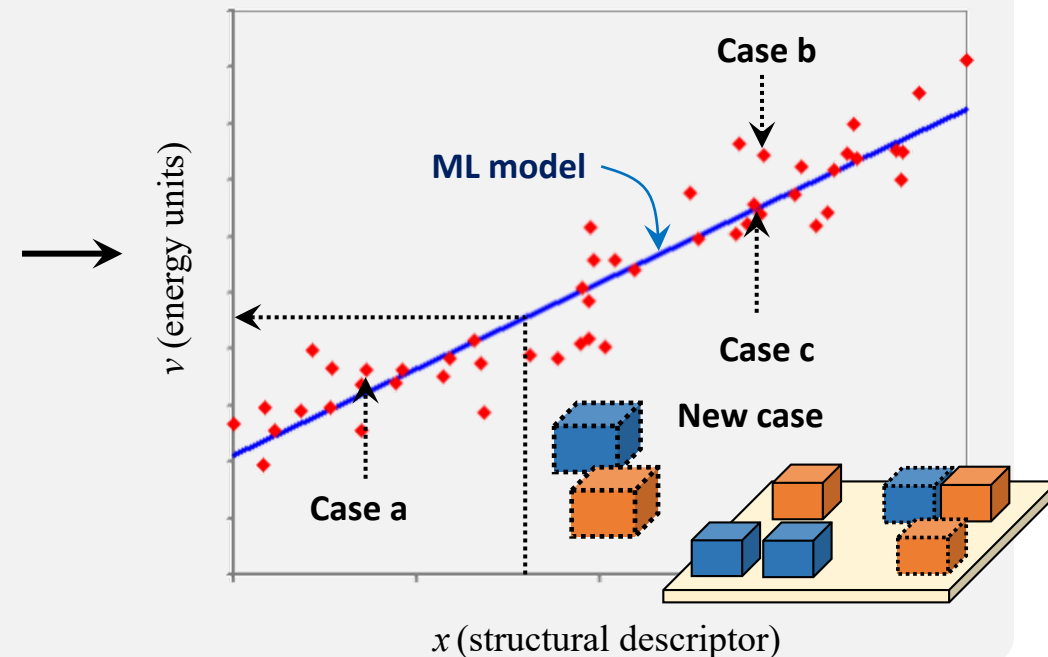
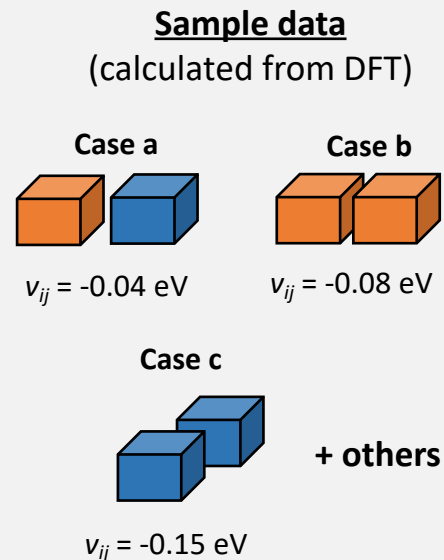
We cannot calculate the energy of each one individually using DFT.

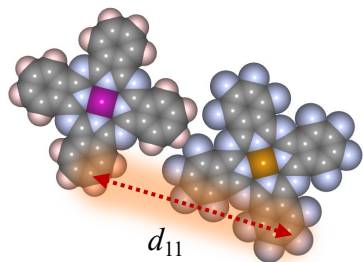
## Our solution

Take a small sample of around 5000 intermolecular interactions, calculate their energies with DFT.

Use the sample data to build a **machine-learning (ML) model**.

Use the model to assign the energies of the remaining cases.

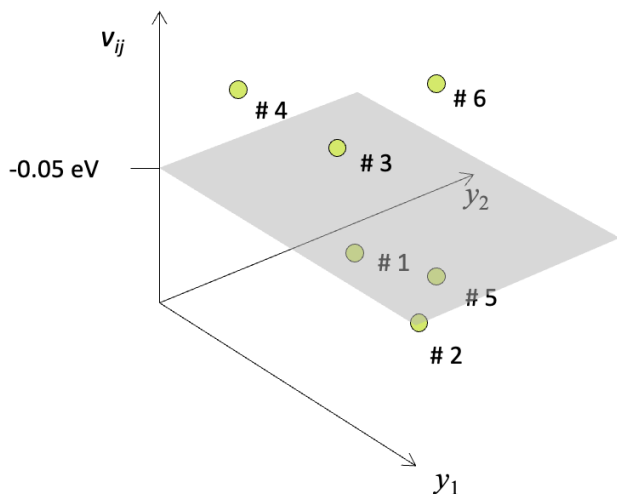




### Coulomb-type descriptors

$$x_1 = 1/d_{11}, x_2 = 1/d_{12}, \dots$$

**Input case** (interaction energy =  $v_{ij}$ )



### Support vector machine

Transform data so that it can be separated by a linear plane  
 $(x_1, x_2, \dots) \Rightarrow (y_1, y_2, \dots)$

$$v_{ij} < -0.05 \text{ eV}$$

$$-0.05 \text{ eV} \leq v_{ij} \leq 0$$

### Kernel ridge regression

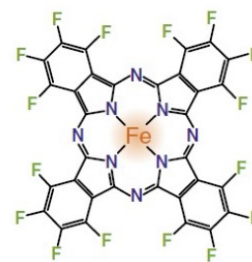
Transform data so that it lies on a linear plane  
 $(y_1, y_2, \dots) \Rightarrow (z_1, z_2, \dots)$

$v_{ij}$  prediction

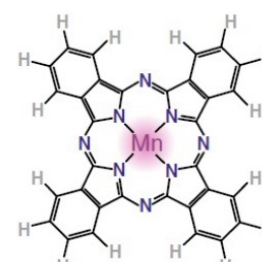
$v_{ij} = 0$

### System under consideration:

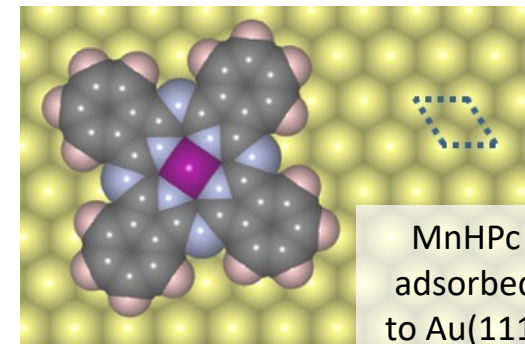
Phthalocyanine molecules adsorbed to gold(111)



FePc

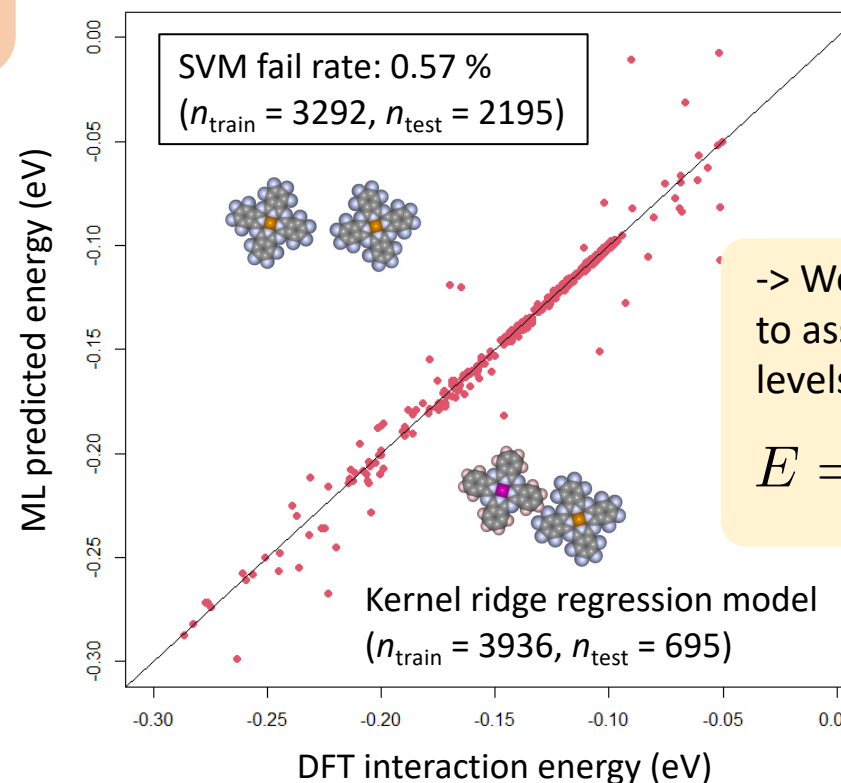


MnHPc



MnHPc  
adsorbed  
to Au(111)

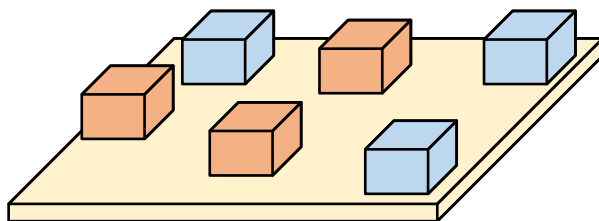
### Performance on test data



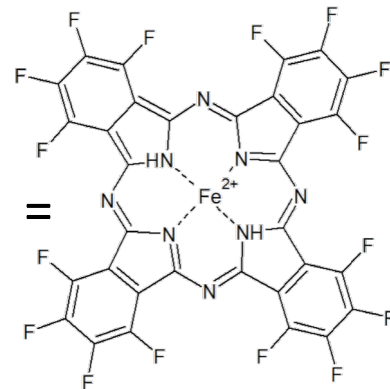
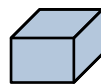
-> We can use the ML model to assign the  $v_{ij}$  with DFT-levels of accuracy:

$$E = \sum u_i + \sum v_{ij}$$

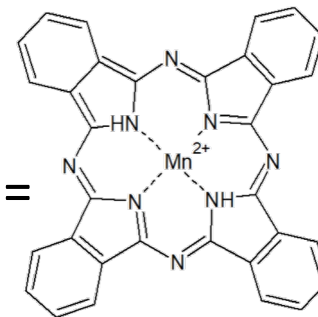
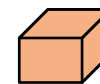
## Q1. How to build a simplified model for the system?



Metal complexes on gold surface (Au(111))

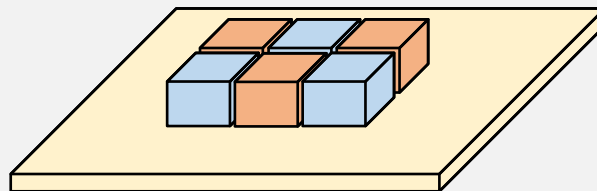
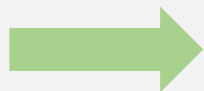
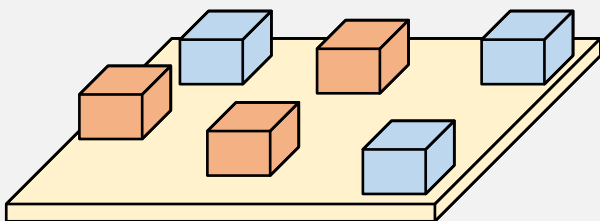


FeFPC



MnHPc

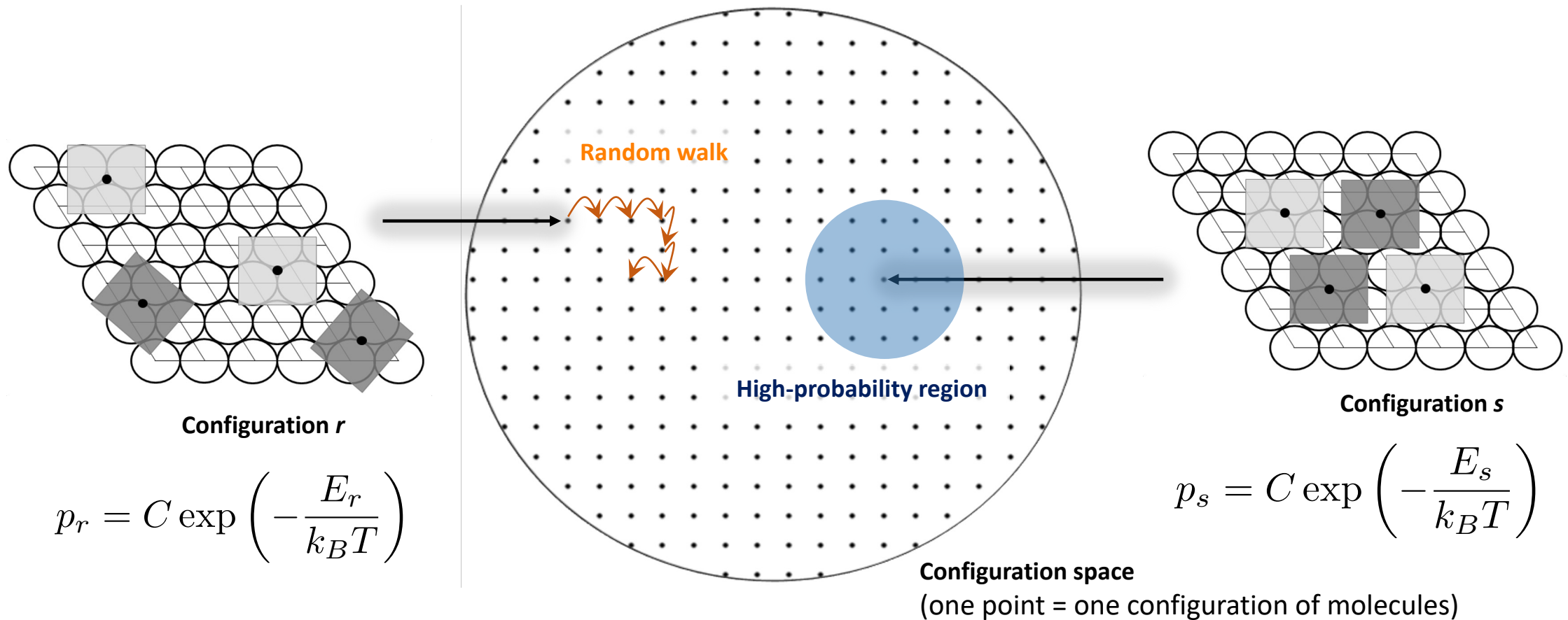
## Q2. How to obtain predictions from the model?



We need to identify molecule configurations with high formation probability  $p$ :

$$p = C \exp \left( -\frac{E}{k_B T} \right)$$

# Markov chain Monte Carlo (MCMC)

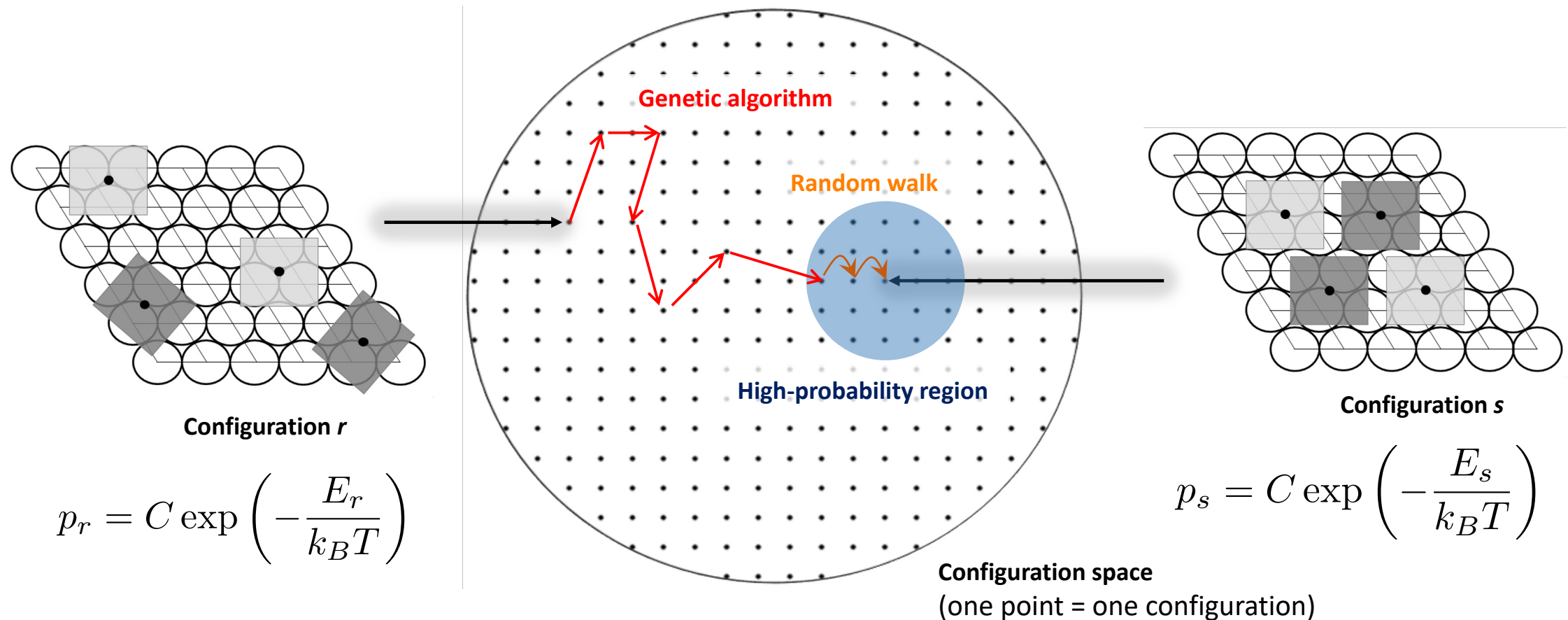


In MCMC, we simulate a random walk over the configuration space. It is simulated in such a way that, after a long length of time, the number of visits to configuration  $s$  is proportional to  $p_s$  (the formation probability).

**Problem:** MCMC can be very inefficient. The configuration space is huge. The high-probability region (where the molecules are closely packed) is very small. Long simulation times are usually required to reach it.



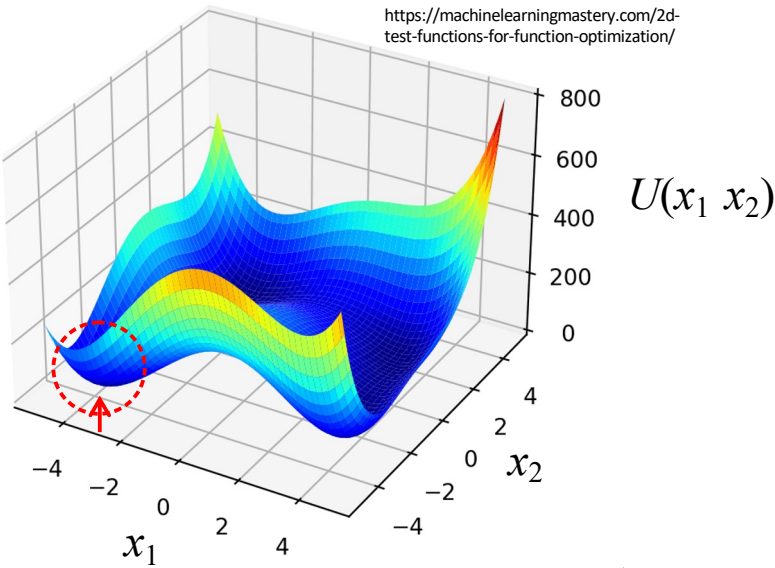
## Our solution: genetic algorithm + Markov chain Monte Carlo



The genetic algorithm makes large steps in the configuration space, quickly bringing us to the high-probability region (coarse search)

Markov chain Monte Carlo (random walk) makes short steps, thoroughly exploring the high-probability region. (fine search)

# Genetic algorithm (general concept)



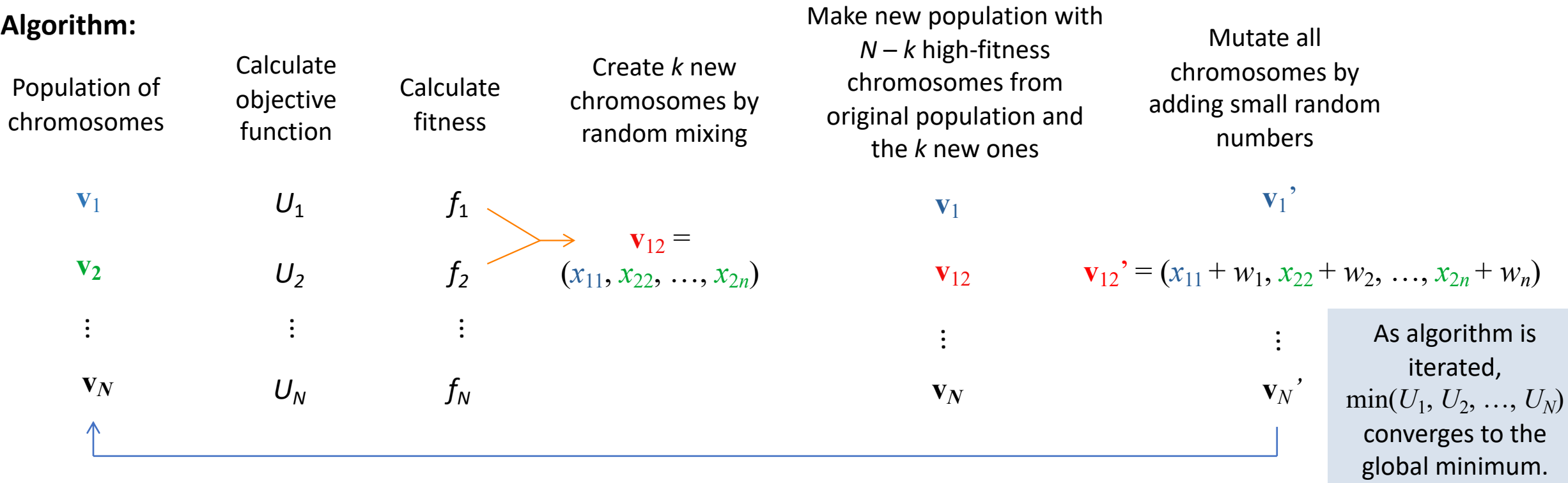
**Task:** Find the values of  $x_1, x_2, \dots, x_n$  which minimize the objective function  $U(x_1, x_2, \dots, x_n)$

**Initial setup:**  $N$  vectors of random numbers  $(x_{ij})$

$$\mathbf{v}_1 = (x_{11}, x_{12}, \dots, x_{1n}) \quad \mathbf{v}_2 = (x_{21}, x_{22}, \dots, x_{2n}) \quad \dots \quad \mathbf{v}_N = (x_{N1}, x_{N2}, \dots, x_{Nn})$$

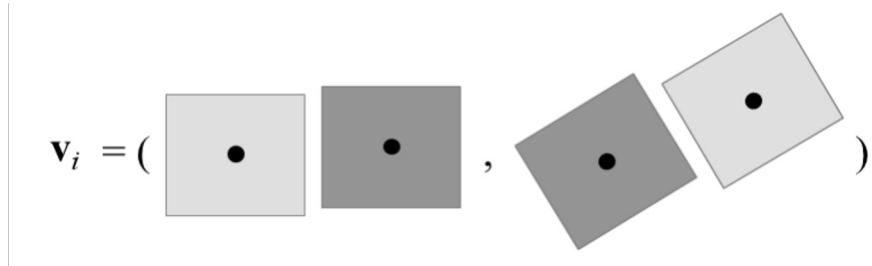
These vectors are called *chromosomes*. The elements  $x_{ij}$  are called *genes*.

## Algorithm:

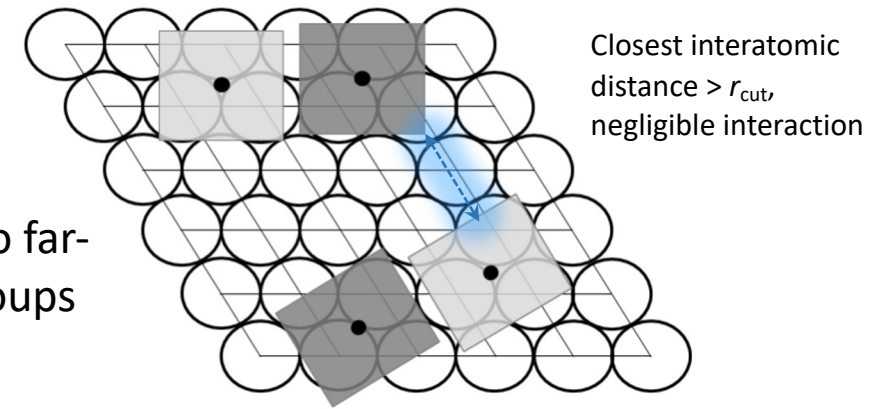


# Genetic algorithm implementation

**Chromosomes:** Vector of molecular clusters



Clusters correspond to far-separated groups of molecules



**Objective function:** a free energy function

$$U(\mathbf{v}_i) = E(\mathbf{v}_i) - JS(\mathbf{v}_i)$$

↑  
Energy

↑  
Pseudo-temperature

$$S(\mathbf{v}_i) = k_B \ln W(\mathbf{v}_i)$$

↑  
Configuration entropy

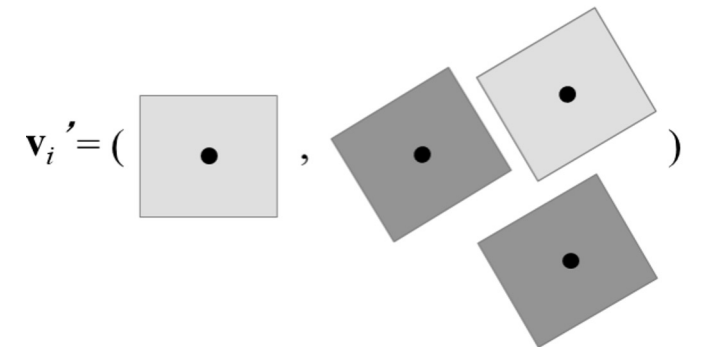
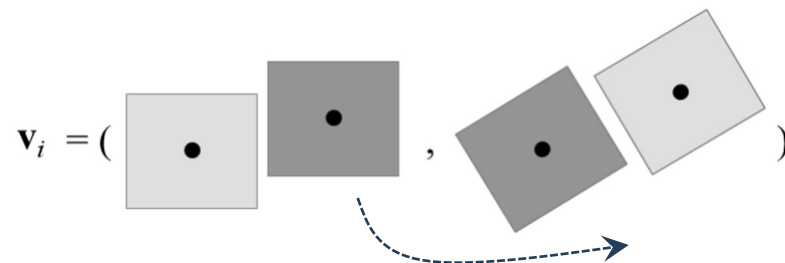
↑  
Number of unique ways molecular clusters can be placed on the surface (approximate formula: *R. Soc. Open. Soc.* **3**, 2016, 150681)

**Chromosome mixing:**

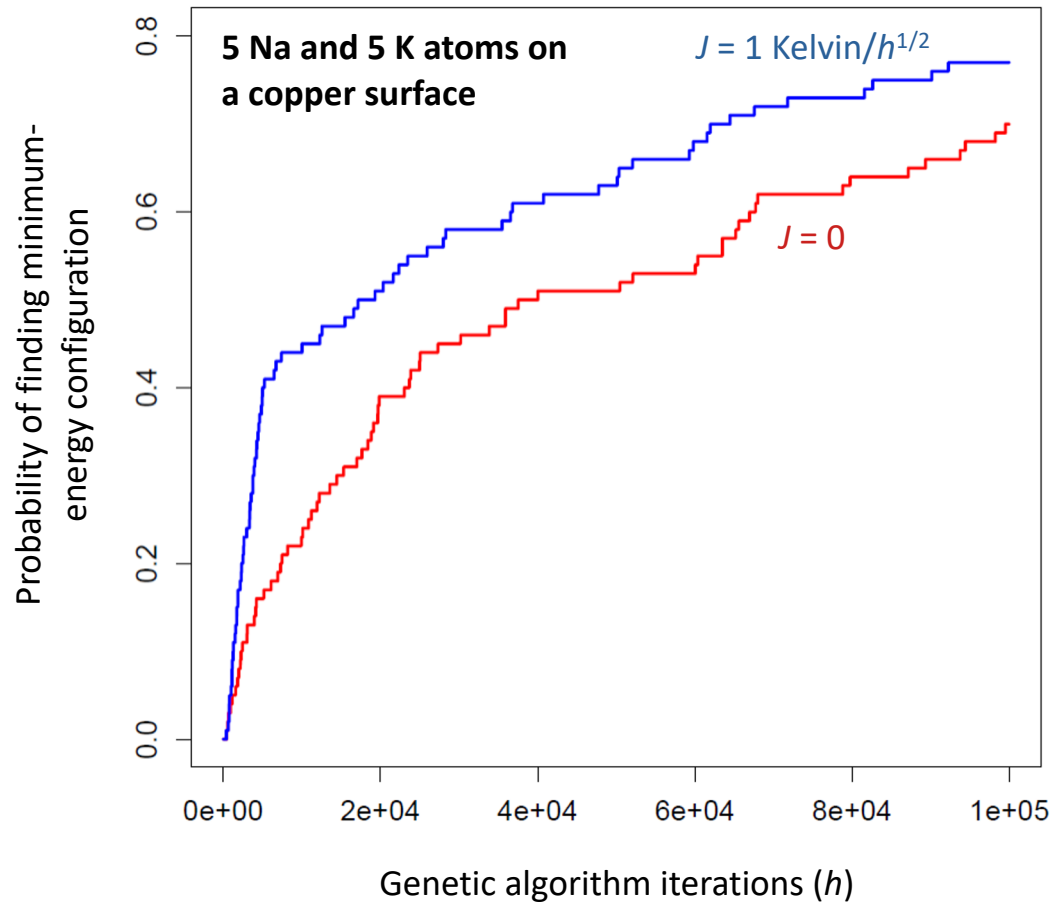
As before, but with conditions (to ensure that number of molecules  $n_1, n_2$ , is constant)

**Chromosome mutations:**

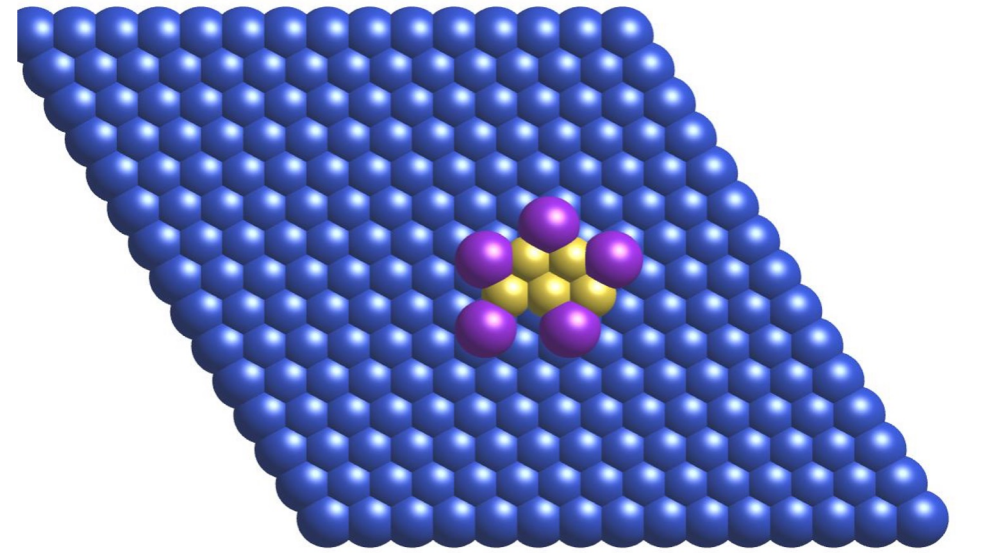
Random shift of a molecule between clusters



## Special feature of our genetic algorithm: simulated annealing



Minimum-energy configuration  
(yellow = Na atom, purple = K atom)



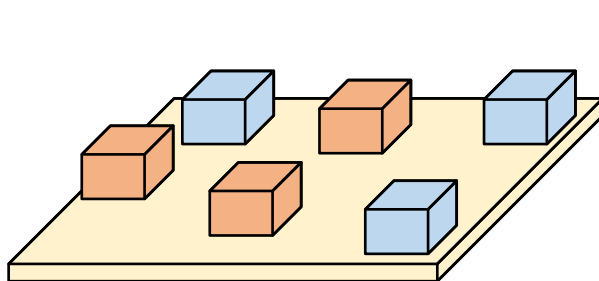
We gradually decrease the pseudo-temperature  $J$  as the algorithm proceeds. This helps the algorithm move out of local energy minima, improving performance.

**Note:**  $J$  is just a parameter to control the algorithm, and not the true surface temperature.

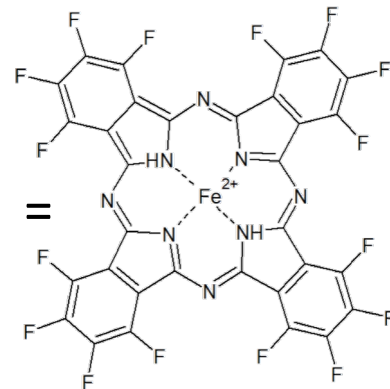
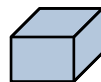
$$U(\mathbf{v}_i) = E(\mathbf{v}_i) - JS(\mathbf{v}_i)$$



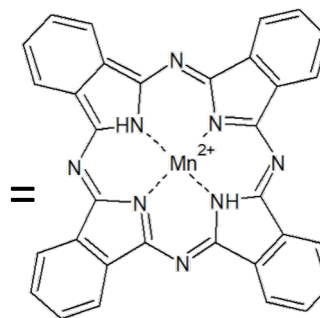
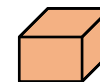
## Q1. How to build a simplified model for the system?



Metal complexes on gold surface (Au(111))

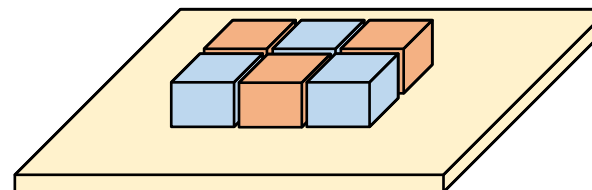
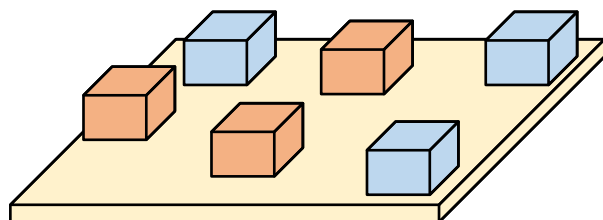


FePc



MnHPc

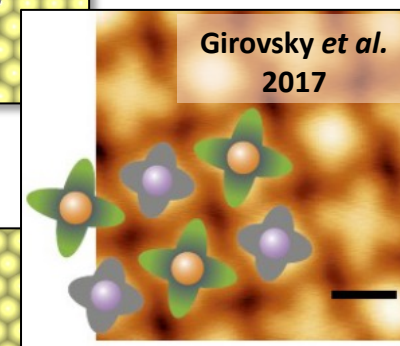
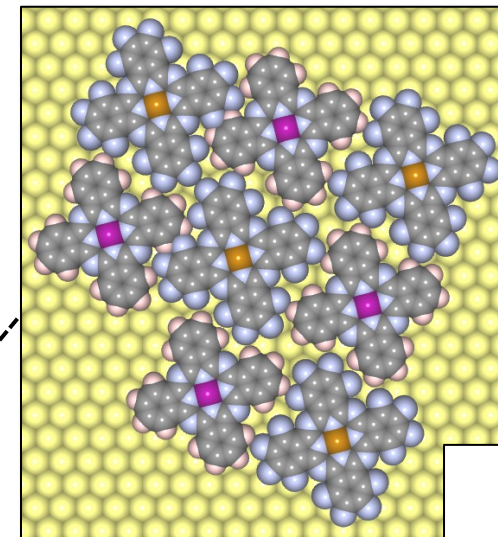
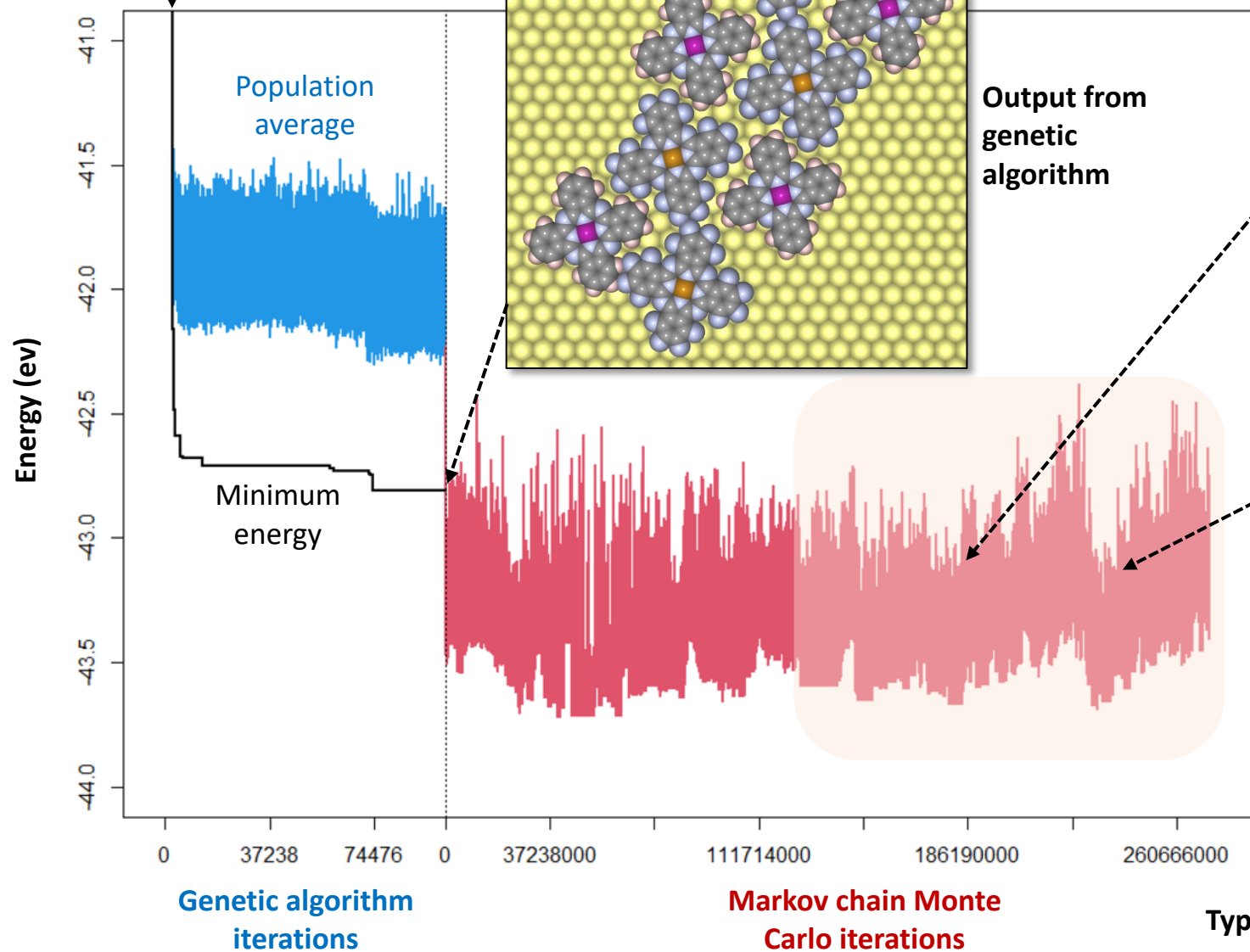
## Q2. How to obtain predictions from the model?



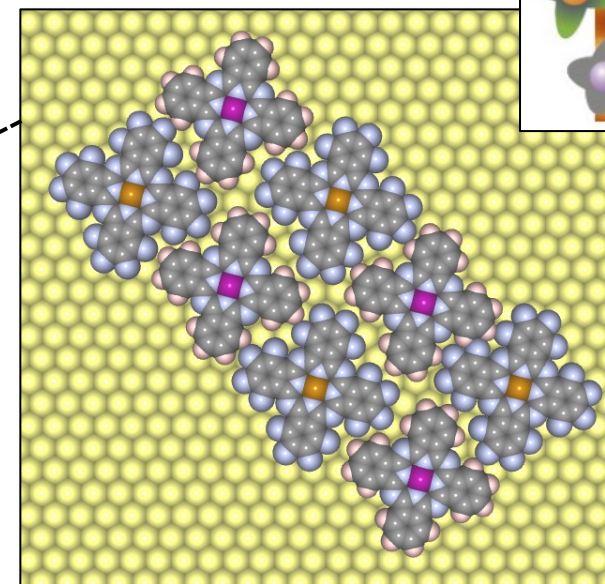
Results...

# Simulation for 4 FeFPc and 4 MnPc molecules on Au(111)

Initial: random dispersion of molecules

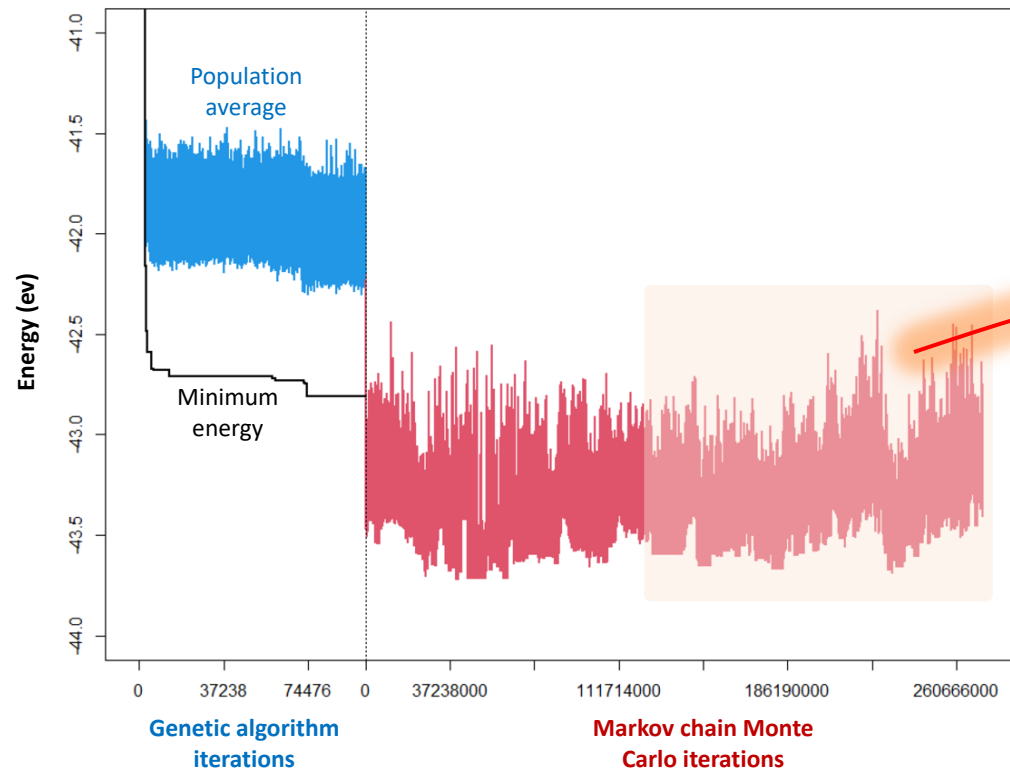


Predictions look similar to STM image?

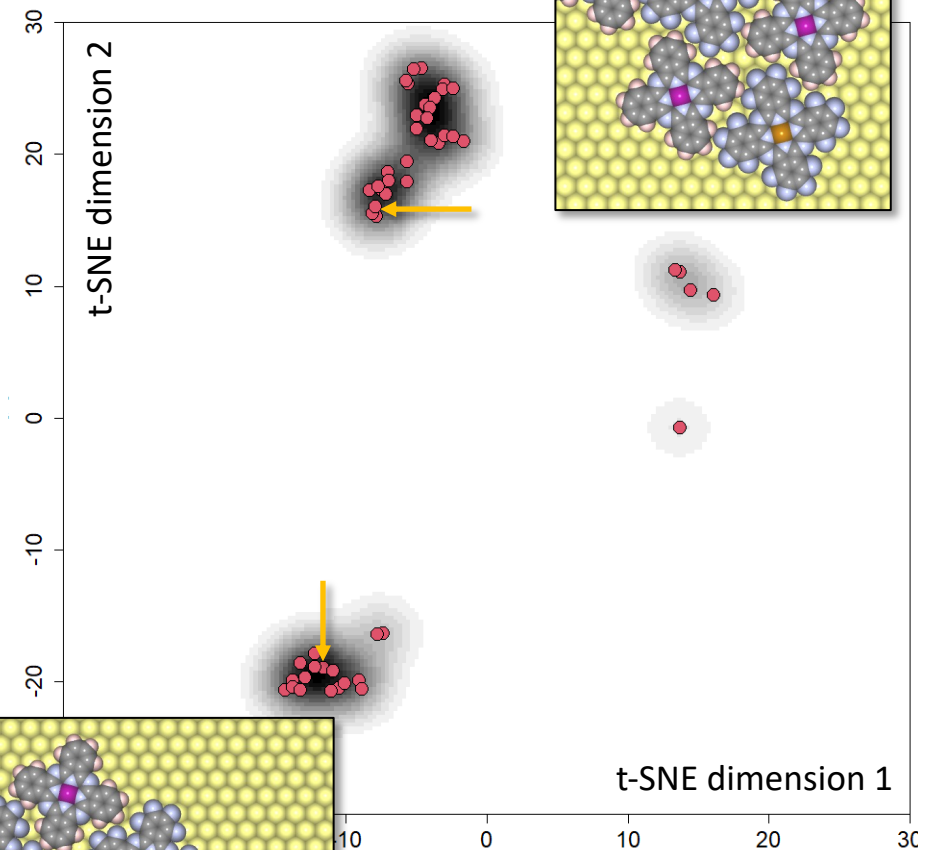


Typical assemblies observed during Monte Carlo (at 300 K)

# Visualization of equilibrium distribution of molecular configurations (300 K)



50  
configurations  
from MCMC

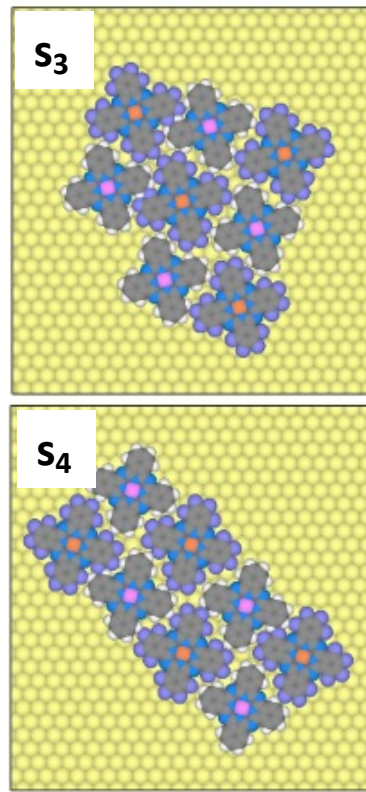
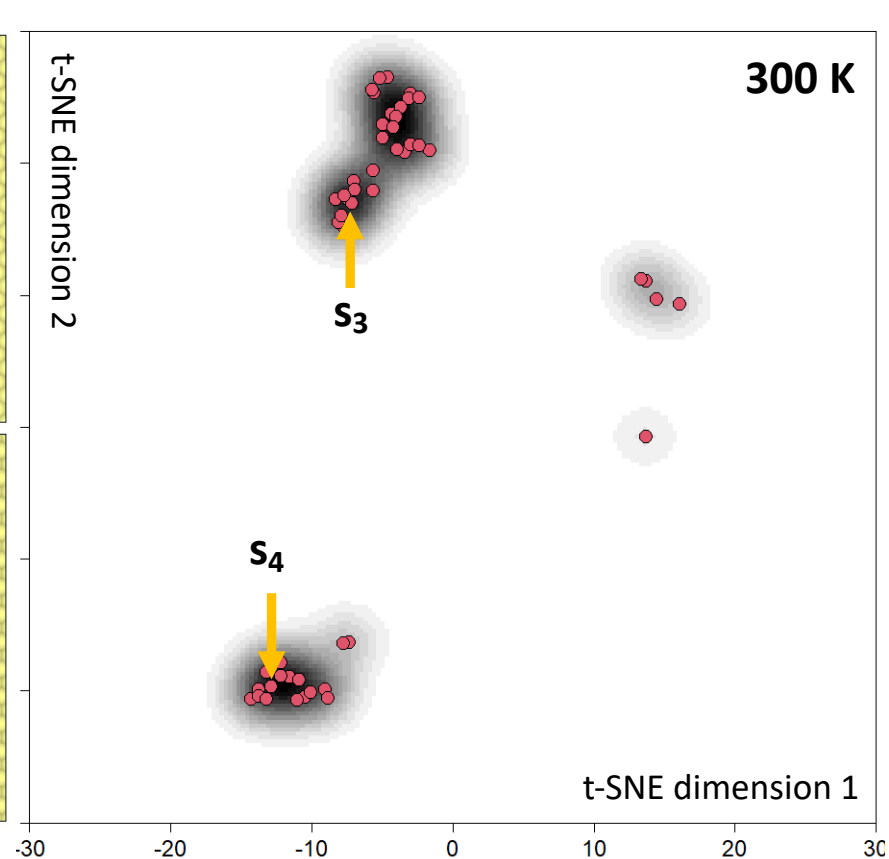
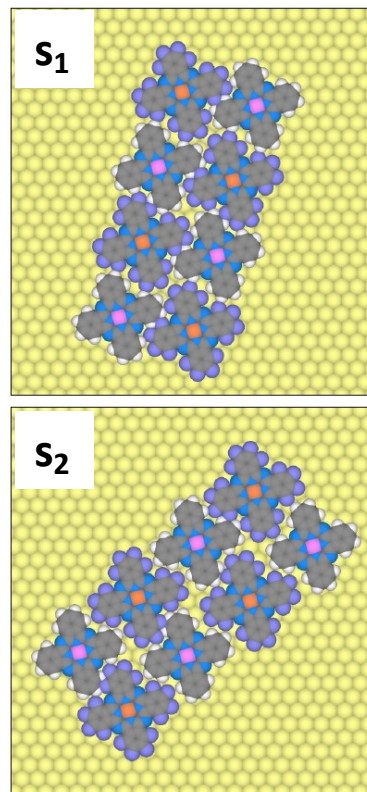
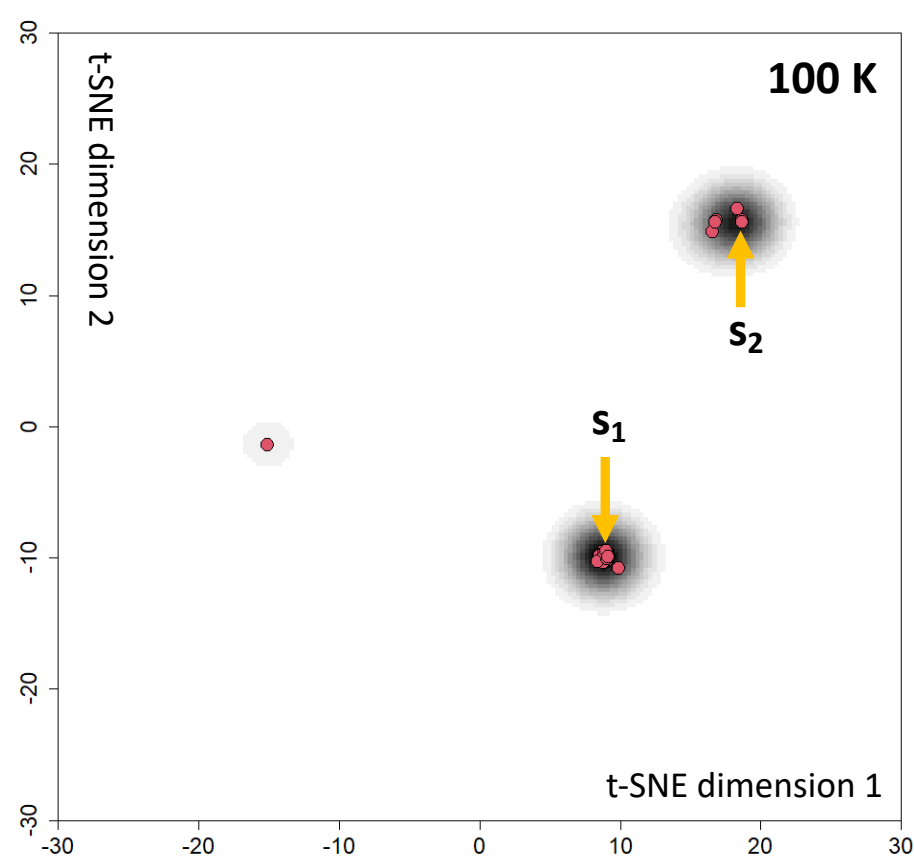


- MCMC gives us a sample of configurations that appear at thermodynamic equilibrium.
- We apply t-distributed stochastic neighborhood embedding (tSNE) to visualize the MCMC sample in 2D. The configurations (red points) are arranged according to their structures.
- The configurations sampled by MCMC are mostly located in two clusters. Within these clusters, the configurations are mostly identical.

## Interpretation.

At 300 K, are the predominantly two types of configurations on the surface.



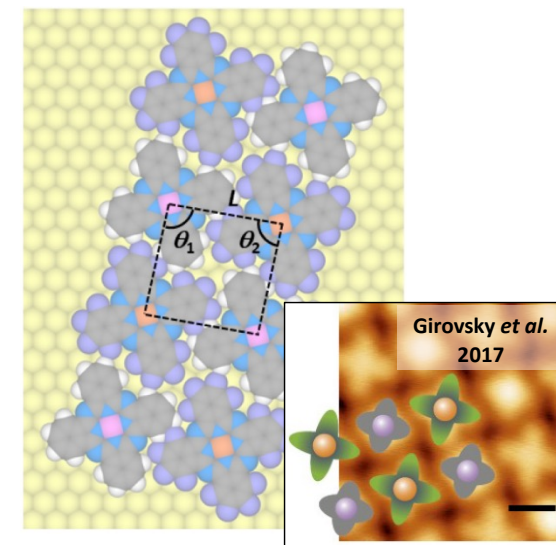


To compare with experiment, we consider simulations performed at low temperature (100 K).

Two types of configurations are also seen at 100 K, although different from the ones at 300 K.

Configuration  $s_1$  achieves quantitative agreement with cryogenic STM images. Another phase also reported experimentally (perhaps  $s_2$ ?) .

	Calculation (mean $\pm$ st err)	Experiment (mean $\pm$ st err)
$L$	$14.14 \pm 0.07$	$14.05 \pm 0.08$
$\theta_1$	$95.27 \pm 1.86$	$95.84 \pm 1.26$
$\theta_2$	$84.73 \pm 1.87$	$84.84 \pm 0.79$





# Predicting magnetic properties

The spin directions of the Mn / Fe ions are very difficult to predict from first-principles. We therefore use a classical Ising model:

$$H = \sum_{i \neq j} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j$$

$\mathbf{S}_i$  = spin vector for spin  $i$

$J_{ij}$  = Exchange constant

For the case of RKKY interactions between spins\*:

$$J_{ij} = Qg(2k_F r_{ij}) \quad g(x) = \left(\frac{2}{x}\right)^4 (\sin x - x \cos x)$$

$Q$  = System-dependent constant

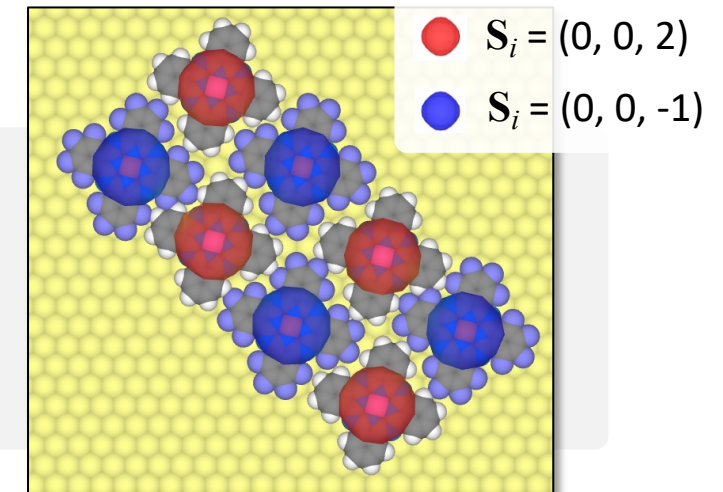
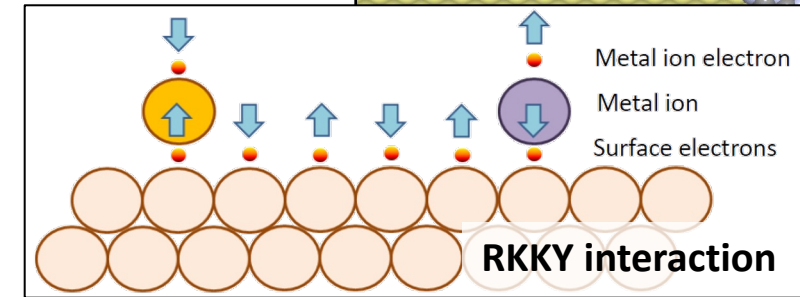
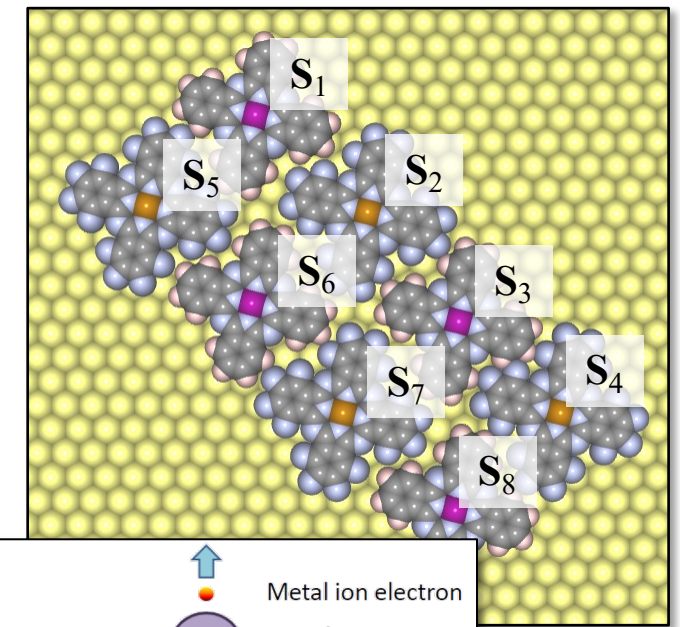
$r_{ij}$  = Distance between spins  $i$  and  $j$

$k_F$  = Fermi vector for surface  
(about  $0.18 \text{ \AA}^{-1}$  for Au(111)\*)

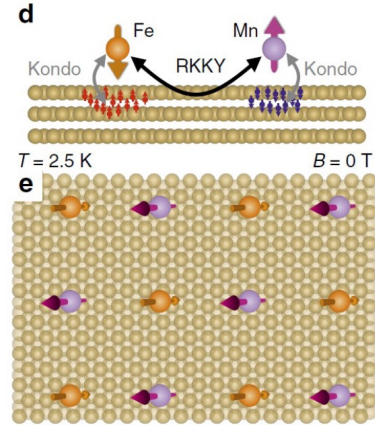
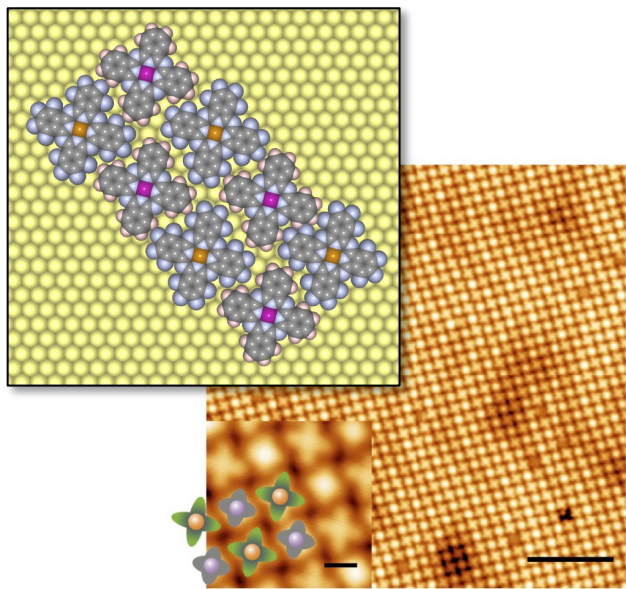
## Result:

This model correctly predicts an antiferromagnetic ground state for the molecular assembly!

(a direct calculation of the spin orientations using density functional theory would be preferred for studying the magnetic ordering in more detail)



# Final comments for part 1



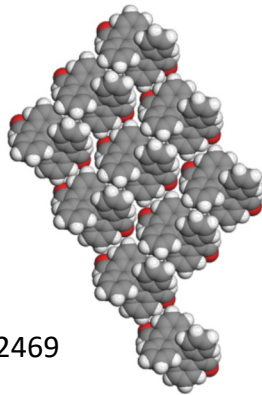
- We succeeded to create a method that can predict how molecules assemble on metal surfaces.
- It combines a machine-learned interaction potential with genetic algorithms and Monte Carlo sampling.

**References:** *Nat. Commun.* **8**, 2017, 14463;

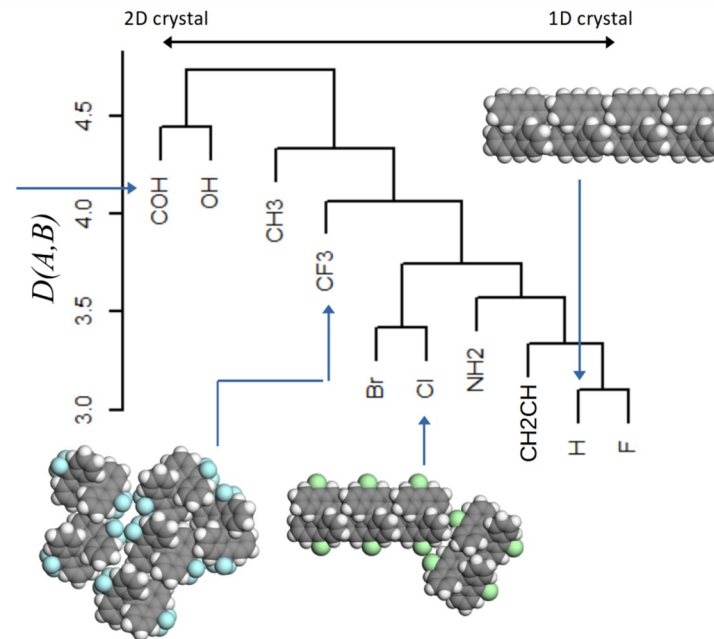
*Adv. Phys. Res.* **1**, 2022, 2200019 <- this one is better!

## We have done some other things with this method:

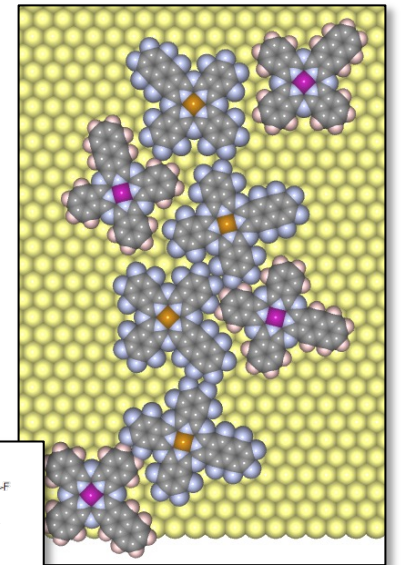
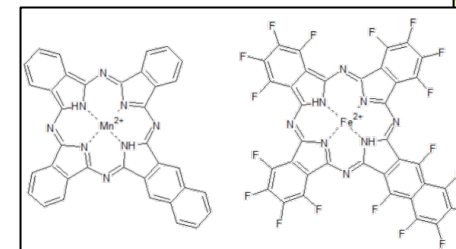
- Discovery of connection between single-molecule properties and molecular assembly shape.



*Nat. Commun.* **9**, 2018, 2469



- Prediction of disordered molecular assemblies using asymmetric molecules



*Adv. Phys. Res.* **1**, 2022, 2200019

## Lecture topics

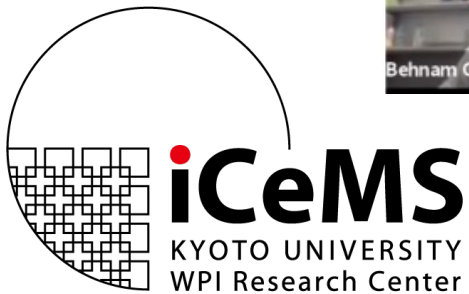
Simulation of on-surface molecular self-assembly

Machine learning for organic photovoltaic materials



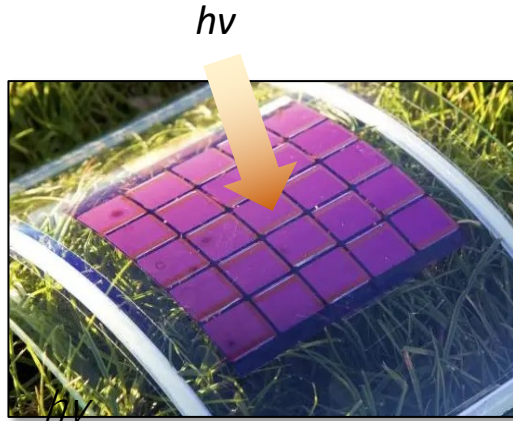
*April 2021*

# iCeMS-MacDiarmid Institute Online Workshop



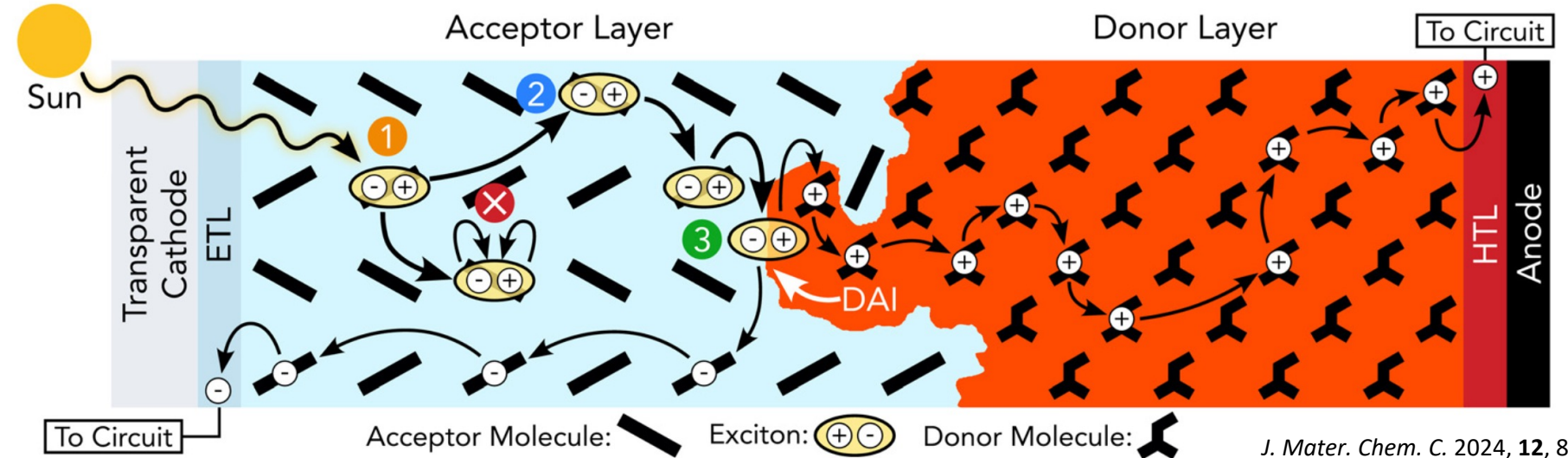


# Organic solar cells and exciton diffusion (Hodgkiss group)



Organic solar cell

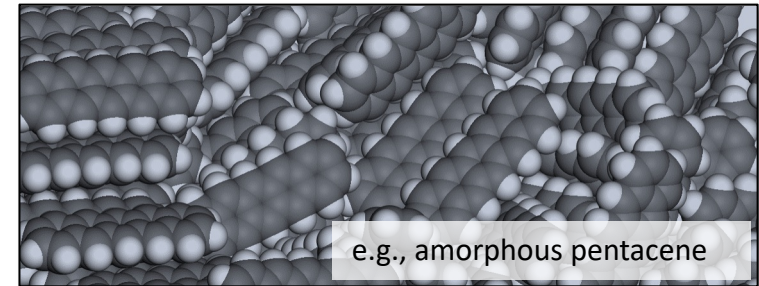
([www.solarreviews.com/blog/organic-solar-cells/](http://www.solarreviews.com/blog/organic-solar-cells/))



*J. Mater. Chem. C.* 2024, **12**, 8747

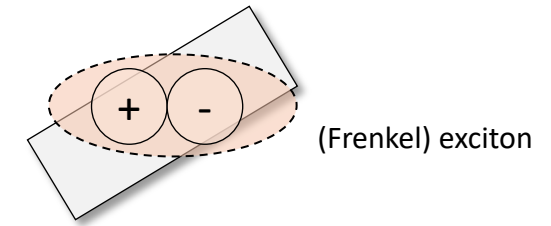
Acceptor layer is an organic semiconductor  
(crystal or amorphous solid of organic molecules)

Donor layer is another semiconductor (can be organic or inorganic)



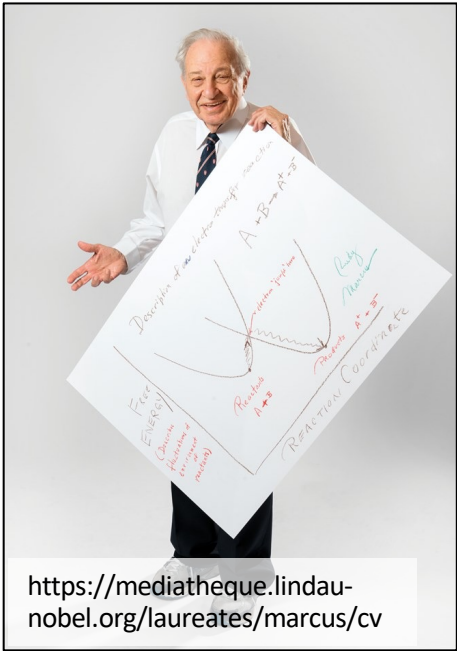
e.g., amorphous pentacene

Light absorption by the acceptor layer results in **Frenkel excitons**  
– tightly-bound electron-hole pairs localized to single molecules.



(Frenkel) exciton

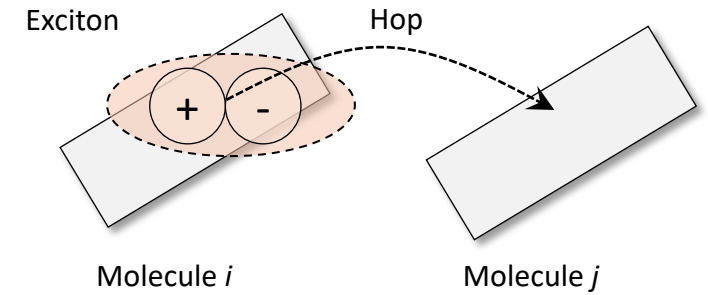
Excitons hop between molecules to the donor-acceptor interface (DAI), where they dissociate into electron and hole pairs and generate electricity. However, if the exciton moves too slowly, it will die by electron-hole recombination.



## Exciton hopping

About 60 years ago, Rudolf Marcus showed that the rate of exciton hopping between molecules is approximately

$$k_{ij} = \left( \frac{\pi}{\lambda k_B T} \right)^{\frac{1}{2}} \frac{v_{ij}^2}{\hbar} \exp \left( - \frac{\lambda}{4k_B T} \right)$$

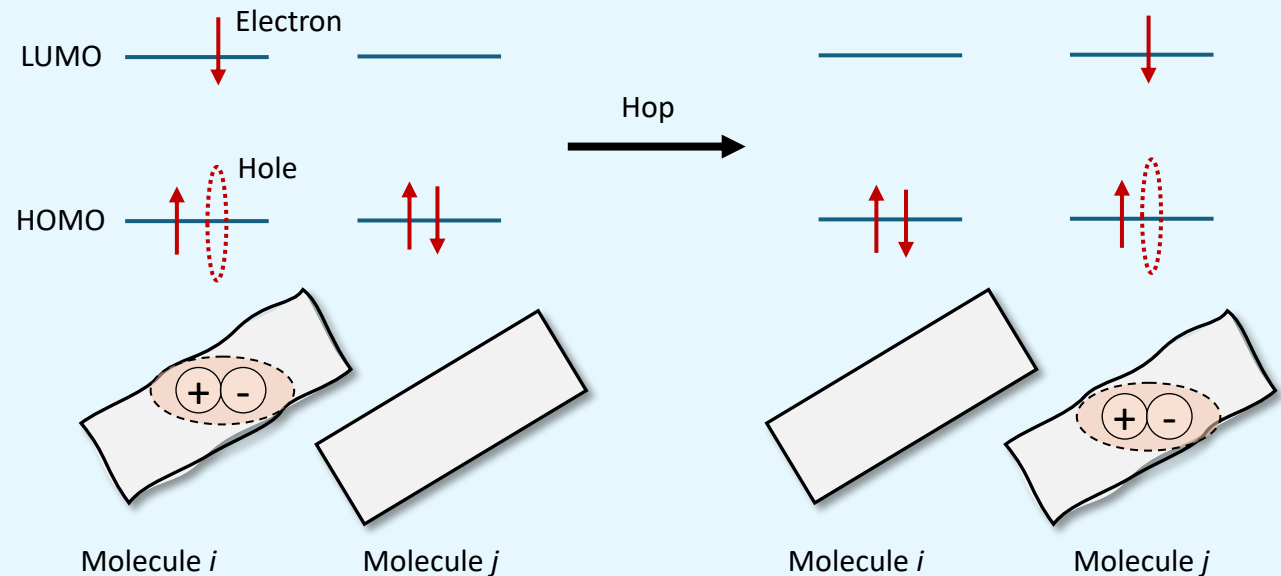


$1/k_{ij}$  is the (average) time taken for the exciton to hop from molecule  $i$  to molecule  $j$ .

**Important parameter 1.**  
**Reorganization energy ( $\lambda$ )**

The exciton distorts the shape (conformation) of the molecule.

The reorganization energy measures the energy required to change the shapes of the molecules as the exciton shifts.

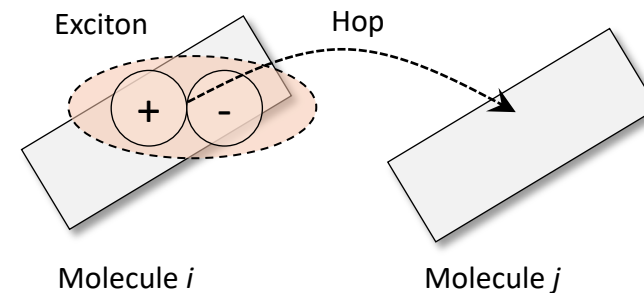


## Important parameter 2. Exciton coupling

Marcus equation:

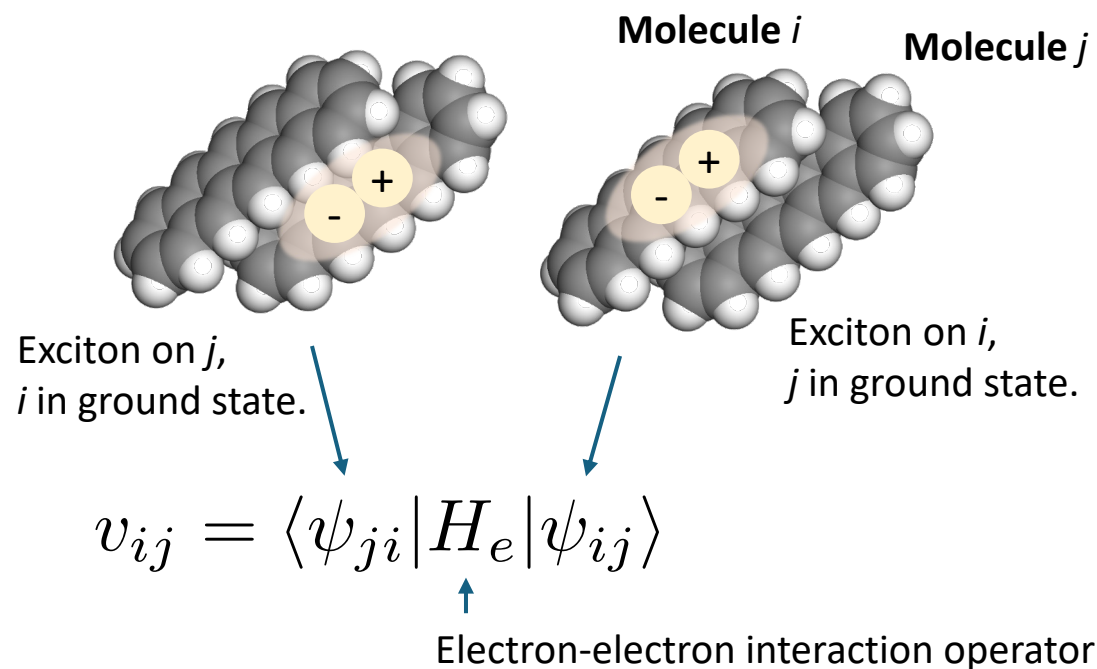
$$k_{ij} = \left( \frac{\pi}{\lambda k_B T} \right)^{\frac{1}{2}} \frac{v_{ij}^2}{\hbar} \exp \left( -\frac{\lambda}{4k_B T} \right)$$

Exciton coupling Reorganization energy



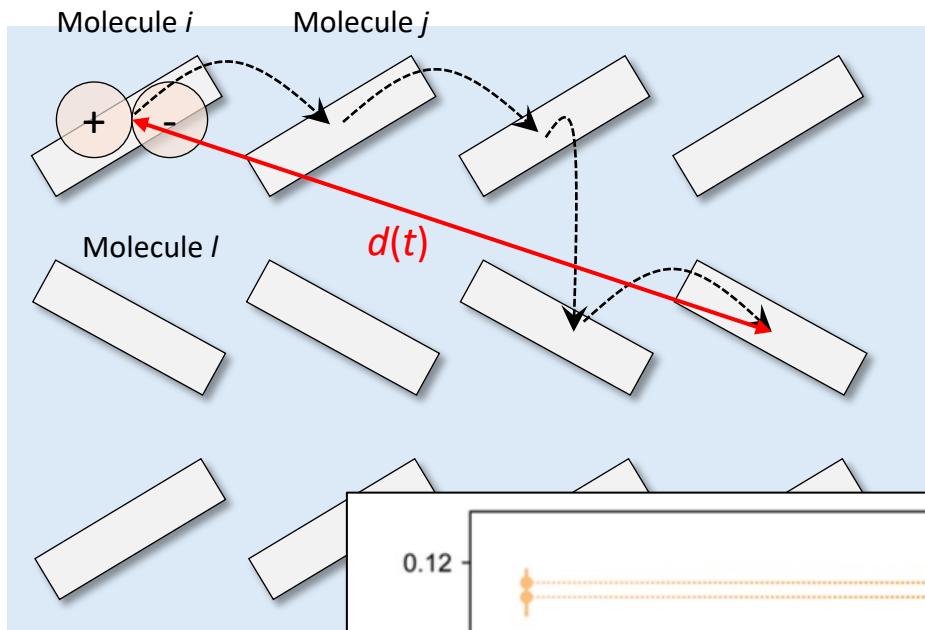
The exciton coupling  $v_{ij}$  is a quantum mechanical parameter. Its interpretation is somewhat abstract:

- Start with  $|\psi_{ij}\rangle$ . This is the wavefunction for the two molecules together, with the exciton on molecule  $i$  (initial state, before the hop).
- Electron-electron interactions between molecules perturbs  $|\psi_{ij}\rangle$ . The perturbed wavefunction is represented as  $H_e|\psi_{ij}\rangle$ .
- Now consider  $\langle\psi_{ji}|$ . This is the final wavefunction, with the exciton on molecule  $j$  (final state, after the hop).
- Finally,  $\langle\psi_{ji}|H_e|\psi_{ij}\rangle$  tells us the overlap between the perturbed initial wavefunction and the final wavefunction.



**Both of these parameters can be calculated from first-principles (time-dependent DFT / TDDFT). However, these calculations are very time consuming!**

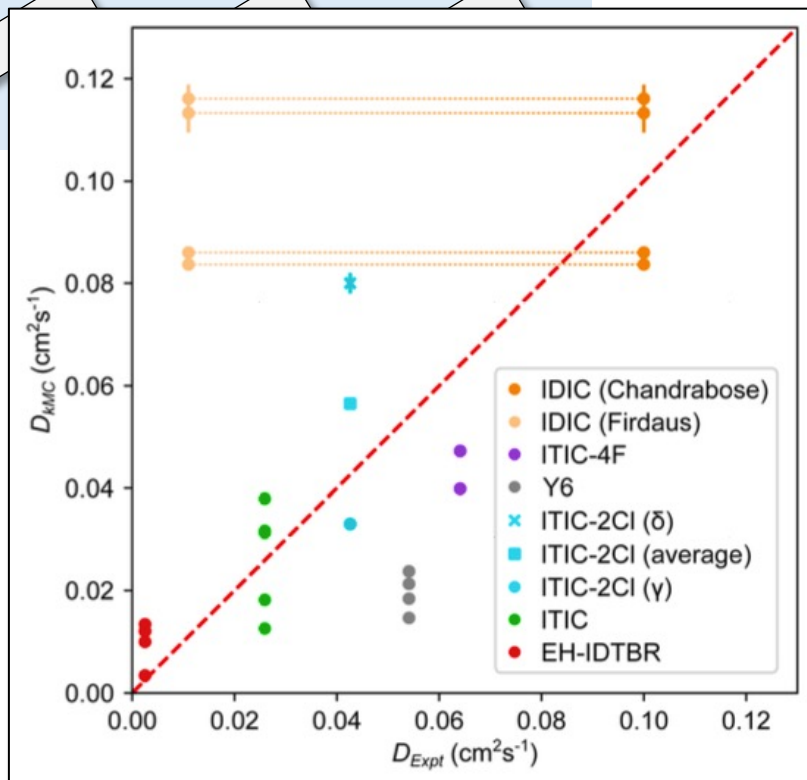
# Kinetic Monte Carlo (kMC) simulation of exciton diffusion



Once the hopping rates are calculated, the hopping process can be simulated using the kinetic Monte Carlo (kMC) method.

Roughly, the exciton hops to its neighbor  $j$  with probability proportional to  $k_{ij}$ . This results in a random walk-type motion through the crystal\*.

Calculations of  $D$  compared to spectroscopic measurements (crystalline materials)



Exciton diffusion coefficient can be estimated as:

Mean-square displacement

$$D = \lim_{t \rightarrow \infty} \frac{1}{6} \frac{\overline{d(t)^2}}{t}$$

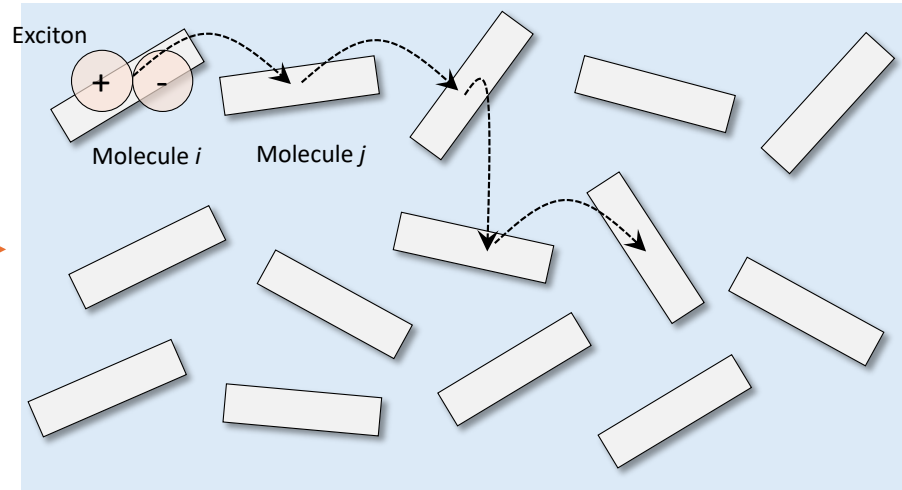
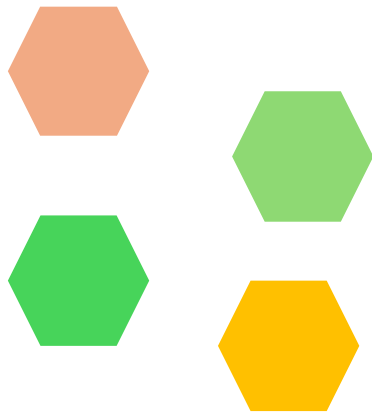
kMC simulations yield reasonable predictions of diffusion coefficients.

Could we discover new organic materials with large diffusion coefficients by high-throughput kMC simulations?

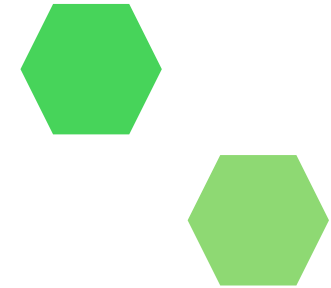


# Kinetic Monte Carlo (kMC) for materials discovery?

Candidate  
organic semiconductors



Candidates with high exciton  
diffusion coefficients



Looks feasible – kMC can be done quite quickly on a laptop (minutes per candidate).

## Problem: kMC has a big computational overhead!

For each candidate, couplings and reorganization energies need to be pre-computed from TDDFT before kMC can start.

Can we use machine learning to predict these parameters quickly?

$$k_{ij} = \left( \frac{\pi}{\lambda k_B T} \right)^{\frac{1}{2}} \overset{\text{Exciton coupling}}{v_{ij}^2} \exp \left( - \frac{\overset{\text{Reorganization energy}}{\lambda}}{4k_B T} \right)$$

**2021 - 2022**

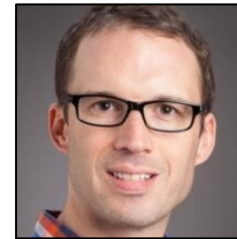
*First attempt at machine learning for exciton couplings*

**Dr. Chayanit  
Wechwithayakhlung**



**Data science**

**Prof. Justin  
Hodgkiss**



**Dr. Paul  
Hume**

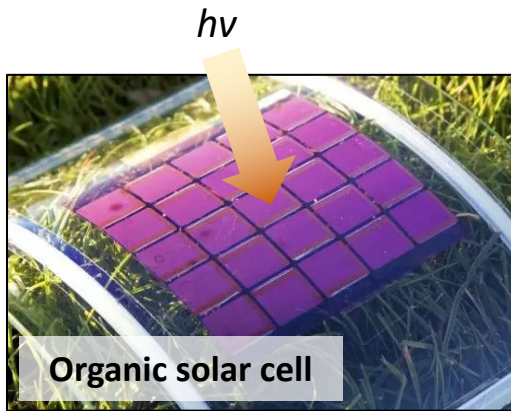
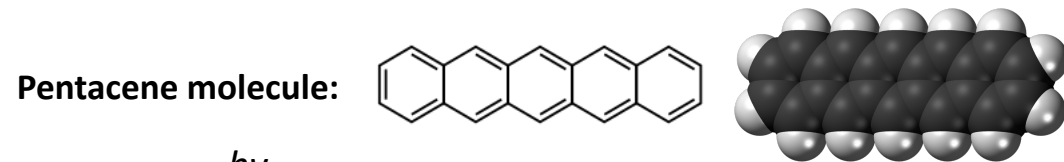
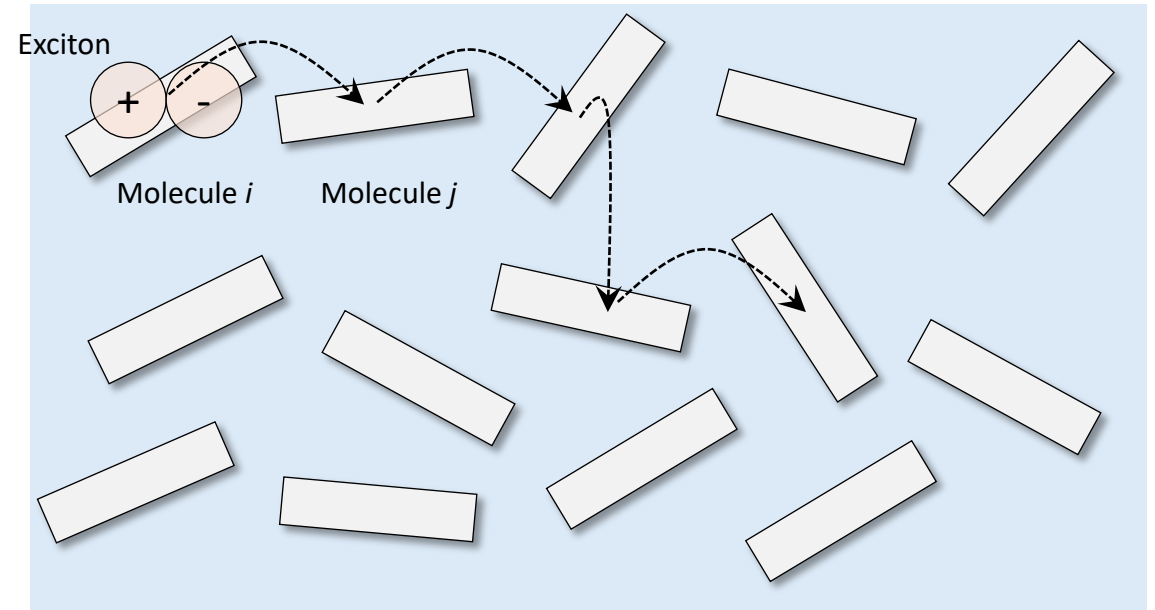
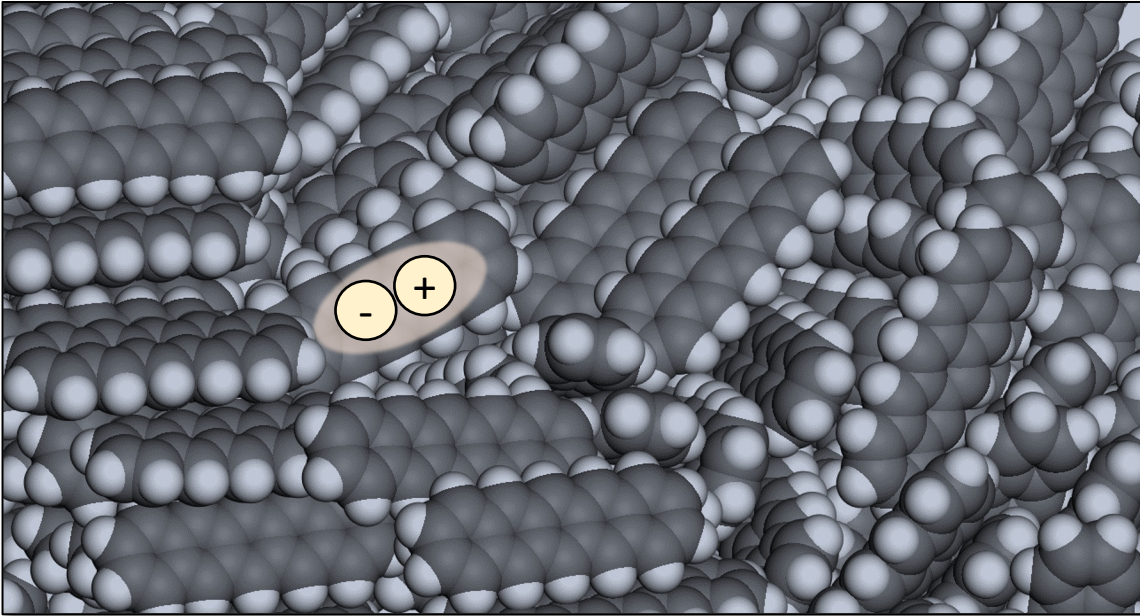


**Dr. Geoff  
Weal**



**Ultrafast spectroscopy  
Excited state quantum chemistry**

# First target - exciton diffusion in amorphous pentacene

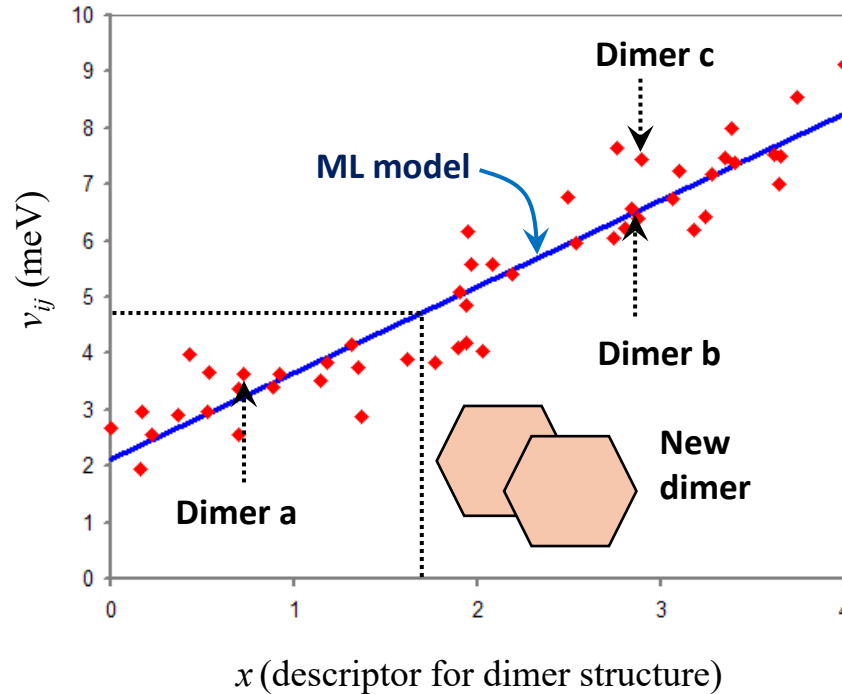
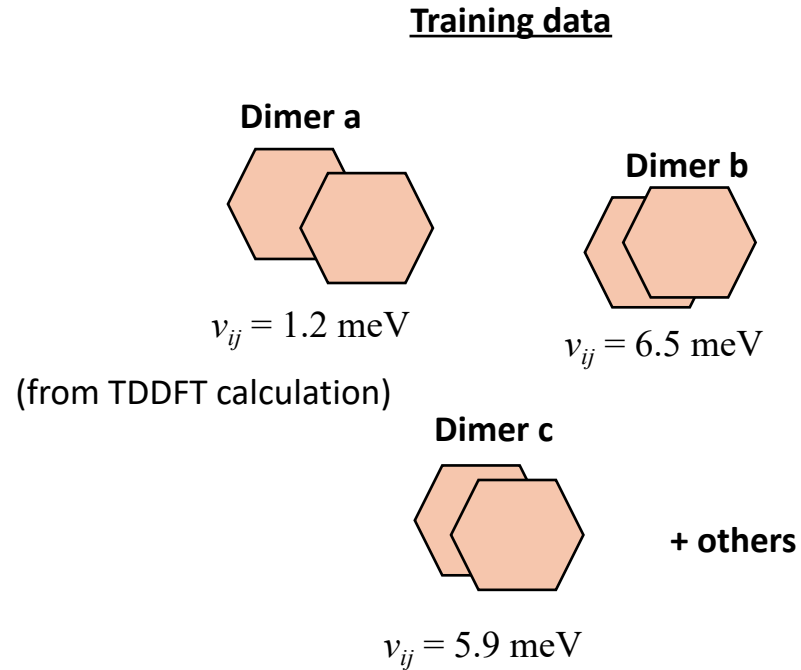


$$k_{ij} = \left( \frac{\pi}{\lambda k_B T} \right)^{\frac{1}{2}} \frac{\overset{\text{Exciton coupling}}{v_{ij}^2}}{\hbar} \exp \left( -\frac{\lambda}{4k_B T} \right)$$

Amorphous pentacene is an organic semiconductor which can be used in organic solar cells.

Simple and relevant molecule, good starting point.

# Machine learning concept

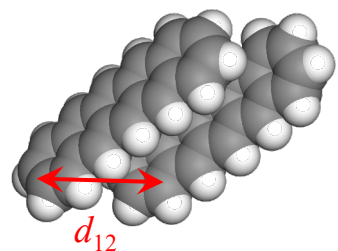


Again, we aim to fit a model which predicts coupling from dimer structure ( $x$ ).

With the fitted model, couplings for new dimers can be quickly predicted.



# ML model for exciton couplings



**Dimer** (coupling =  $v_{ij}$ )

Pair of molecules within a 5 Å interatomic cut-off

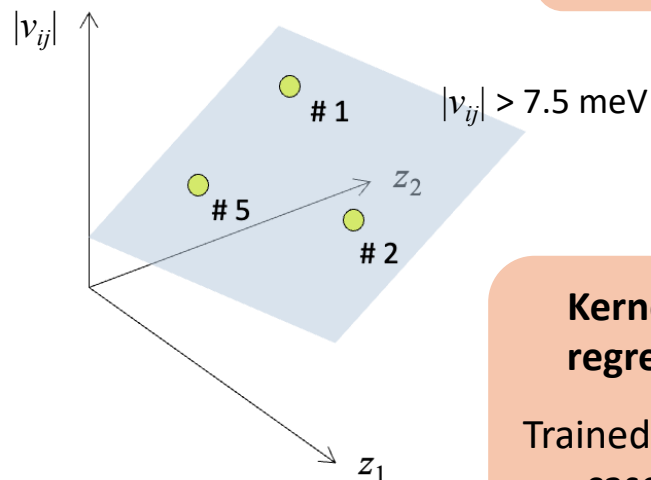
**Coulomb-type descriptors**

$$x_1 = 1/d_{11}, x_2 = 1/d_{12}, \dots$$

Sorting  
PCA

**Support vector machine**

Classify coupling as strong or weak



**Kernel ridge regression 1**

Trained on strong cases only

$v_{ij}$  prediction

**Kernel ridge regression 2**

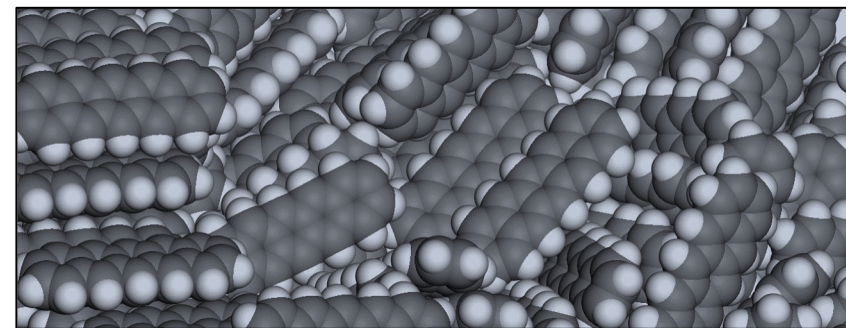
Trained on weak cases only

$v_{ij}$  prediction

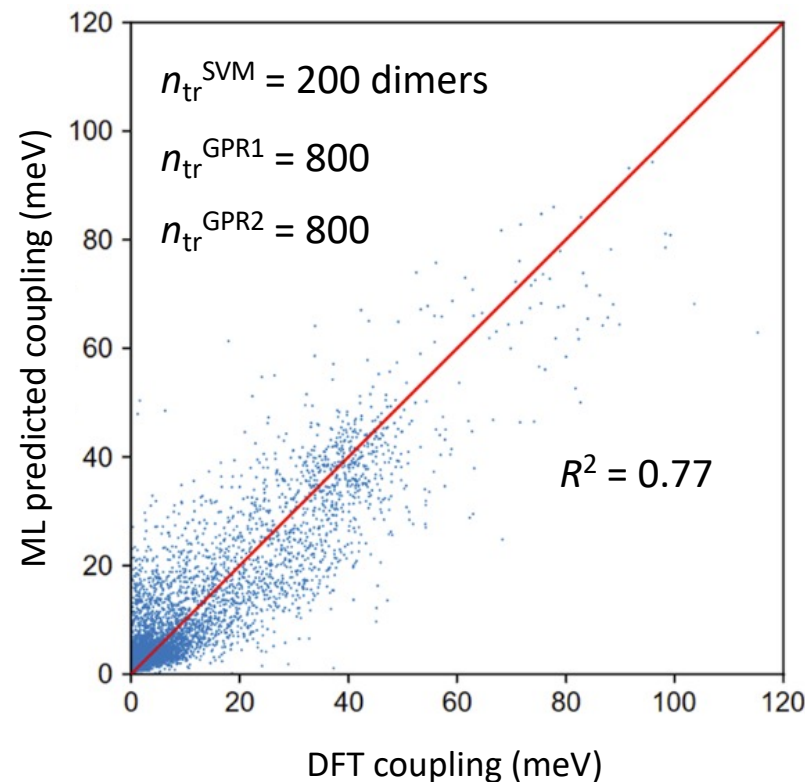
DFT settings: CAM-B3LYP xc functional with 6-3111+G(2d,p) basis set as implemented in Gaussian.

**System under consideration:**

Amorphous pentacene (generated by Dr. Yu Kaneko (DAICEL))



**Performance on 4127 dimers**



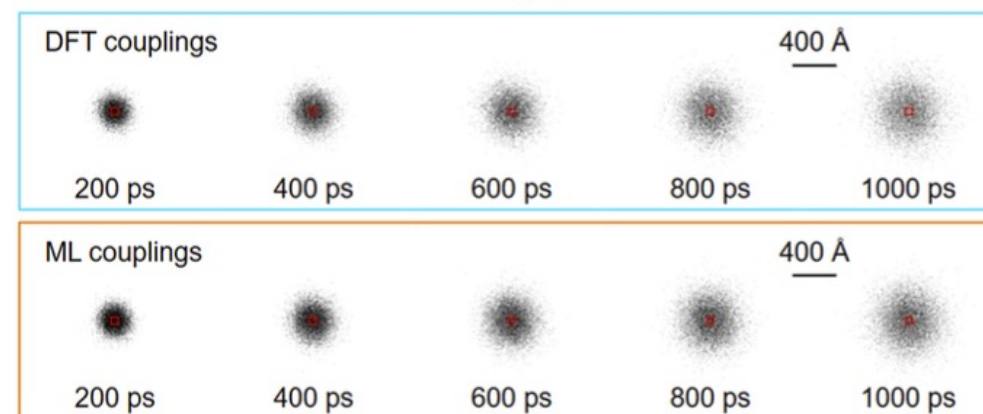
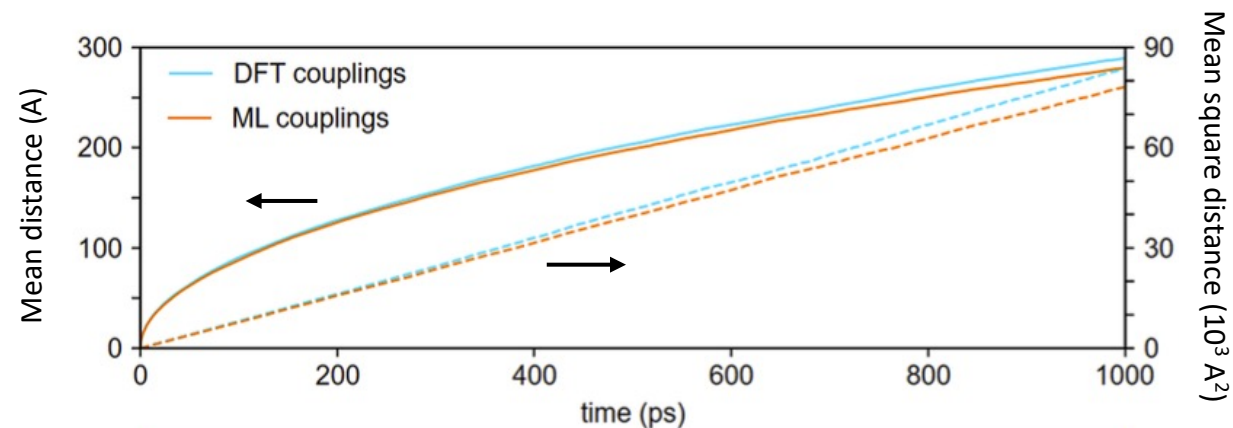
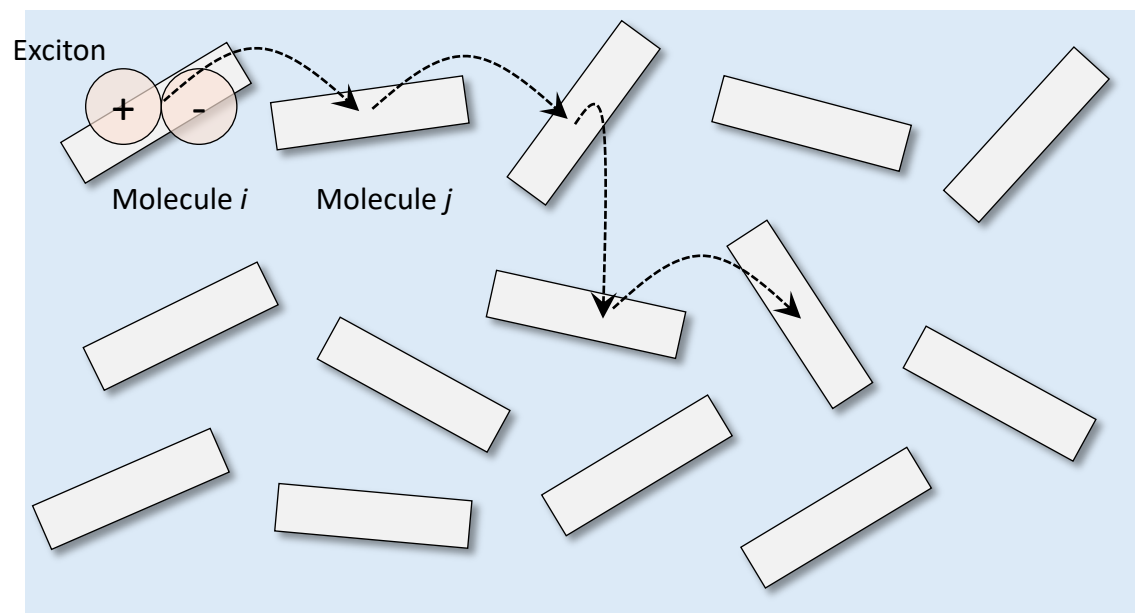
Average prediction time around 7 ms on an office workstation\*

(cf. 7.4 hours for TDDFT on supercomputer\*\*)

\* Single 3.50 GHz Intel Xeon E5-1620 core

\*\* 64 cores (128 threads), 2.0 – 3.35 GHz AMC Epyc 7702 processors.

# Kinetic Monte Carlo simulation in amorphous pentacene



(Statistics from  $10^4$  trials)

	<i>Ab initio</i> couplings	Model-predicted couplings
Diffusion coefficient ( $\times 10^{-3} \text{ cm}^2 \text{ s}^{-1}$ )	$1.630 \pm 0.011$	$1.547 \pm 0.005$
Diffusion tensor eigenvalues ( $\times 10^{-3} \text{ cm}^2 \text{ s}^{-1}$ )		
Major	$1.815 \pm 0.014$	$1.686 \pm 0.017$
Middle	$1.551 \pm 0.016$	$1.492 \pm 0.007$
Minor	$1.525 \pm 0.012$	$1.462 \pm 0.014$

(Experimental diffusion coefficient for multicrystalline pentacene:  $0.5 \times 10^{-3} \text{ cm}^2 \text{ s}^{-1}$ )

**Good start, but a serious problem remains:**

Coupling model restricted to only one type of molecule (pentacene). Cannot use this for virtual screening!

**2023 - 2024**

# *General exciton coupling model*

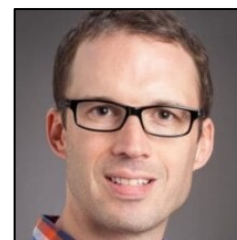
**Dr. Geoff  
Weal**



**Data science**



**Prof. Justin  
Hodgkiss**

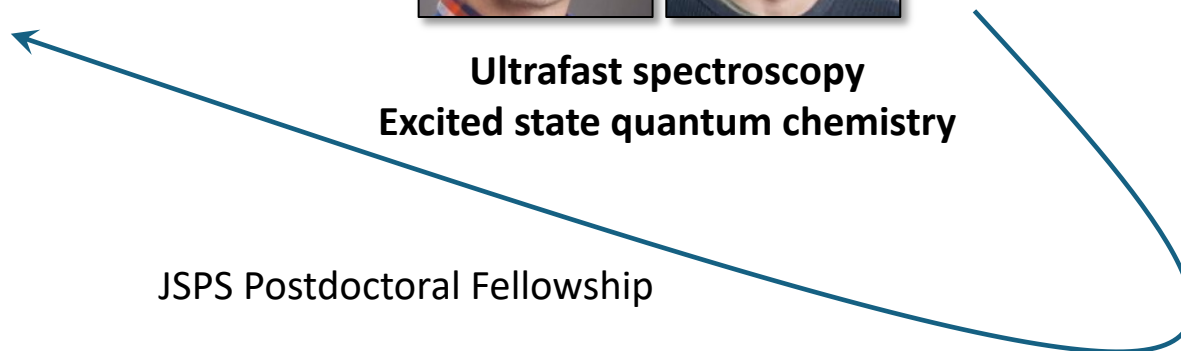


**Dr. Paul  
Hume**

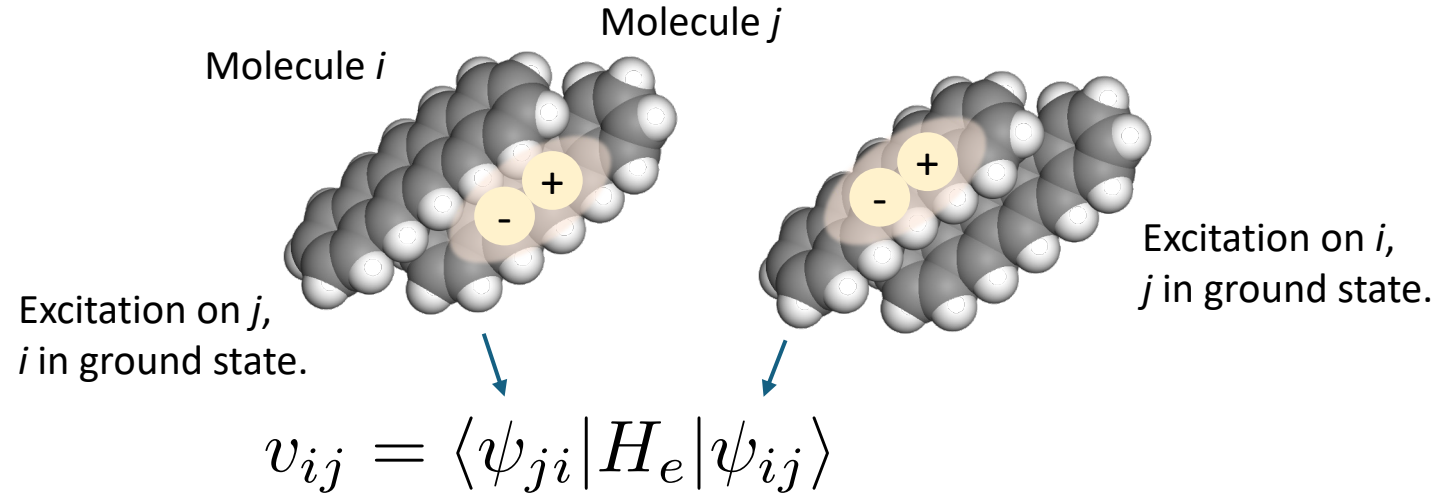
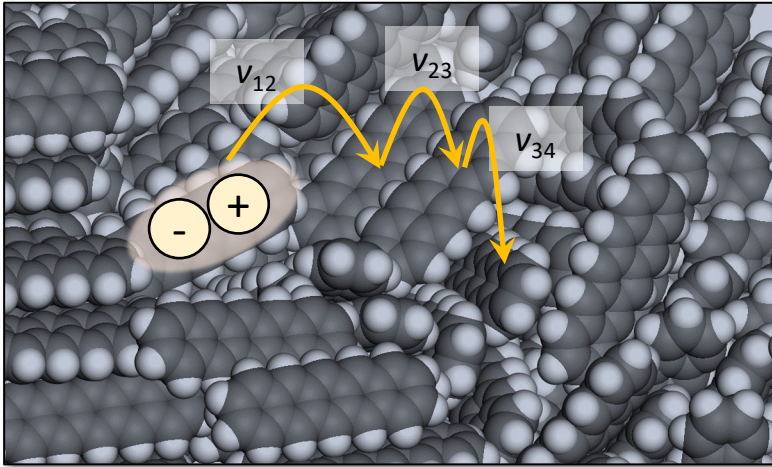


**Ultrafast spectroscopy  
Excited state quantum chemistry**

JSPS Postdoctoral Fellowship



# Towards a new coupling model that works for all types of molecules?



Longuet-Higgins showed that  $v_{ij}$  can be approximated as a sum over simple, Coulomb-like terms\*:

$Q_\alpha^i$  is the **atomic transition charge (ATC)** for atom  $\alpha$ .

It measures the change in the electron density on atom  $\alpha$  when molecule  $i$  transitions from the ground to the excited state ( $\rho_{\text{tr}}^i(\mathbf{r})$  is called the transition density)

This is known as the **ATC approximation**. It has never been tested extensively, so we put it to work...

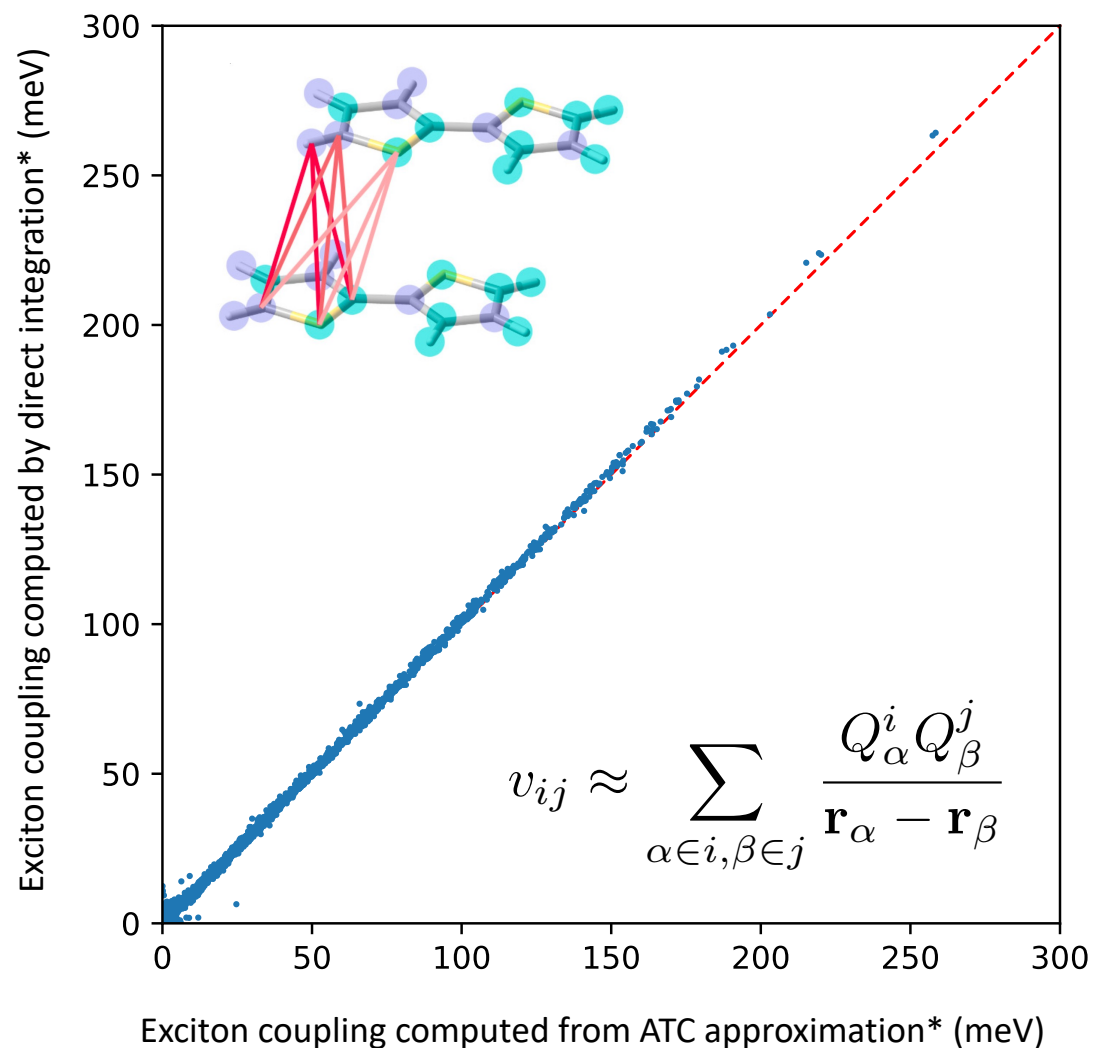
$$v_{ij} \approx \sum_{\alpha \in i, \beta \in j} \frac{Q_\alpha^i Q_\beta^j}{|\mathbf{r}_\alpha - \mathbf{r}_\beta|}$$

$$Q_\alpha^i = \int_{\mathbf{r} \text{ on atom } \alpha} \rho_{\text{tr}}^i(\mathbf{r}) d\mathbf{r}$$

\*Longuet-Higgins. *Proc. Roy. Soc.* **235**, 1956, 537



# Towards a new coupling model that works for all types of molecules?



We tested the ATC approximation for molecular dimers extracted from 1989 organic crystals in the CCDC database (Cambridge Crystal Data Center)

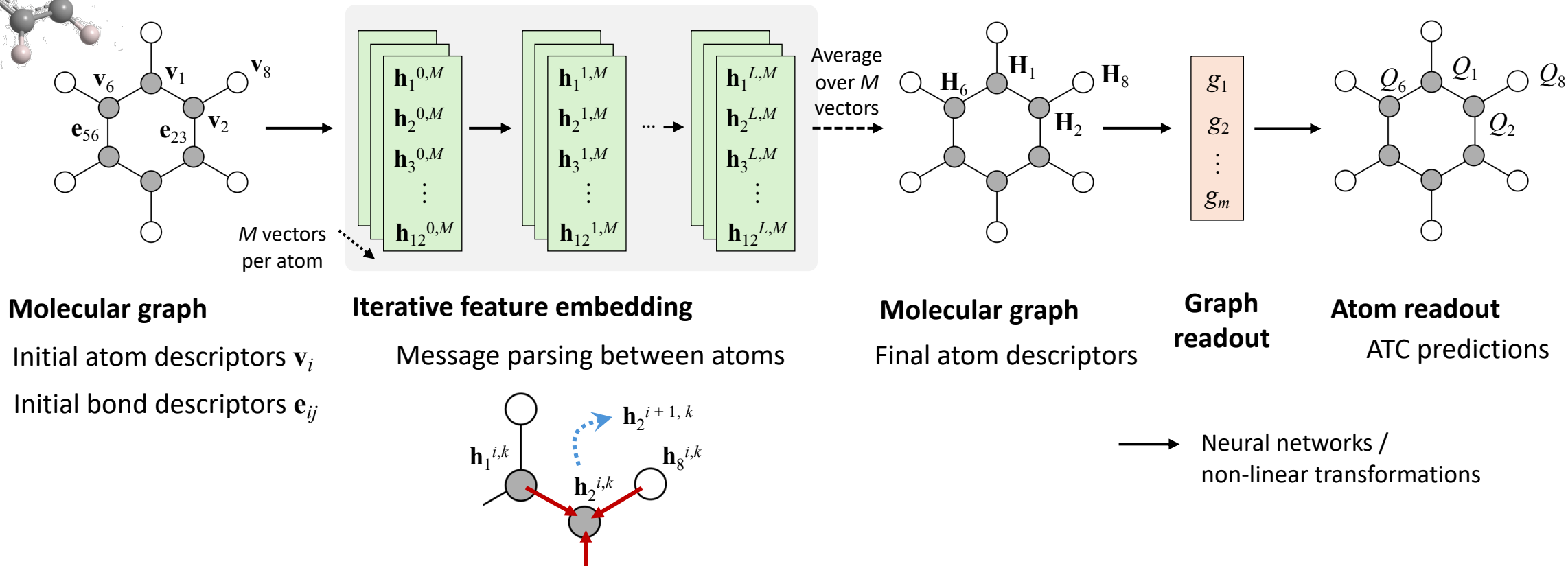
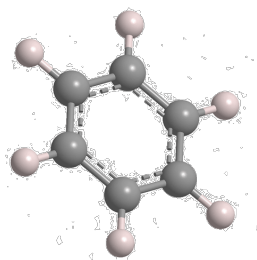
Good agreement, with only small deviations for strong coupling cases (ATC approximation neglects some close-range quantum effects)

=> Instead of creating an ML model for couplings directly, let's **make an ML model for atomic transition charges ( $Q_{\alpha}$ )**.

We could then compute  $v_{ij}$  for any type of molecule using the ATC approximation.

\* TDDFT with  $\omega$ -B97XD xcf with 6-31+G(d,p) basis set as implemented in Gaussian 16.

# Graph neural network to predict atomic transition charges (ATCs)



## Implementation

SOAP (Smooth Overlap of Atomic Positions) descriptors used for  $\mathbf{v}_i$ .

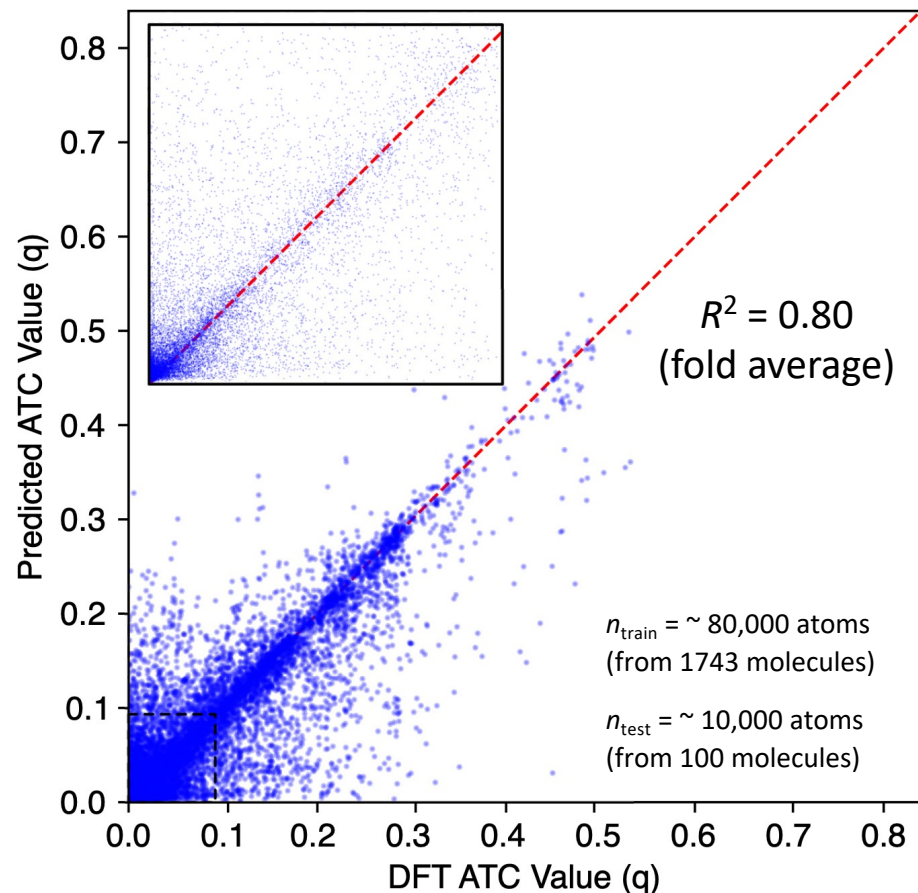
An integer-valued descriptor corresponding to hybridization state also incorporated (sp, sp<sup>2</sup>, sp<sup>3</sup>).

Integer-valued descriptor for bond type (single, double, etc) used for  $\mathbf{e}_{ij}$ .

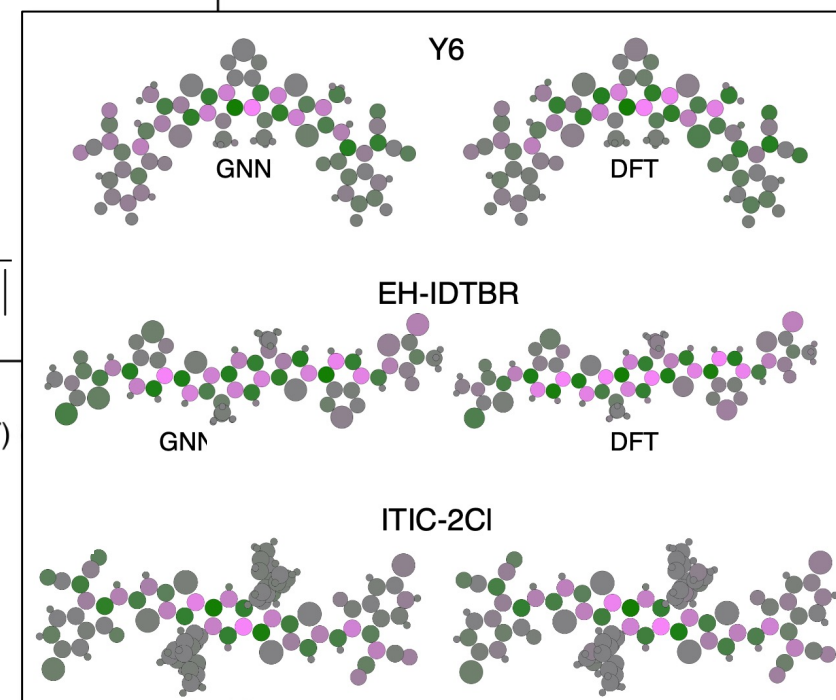
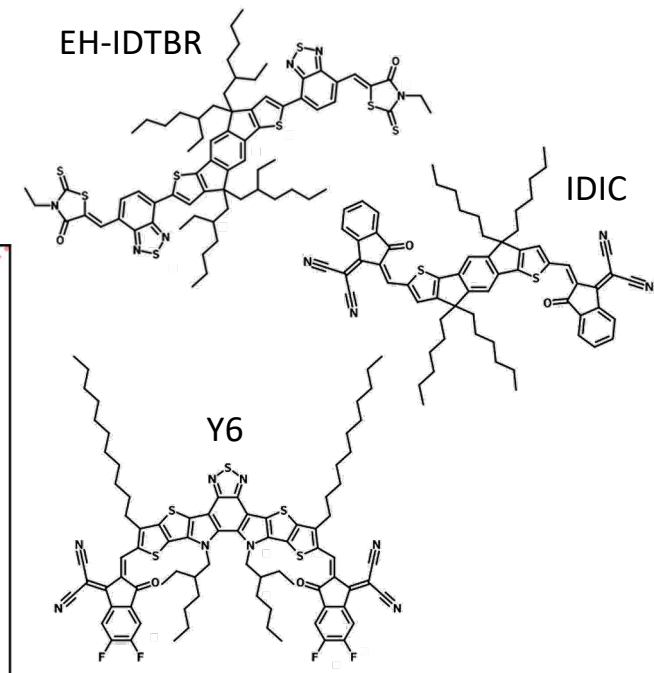
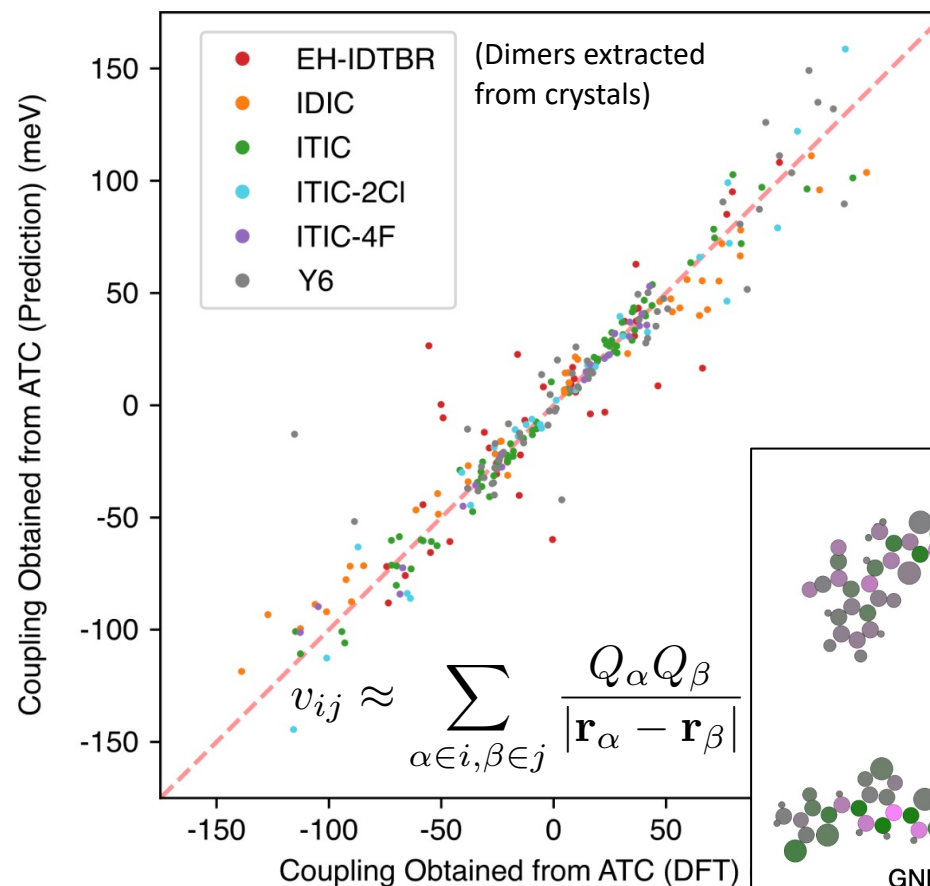
Other settings and training procedure followed Han *et al. Phys. Chem. Chem. Phys.* **24**, 2022, 26870.

# Graph neural network performance

Performance on test set (10xCV fold)



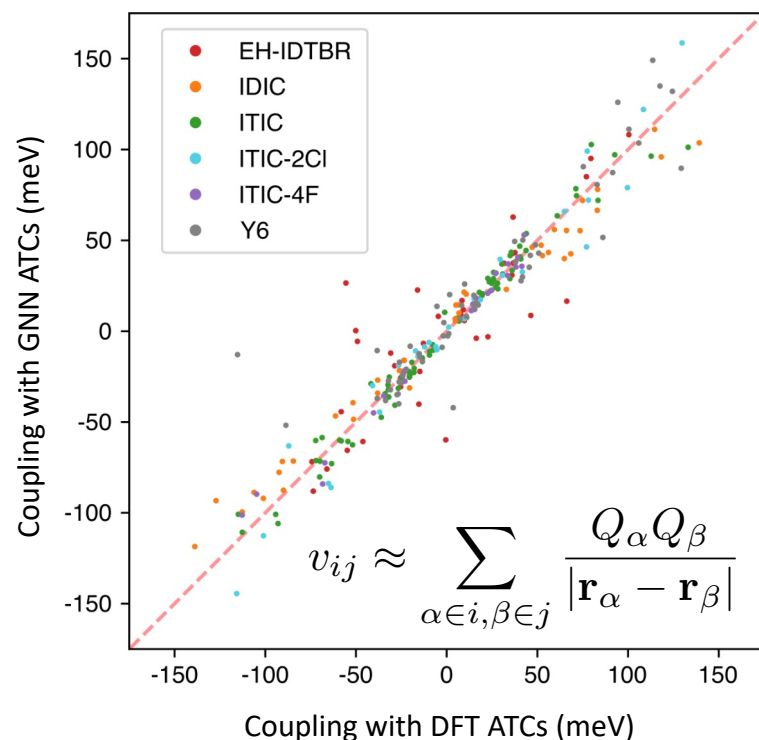
Exciton coupling predictions for important organic photovoltaic materials (FREAs)



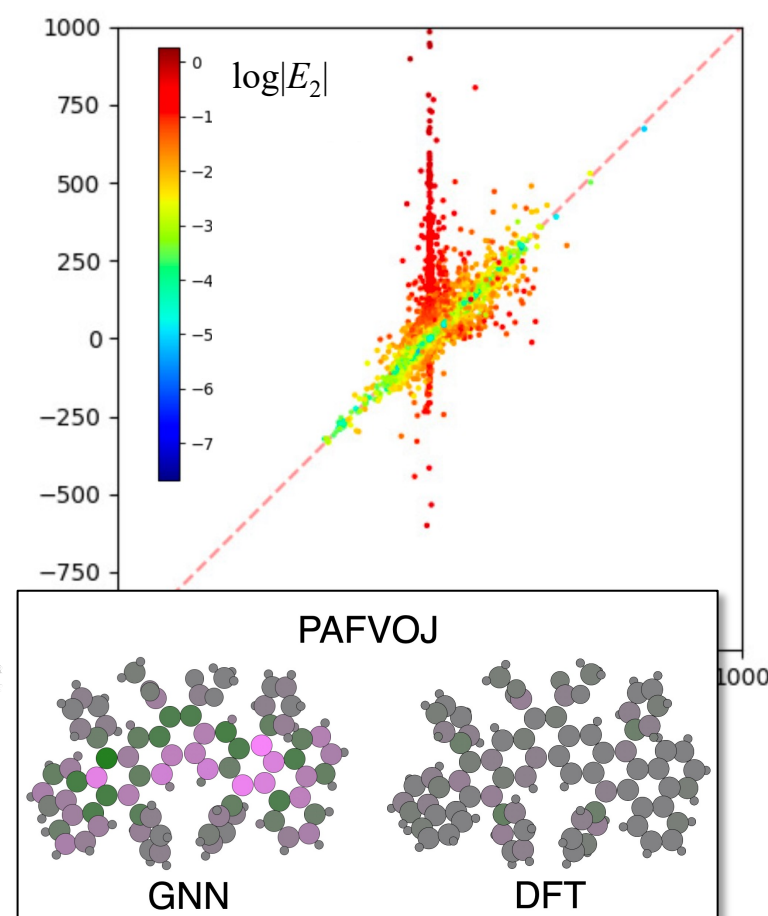
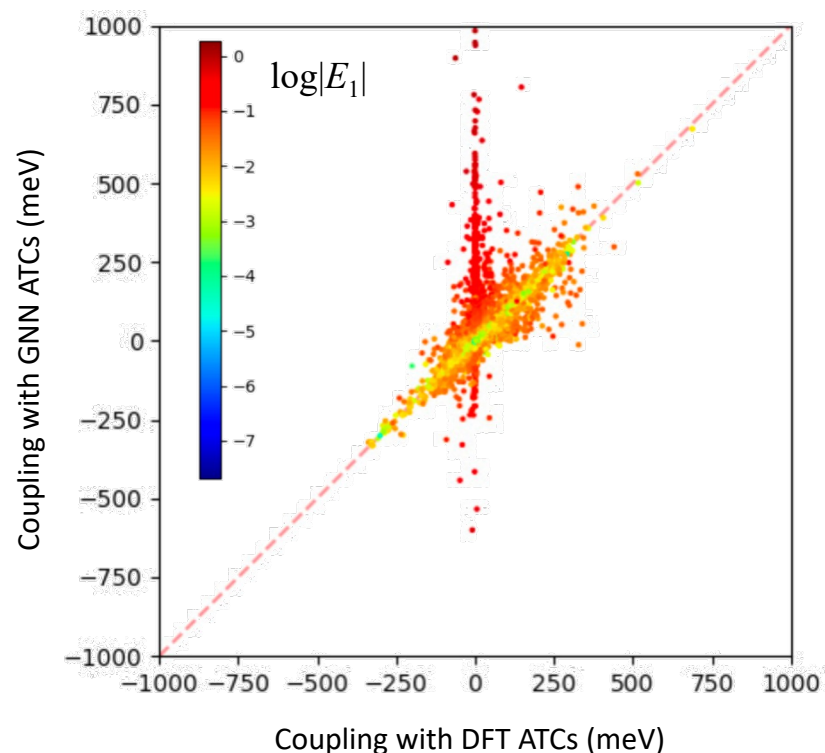
- Reasonable accuracy when predicting ATCs.
- **Generality achieved** - couplings predicted with good accuracy for a range of important organic photovoltaic materials!

# How general? Error propagation limits accuracy for weakly coupled cases

Coupling predictions for important organic photovoltaic materials (FREAs)



Couplings for dimers extracted from 1000 organic crystals



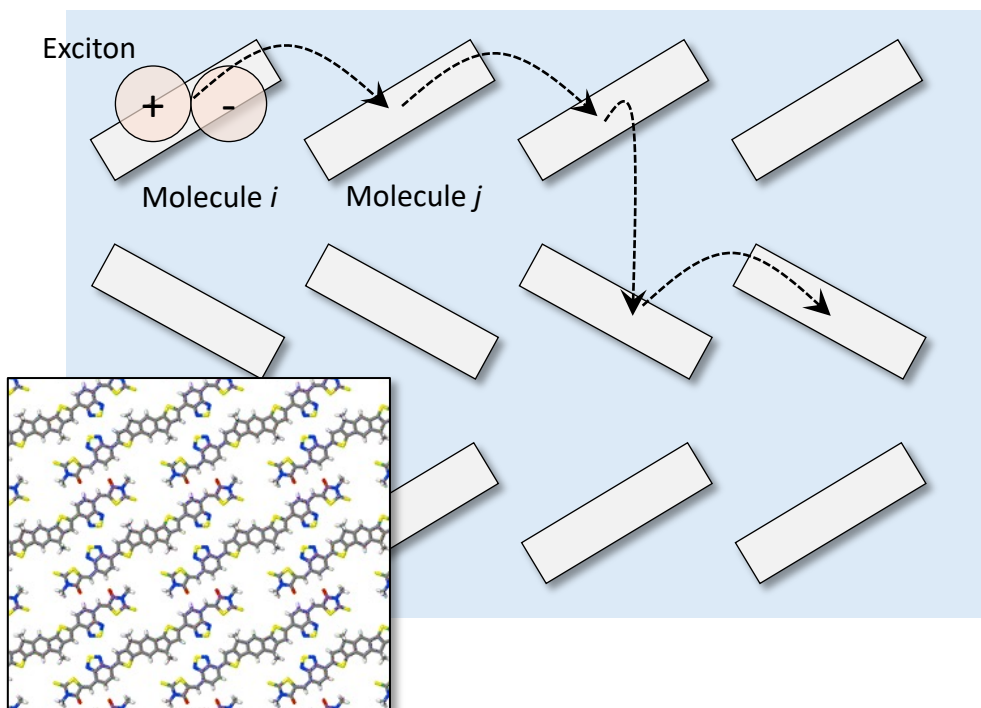
Let  $\varepsilon_{\alpha}$  be the error of the  $Q_{\alpha}$  prediction ( $\varepsilon_{\alpha} = Q_{\alpha} - Q_{\alpha}^{\text{GNN}}$ ).

Then  $v_{ij} = v_{ij}^{\text{GNN}} + E_1 + E_2$ , where  $E_1 = \sum \varepsilon_{\alpha} Q_{\beta}^{\text{GNN}} / |\mathbf{r}_{\alpha} - \mathbf{r}_{\beta}|$  and  $E_2 = \sum \varepsilon_{\alpha} \varepsilon_{\beta} / |\mathbf{r}_{\alpha} - \mathbf{r}_{\beta}|$ .

First- and second-order errors can be large for dimers where coupling is in fact weak ( $|v_{ij}|$  less than about 75 meV). For other cases, the method seems reliable.



# Kinetic Monte Carlo simulations using GNN-predicted ATCs

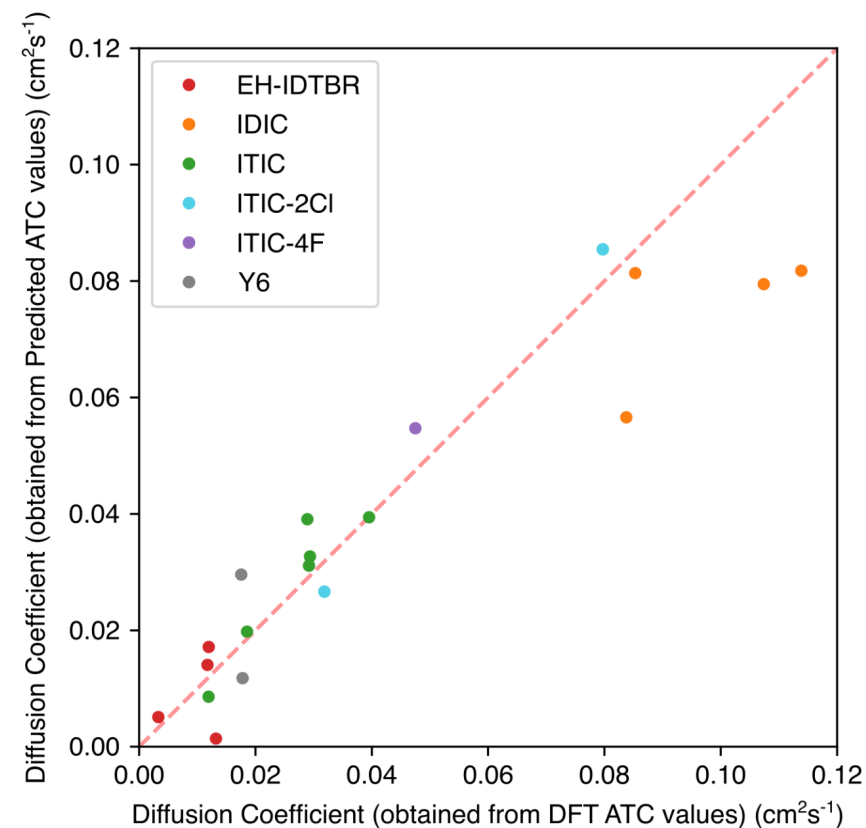


Marcus hopping rates:

$$k_{ij} = \left( \frac{\pi}{\lambda k_B T} \right)^{\frac{1}{2}} \frac{v_{ij}^2}{\hbar} \exp \left( -\frac{\lambda}{4k_B T} \right)$$

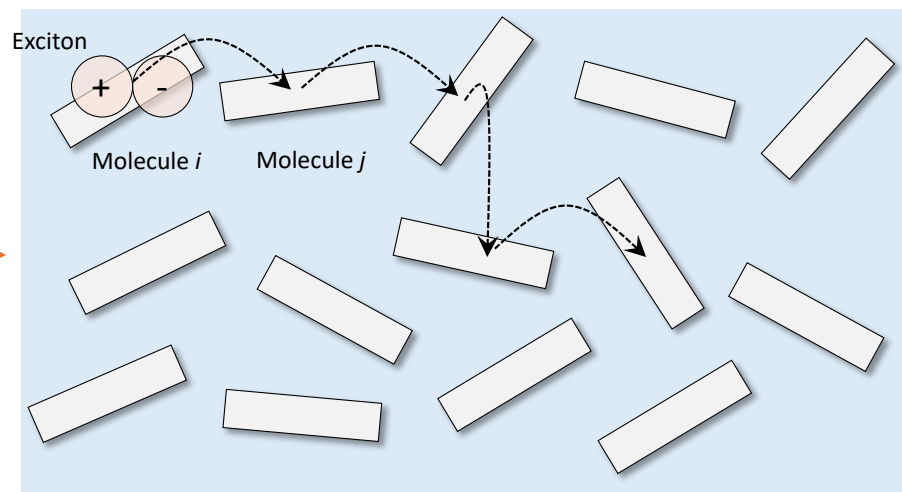
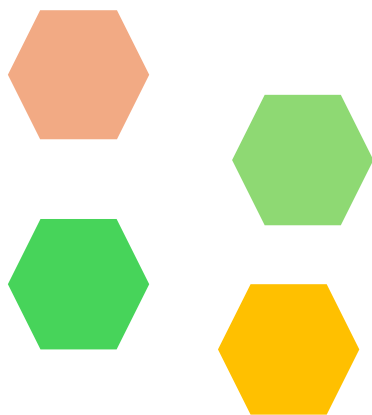
GNN-predicted exciton coupling Reorganization energy (TDDFT-calculated)

**Good predictions of diffusion coefficient obtained for multiple materials!**

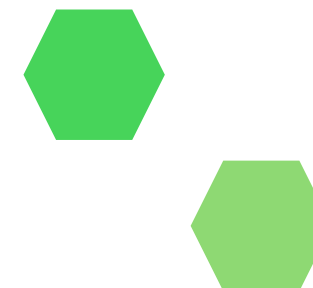


# Closer to high-throughput kinetic Monte Carlo (kMC)?

Candidate  
organic semiconductors



Candidates with high exciton  
diffusion coefficients

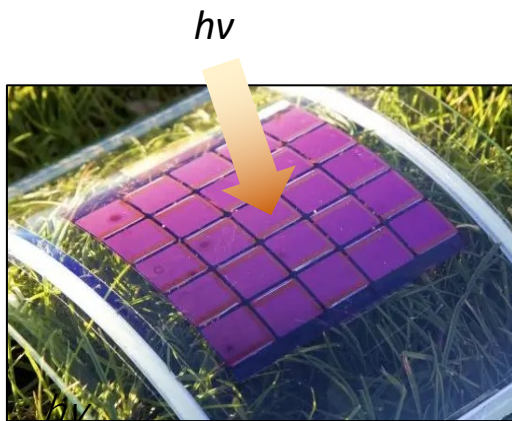


## Computational overhead remains!

We succeeded to create a general ML scheme to predict couplings.

But reorganization energies still require expensive time-dependent DFT calculations. We still have work to do!

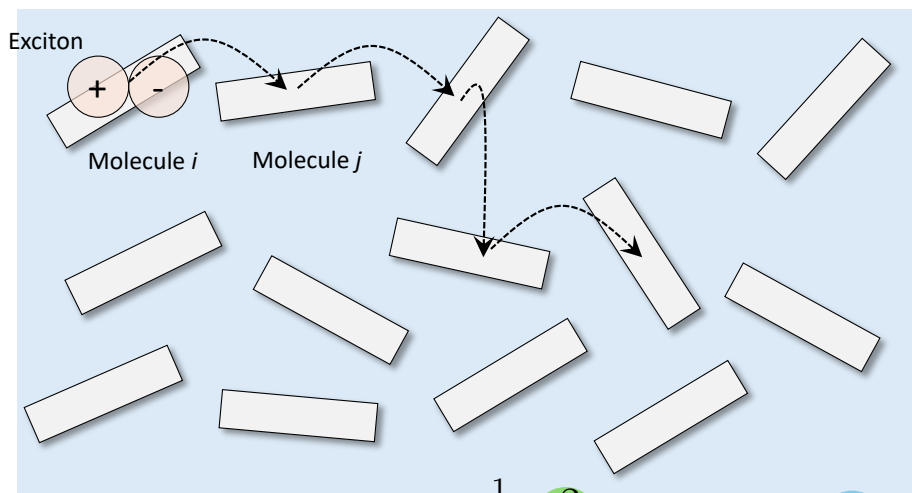
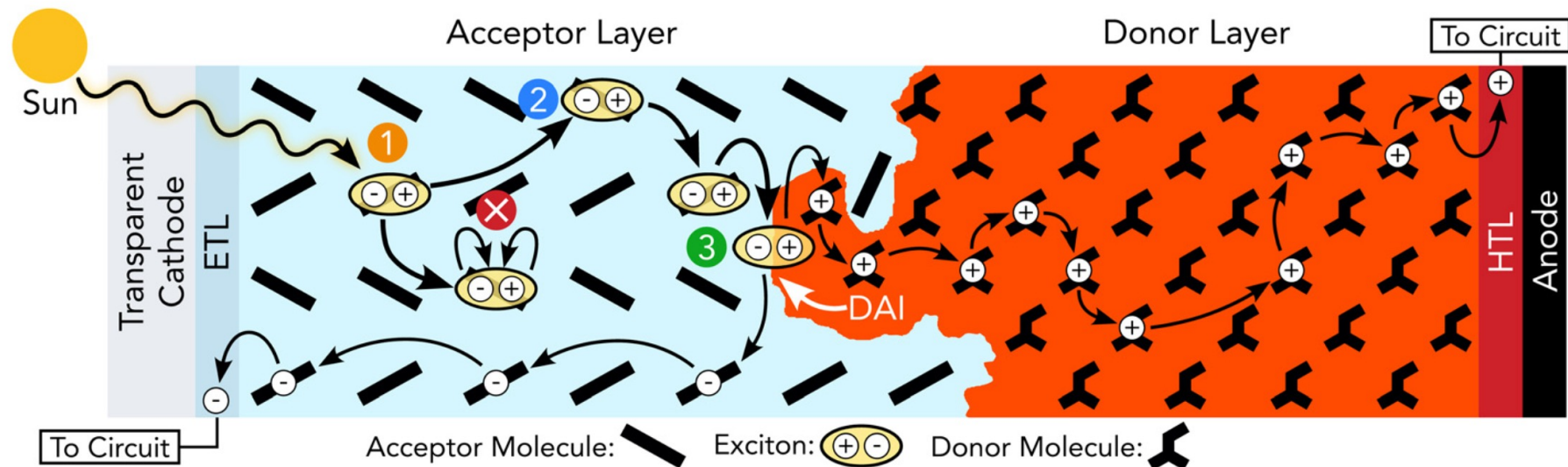
$$k_{ij} = \left( \frac{\pi}{\lambda k_B T} \right)^{\frac{1}{2}} \overset{\text{GNN-predicted exciton coupling}}{v_{ij}^2} \exp \left( - \frac{\overset{\text{Reorganization energy (TDDFT-calculated)}}{\lambda}}{4 k_B T} \right)$$



Organic solar cell

([www.solarreviews.com/blog/organic-solar-cells/](http://www.solarreviews.com/blog/organic-solar-cells/))

## Summary of part 2



$$k_{ij} = \left( \frac{\pi}{\lambda k_B T} \right)^{\frac{1}{2}} \frac{v_{ij}^2}{\hbar} \exp \left( -\frac{\lambda}{4 k_B T} \right)$$

- Kinetic Monte Carlo simulations *might* be used for screening organic semiconducting materials. However, the computational times required for exciton hopping rates need to be significantly reduced first.
- We created a new method for quickly computing exciton coupling parameters. It combines the atomic transition charge approximation and a graph neural network.
- The method generalizes widely across different molecule types.

**References:** Wechwithayakhlung *et al.* *J. Chem. Phys.* **158**, 2023, 204106  
 Weal *et al.* *J. Mater. Chem. C.* **24**, 2024, 8748  
 Weal *et al.* *J. Chem. Phys.* **163**, 2025, 024125

## Lecture topics

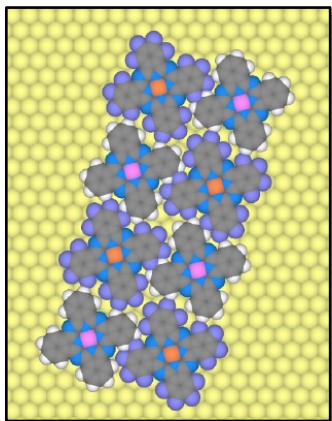
Simulation of on-surface molecular self-assembly

Machine learning for organic photovoltaic materials

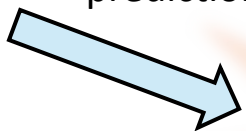
**What's next?**



## Molecular self-assembly research

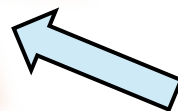


Interface structure  
prediction

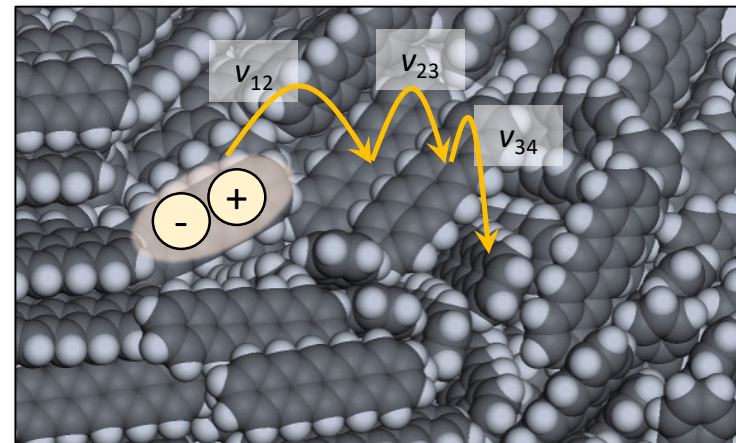


**New  
challenge!**

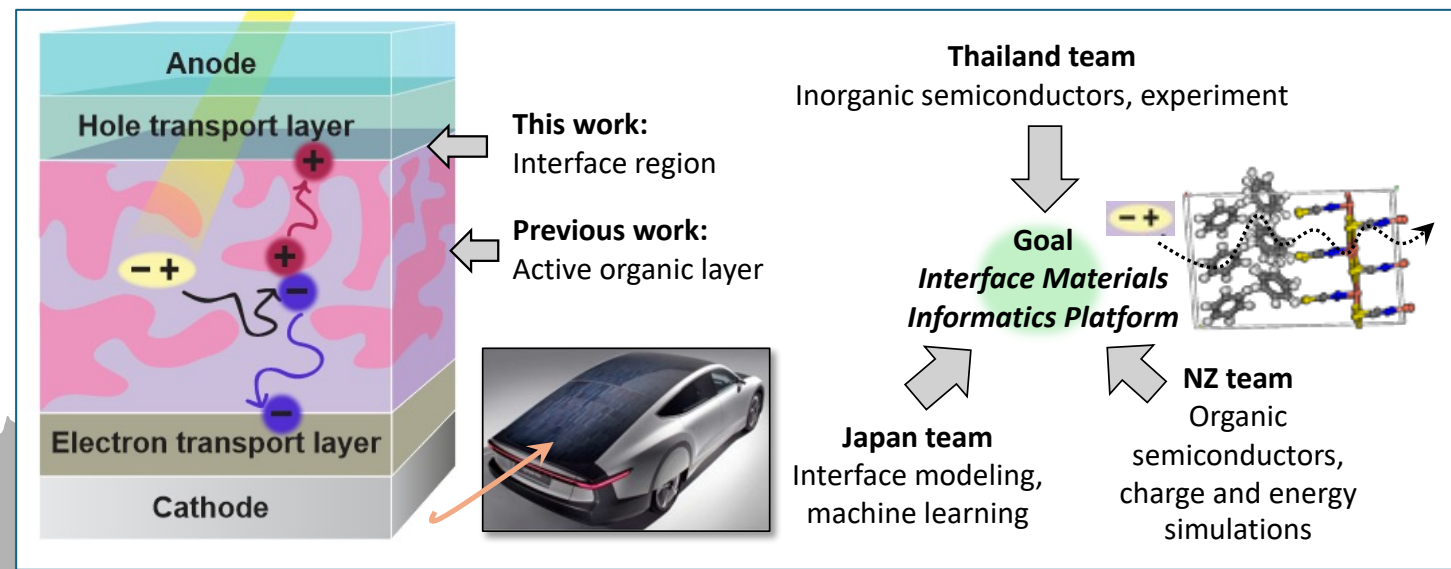
Exciton transport  
simulation



## Organic semiconductor research



# e-ASIA JRP Project: Interface Materials Informatics



**Pichaya Pattanasattayavong**  
(inorganic semiconductors)



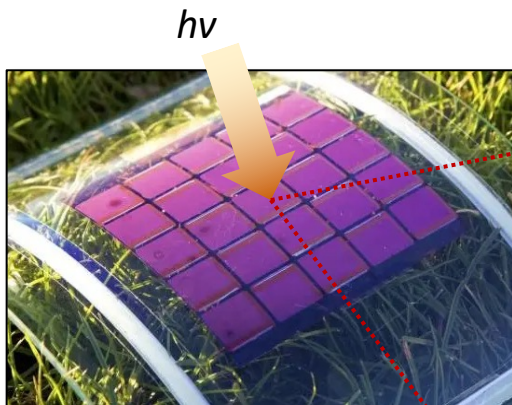
**Justin Hodgkiss**  
(organic semiconductors)



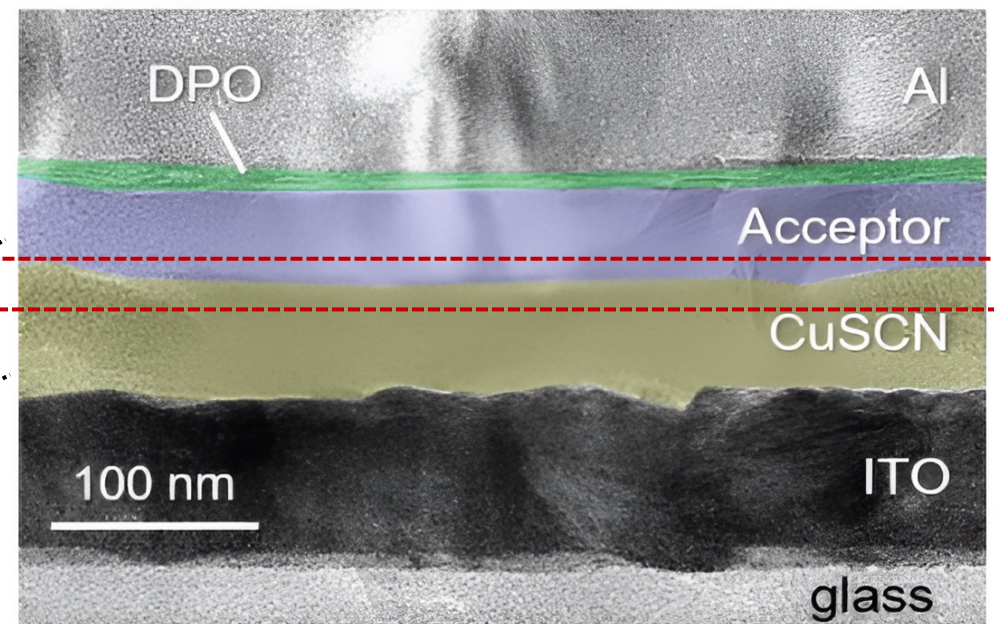
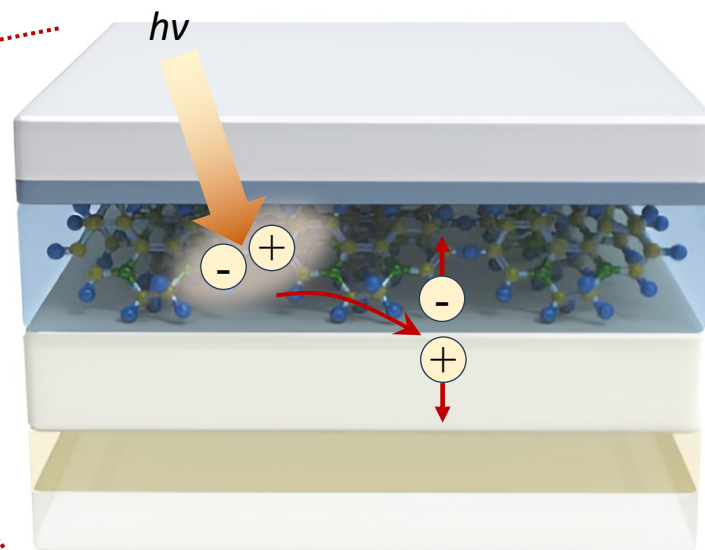
Started April 2025  
(Japan side)



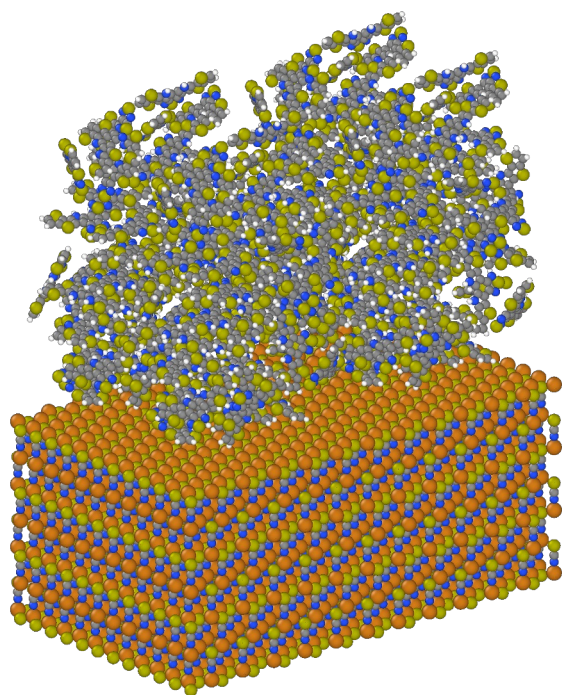
# Project focus: organic-inorganic interface in organic solar cells



([www.solarreviews.com/blog/organic-solar-cells/](http://www.solarreviews.com/blog/organic-solar-cells/))

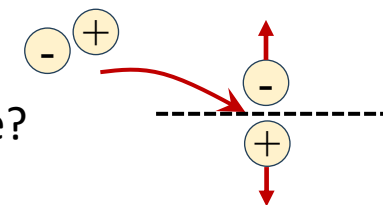


(<https://discovery.kaust.edu.sa/en/article/6432/guiding-the-way-to-improved-solar-cell-performance/>)



Can we predict the atomic-level structure of the interface?

Can we simulate charge separation at the interface?



Can we use these simulations to guide experimental solar cell fabrication?

Can experiment guide simulations?

(<https://www.solar.fau.de/research/devices/>)



**We hope to report good results soon!**



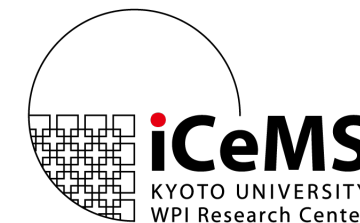
## Group members (2025)

Dr. Priya Dey  
Chio Hayashi  
Dr. Bandon Meza  
Dr. Maryam Nurhuda  
Dr. Shreya Rastogi  
Dr. James Scott



Geoff Weal

2024



2022

## Previous members

Dr. Geoff Weal  
Dr. Chayanit Wechwithayakhlung

Chayanit  
Wechwithayakhlung



## Collaborators (on-surface self-assembly)

Dr. Patrick Han  
Prof. Taro Hitosugi

## Collaborators (organic photovoltaics)

Dr. Yu Kaneko (DAICEL)  
Dr. Paul Hume (MacDiarmid Institute)  
Dr. Joshua Sutton (MacDiarmid Institute)  
Prof. Justin Hodgkiss (MacDiarmid Institute)

## Funding (on-surface self-assembly, organic photovoltaics)

JST PRESTO (2014 – 2018)  
JSPS Kakenki Kiban C (21K05003)  
JSPS Kakenhi Shingakujiyutsu Koubo (19H04574)  
JSPS Kakenhi Wakate (18K14126)  
JSPS Kakenhi Shingakujiyutsu Koubo (16H00879)  
JSPS Bilateral Projects / MBIE CATALYST  
JSPS Postdoctoral Fellowship for Foreign Researchers