

Computational Materials Design (CMD®) Workshop
Spintronic Design Course

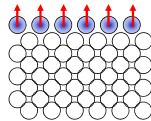
Spintronic · Design · Magnetic control II

Materials Design based on band structures

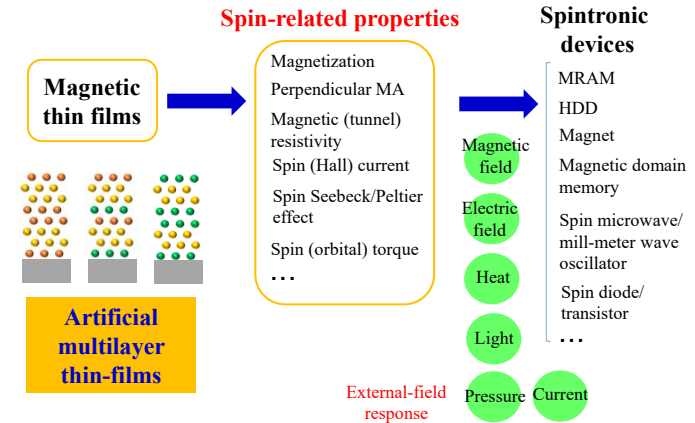
Kohji Nakamura
Mie University

Outline

- Introduction
- Calculation method
- Electronic structures and magnetism in bulk and at surfaces/thin films
- Control of magnetism by atomic-layer alignments and external electric field



Role of materials in spintronics



- Introduction
- **Calculation method**
- Electronic structures and magnetism in bulk and at surfaces/thin films
- Control of magnetism by tuning atomic-layer alignments and by external electric field

Basic: Density functional theory and Kohn-Sham equation

Many-body problem for electrons

$$\hat{\mathcal{H}}_e \Psi(\mathbf{R}_1, \mathbf{R}_2, \dots; \mathbf{r}_1, \mathbf{r}_2, \dots) = E \Psi(\mathbf{R}_1, \mathbf{R}_2, \dots; \mathbf{r}_1, \mathbf{r}_2, \dots)$$



① Adiabatic (Born-Oppenheimer) approximation

② Non-interacting one particle approximation

Kohn-Sham equation (one-electron equation)

$$\left[-\frac{\nabla^2}{2} + v_H(n(\mathbf{r})) + v_{XC}(n(\mathbf{r})) \right] \phi_b(\mathbf{r}) = \varepsilon_b \phi_b(\mathbf{r})$$

Basic: Case of non spin-polarized systems

Kohn-Sham equation in one-electron approximation

$$\left[-\frac{\nabla^2}{2} + v_H(n(\mathbf{r})) + v_{xc}(n(\mathbf{r})) \right] \phi_b(\mathbf{r}) = \varepsilon_b \phi_b(\mathbf{r})$$

↑
Local Density Approximation (LDA)

$$E_{xc}^{LDA}[n] \approx \int n(\mathbf{r}) \varepsilon_{xc}^{LDA}(n(\mathbf{r}))$$

$$E_{xc}^{GGA}[n] \approx \int n(\mathbf{r}) \varepsilon_{xc}^{GGA}(n(\mathbf{r}), \nabla n(\mathbf{r}))$$

Electron density

$$n(\mathbf{r}) = \langle \psi(\mathbf{r}) | \hat{n} | \psi(\mathbf{r}) \rangle$$

$$\sim |\phi_1(\mathbf{r})|^2 + |\phi_2(\mathbf{r})|^2 + \dots = \sum_{b=1}^N |\phi_b(\mathbf{r})|^2$$

Basic: Case of spin-polarized systems

Kohn-Sham equation in one-electron approximation

$$\left[-\frac{\nabla^2}{2} + v_H(n(\mathbf{r})) + v_{xc}(n(\mathbf{r}), m(\mathbf{r})) \right] \phi_{b,\sigma}(\mathbf{r}) = \varepsilon_{b,\sigma} \phi_{b,\sigma}(\mathbf{r})$$

↑
Local Spin Density Approximation (LSDA)

Electron density

$$n(\mathbf{r}) \sim \sum_{b=1}^{N_\uparrow} |\phi_{b,\uparrow}(\mathbf{r})|^2 + \sum_{b=1}^{N_\downarrow} |\phi_{b,\downarrow}(\mathbf{r})|^2$$

$$= n_\uparrow(\mathbf{r}) + n_\downarrow(\mathbf{r})$$

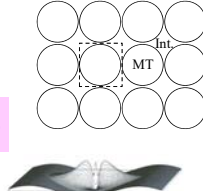
$$m(\mathbf{r}) = n_\uparrow(\mathbf{r}) - n_\downarrow(\mathbf{r})$$

Basic: Variational principle method

$$\psi_{k,\sigma}(\mathbf{r}) = \sum_{\mathbf{G}}^{<K_{\max}} C_{\mathbf{k}+\mathbf{G},\sigma} |\phi_{\mathbf{k}+\mathbf{G},\sigma}\rangle \quad \Rightarrow \quad \text{Eigenvalue problem} \quad \hat{H}_{KS} |\psi_{k,\sigma}\rangle = \varepsilon_b |\psi_{k,\sigma}\rangle$$

Example: Linearized Augmented Plane Wave basis

$$\phi_{\mathbf{k}+\mathbf{G},\sigma}(\mathbf{r}) = \begin{cases} \text{Interstitial} \\ e^{i(\mathbf{k}+\mathbf{G})\cdot\mathbf{r}} \\ \text{MT sphere} \\ \sum_{l=0}^{l_{\max}} [A_{l,\sigma}(\mathbf{k}+\mathbf{G}) u_{l,\sigma}(r) + B_{l,\sigma}(\mathbf{k}+\mathbf{G}) \bar{u}_{l,\sigma}(r)] iY_l(r) \\ \left[-\frac{\partial}{\partial r^2} + \frac{l(l+1)}{r^2} + v_{l=0,\sigma}(r) - E_{l,\sigma} \right] r u_{l,\sigma}(r) = 0 \end{cases}$$



E. Wimmer *et al.*, PRB 24, 864, (1981); PRB 26, 4571 (1982)

Basic: SOC and Second variational method

Spin-orbit coupling

$$\hat{H}_{SOC} = \frac{1}{4c^2} \frac{1}{r} \frac{dV_{l=0}^\sigma}{dr} \hat{l} \cdot \hat{\sigma} = \xi(r) \begin{pmatrix} \hat{l}_z & \hat{l}_- \\ \hat{l}_+ & -\hat{l}_z \end{pmatrix} \quad \begin{array}{l} \hat{l} : \text{Angular momentum operator} \\ \hat{\sigma} : \text{Dirac operator} \end{array}$$

① Redefine basis by the KS eigenfunctions

$$|\psi_b\rangle = C_{b,\uparrow} |\psi_{b,\uparrow}\rangle + C_{b,\downarrow} |\psi_{b,\downarrow}\rangle$$



② Generate SOC matrix

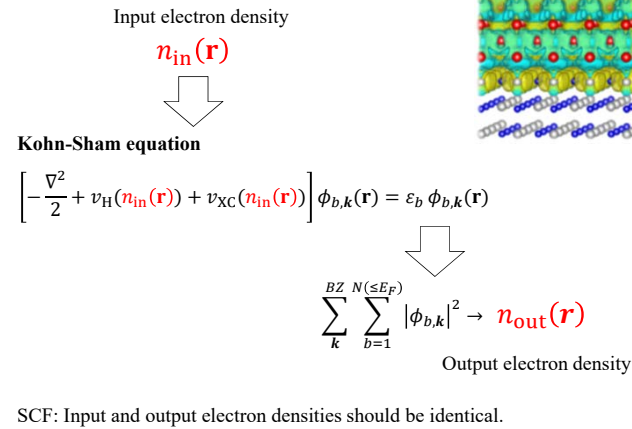
$$\hat{H} = \hat{H}_{KS} + \hat{H}_{SOC} \\ = \begin{pmatrix} \hat{H}_{KS}^{\uparrow} & 0 \\ 0 & \hat{H}_{KS}^{\downarrow} \end{pmatrix} + \xi(r) \begin{pmatrix} \hat{l}_z & \hat{l}_- \\ \hat{l}_+ & -\hat{l}_z \end{pmatrix}$$

③ Eigenvalue problem

$$(\hat{H}_{KS} + \hat{H}_{SOC}) |\psi_b\rangle = \varepsilon_b |\psi_b\rangle$$

Li *et al.*, PRB 42, 5433 (1990).

Basic: Self-consistent field (SCF)



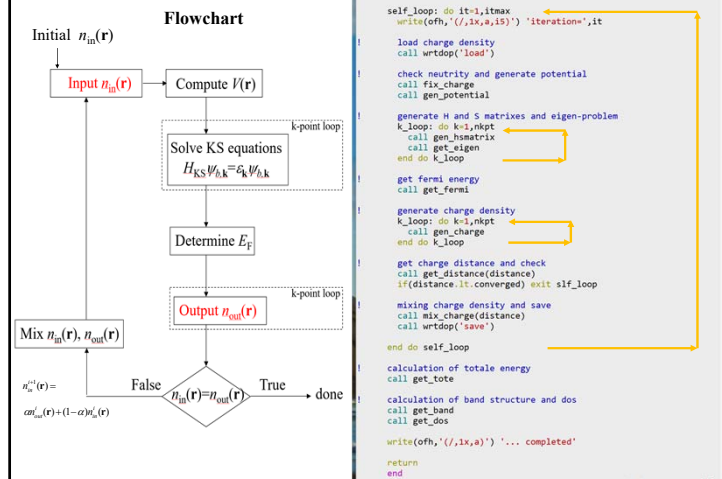
Basic: Total energy

$$E_{\text{Total}} = \sum_{i=1}^{N(\leq E_F)} \varepsilon_b \quad \text{Kohn-Sham eigenvalues}$$

$$-\frac{1}{2} \left[\int v_H(\mathbf{r}) n(\mathbf{r}) d\mathbf{r} + \sum_{\sigma} \int \{v_{XC,\sigma}(\mathbf{r}) - \varepsilon_{XC,\sigma}(\mathbf{r})\} n_{\sigma}(\mathbf{r}) d\mathbf{r} \right]$$

$$-\frac{1}{2} \sum_{\mu} Z_{\mu} V_{C,\mu}(\mathbf{R}_{\mu}) \quad \text{Double counting term}$$

Coding: Program



Coding: Execution

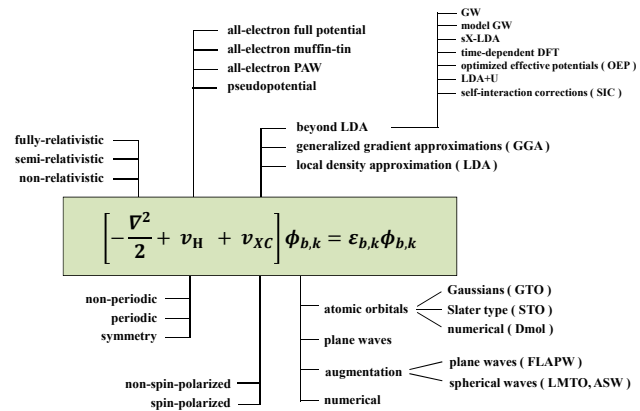
Charge distance :

$$\Delta = \frac{1}{\Omega} \int_{\text{cell}} \sqrt{(n_{out}(\mathbf{r}) - n_{in}(\mathbf{r}))^2} d\mathbf{r}$$

```

[kohji@red27 work]$ FLchk
** Total energy
total energy for it= 9: -75.9859905662 htr
*** Charge distance
distance of charge densities for it= 1: 15.418023 me/bohr**3
distance of charge densities for it= 2: 12.714097 me/bohr**3
distance of charge densities for it= 3: 10.269848 me/bohr**3
distance of charge densities for it= 4: 2.465512 me/bohr**3
distance of charge densities for it= 5: 1.862776 me/bohr**3
distance of charge densities for it= 6: 0.363891 me/bohr**3
distance of charge densities for it= 7: 0.086428 me/bohr**3
distance of charge densities for it= 8: 0.079364 me/bohr**3
distance of charge densities for it= 9: 0.025358 me/bohr**3
***
FLAPW calculations were done
  
```

Coding: Implementation



➤ Introduction

➤ Calculation method

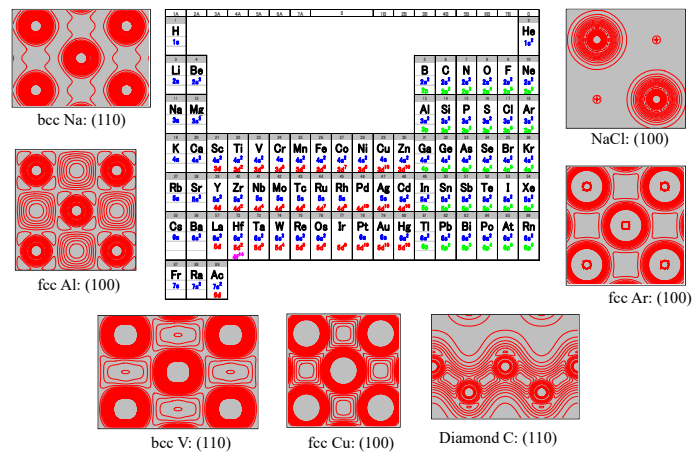
➤ Electronic structures and magnetism

in bulk and at surfaces/thin films

➤ Control of magnetism

by tuning atomic-layer alignments and by external electric field

Demo: Electron density



Demo: Band structure of non-magnetic system

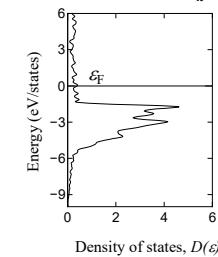
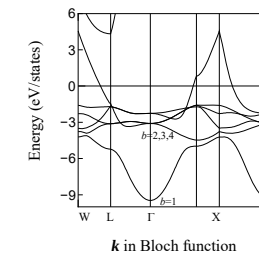
fcc Cu (non spin-polarized system)

Kohn-Sham equation

$$-\frac{\nabla^2}{2} + v_H + v_{xc} |\phi_{b,k}\rangle = \epsilon_{b,k} |\phi_{b,k}\rangle$$

Wave function of system

$$|\psi\rangle = \sum_k \sum_b c_{b,k} |\phi_{b,k}\rangle$$



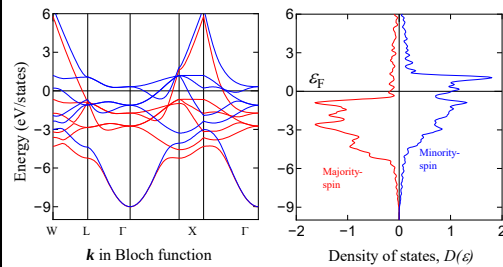
Demo: Band structure of ferromagnetic system

Kohn-Sham equation

fcc Co (spin-polarized system)

$$-\frac{\nabla^2}{2} + v_H + v_{xc,\uparrow} |\phi_{b,k,\uparrow}\rangle = \varepsilon_{b,k,\uparrow} |\phi_{b,k,\uparrow}\rangle \quad \text{spin-up state}$$

$$-\frac{\nabla^2}{2} + v_H + v_{xc,\downarrow} |\phi_{b,k,\downarrow}\rangle = \varepsilon_{b,k,\downarrow} |\phi_{b,k,\downarrow}\rangle \quad \text{spin-down state}$$



Magnetic moment

$$m = \int_{-\infty}^{\varepsilon_F} \mathbf{D}_{\uparrow}(\varepsilon) d\varepsilon - \int_{-\infty}^{\varepsilon_F} \mathbf{D}_{\downarrow}(\varepsilon) d\varepsilon$$

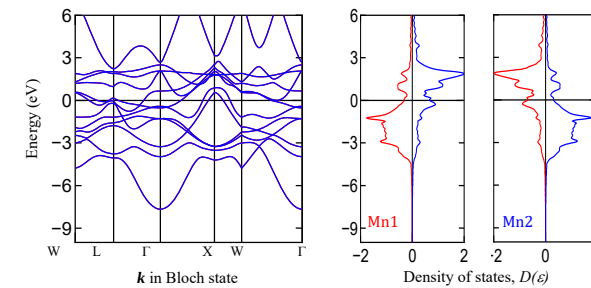
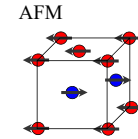
Demo: Band structure of antiferromagnetic system

Kohn-Sham equation

fcc Mn (spin-polarized system)

$$-\frac{\nabla^2}{2} + v_H + v_{xc,\uparrow} |\phi_{b,k,\uparrow}\rangle = \varepsilon_{b,k,\uparrow} |\phi_{b,k,\uparrow}\rangle$$

$$-\frac{\nabla^2}{2} + v_H + v_{xc,\downarrow} |\phi_{b,k,\downarrow}\rangle = \varepsilon_{b,k,\downarrow} |\phi_{b,k,\downarrow}\rangle$$

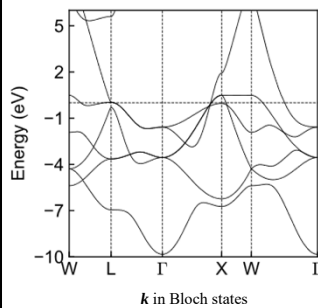


Demo: Band structure of Pt with SOC

Pt (non spin-polarized system)

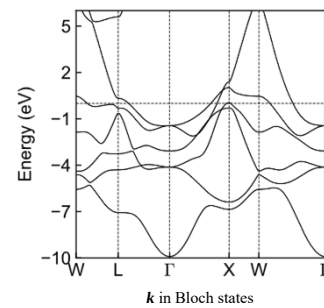
Without SOC

$$\hat{H}_{KS} |\phi_{b,k}\rangle = \varepsilon_{b,k} |\phi_{b,k}\rangle$$



With SOC

$$(\hat{H}_{KS} + \hat{H}_{SOC}) |\psi_{b,k}\rangle = \varepsilon_{b,k} |\psi_{b,k}\rangle$$



Spin-orbit coupling-induced properties

Fundamental properties

- Magnetic spin/orbital and dipole moments
- Exchange interaction
- Magnetocrystalline anisotropy (MCA)
- Dzyaloshinskii-Moriya interaction (DMI)

Conductivity

- Anomalous hall conductivity
- Spin/orbital Hall conductivity
- Magneto-optical conductivity

Dynamics

- Magnetic damping constant
- Spin-orbit torque

External field

- Magnetic field
- Electric field

Spin textures

- Non-collinear magnetic structures

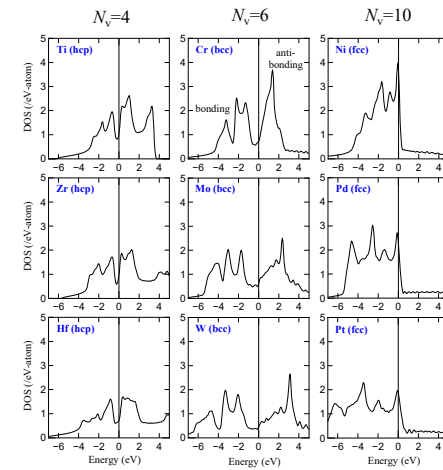
Magnetism

Question:
Why do ferromagnetism
appear only in Fe, Co
and Ni?

1A	2A	3A	4A	5A	6A	7A	8	10	2B	3B	4B	5B	6B	7B	0		
H															He		
Li	Be											B	C	N	O	F	Ne
Na	Mg											Al	Si	P	S	Cl	Ar
K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr
Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe
Cs	Ba	La	Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi	Po	At	Rn

	Magnetic order	Transition Temperature(K)	Crystal structure	Number of valence
Fe	Ferromagnetic	1043	bcc	8
Co	Ferromagnetic	1388	hcp	9
Ni	Ferromagnetic	627	fcc	10

Band structures of transition-metals



Trends in band structures in transition metals

1) Bandwidth

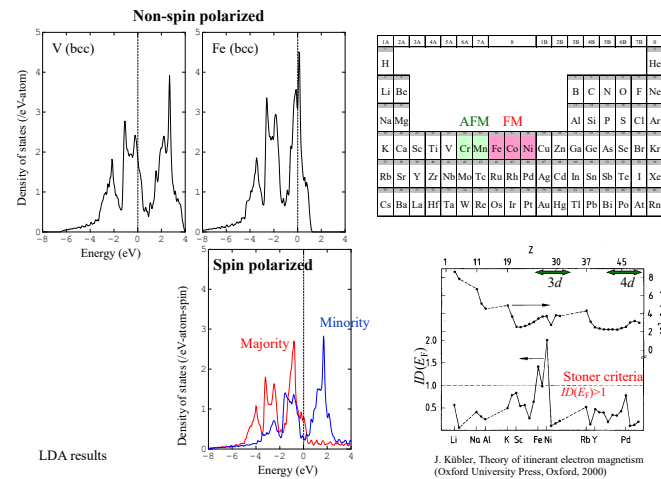
$$W^{3d} < W^{4d} < W^{5d}$$

2) Density of states at E_F

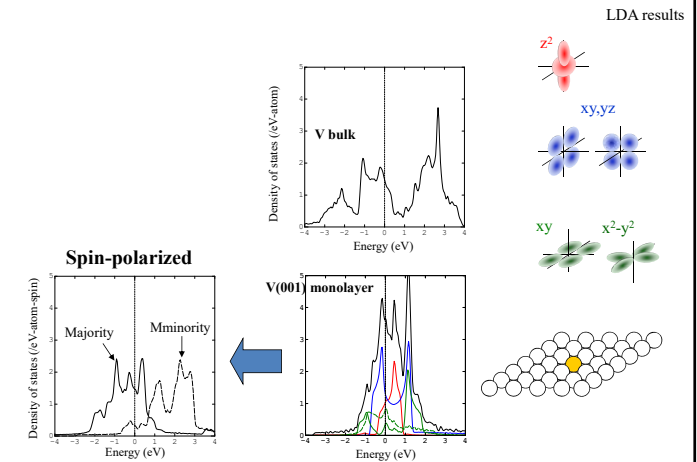
$$D^{3d}(E_F) > D^{4d}(E_F) > D^{5d}(E_F)$$

J. Kübler, Theory of itinerant electron magnetism (Oxford University Press, Oxford, 2000)

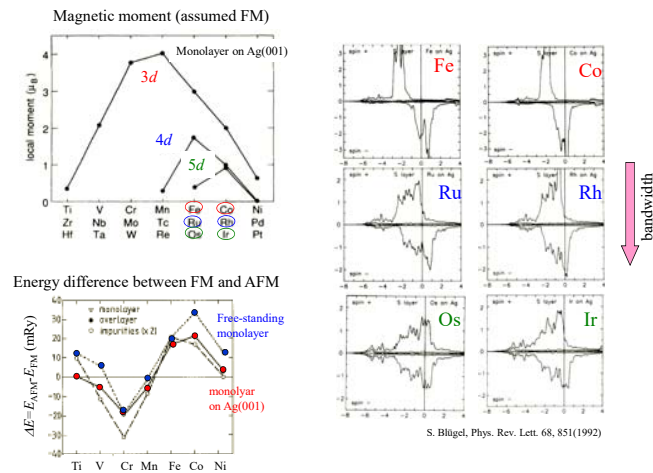
Magnetism in transition metals



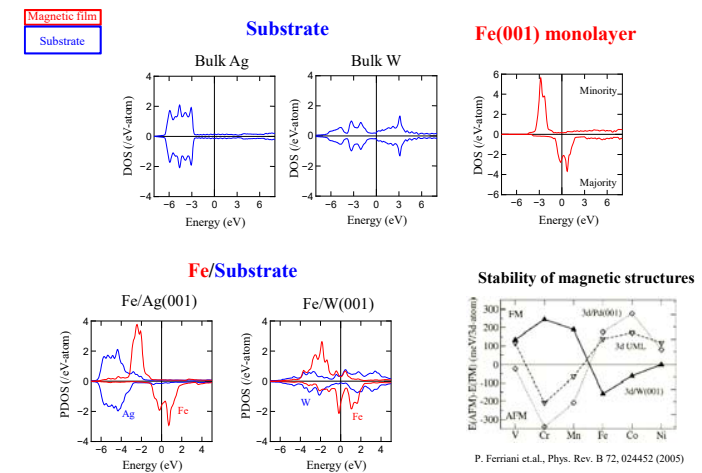
Band structures and spin polarization in transition-metal layers



Ferromagnetism in transition-metal layers

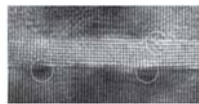


Effect of substrates to magnetism



Magnetocrystalline anisotropy (MCA)

MTJ device



[1]

Ferromagnetic
Insulator
Ferromagnetic

- ✓Non-volatility
- ✓Large capacity
- ✓High-speed

Many Applications

- MRAM
- Magnetic sensor

[1] S. Yuasa, et al., Nature Materials 3, 868 (2004).

In-plane MCA



Advanced applications

Out-of-plane MCA

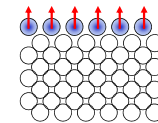


How to design thin films with large PMCA

Magnetocrystalline anisotropy

Spin-orbit coupling

$$H_{SOC} = \xi(r) \mathbf{l} \cdot \boldsymbol{\sigma}$$

SOC strength, $\xi \sim 50$ meV for 3d-metals

Magnetocrystalline anisotropy energy

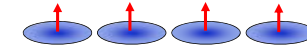
$$E_{MCA} = E(\rightarrow) - E(\uparrow)$$

 ~ 0.1 meV/atom $E_{MCA} < 0$: in-plane MCA

Force theorem

$$\approx \langle E[\rho_0, \mathbf{m}_0(\rightarrow)] - E[\rho_0, \mathbf{m}_0(\uparrow)] \rangle$$

$$= \sum_{\mathbf{k}} \left[\sum_{i \in \sigma} \epsilon_{i,\mathbf{k}}(\rightarrow) - \sum_{i \in \sigma} \epsilon_{i,\mathbf{k}}(\uparrow) \right]$$

 $E_{MCA} > 0$: Perpendicular MCA

MCA energy in 3d transition metal monolayers

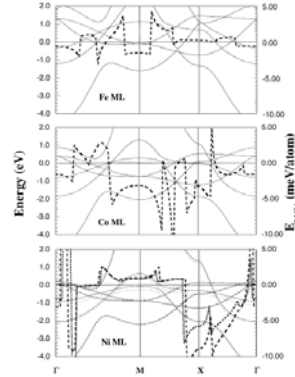
Perpendicular MCA



System	Method	E_{MCA} (meV)
★ Fe(0 0 1)	FLAPW-GGA	+ 0.50
	FLAPW-LDA	+ 0.21
Co(0 0 1)	FLAPW-GGA	- 1.49
	FLAPW-LDA	- 1.42
Ni(0 0 1)	FLAPW-GGA	- 0.77
	FLAPW-LDA	- 1.64

($a=a_{Cu}$)

R. Wu, A.J. Freeman, J. Magn. Magn. Mater. 200, 498(1999)



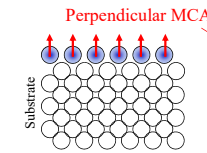
Effect of substrates to MCA

$$E_{MCA} = E(\rightarrow) - E(\uparrow)$$

Table 2
Calculated MCA energies for selected transition magnetic thin films

System	Method	E_{MCA} (meV)
Co/Cu(0 0 1)	FLAPW-LDA-ST-UR [46]	- 0.38
	FLAPW-LDA-SC-RL [44,45]	- 0.36
	FLAPW-LDA-ST-RL [42,43]	- 0.09
	FLAPW-GGA-TQ-RL [41]	- 0.61
	SKKR-LDA-UR [47]	- 0.38
	Experiment [35]	- 0.37
★ Cu/Cu/Cu(0 0 1)	FLAPW-LDA-ST-UR [46]	- 0.01
	FLAPW-GGA-TQ-RL [41]	+ 0.54
	SKKR-LDA-UR [48]	+ 0.85
	Experiment [35]	+ 0.1
Co/Cu(1 1 1)	FLAPW-LDA-ST-RL [48]	- 0.30
Co/Pd(0 0 1)	FLAPW-LDA-ST-RL [76]	- 0.91
★ Co/Pd(1 1 1)	FLAPW-LDA-ST-RL [76]	+ 0.25
★ Pd/Co/Pd(0 0 1)	FLAPW-LDA-ST-UR [118]	+ 0.56
Ni/Cu(0 0 1)	FLAPW-LDA-ST-UR	- 0.69
★ 2Ni/Cu(0 0 1)	FLAPW-LDA-ST-UR	+ 0.33
★ 3Ni/Cu(0 0 1)	FLAPW-LDA-ST-UR	+ 0.08
4Ni/Cu(0 0 1)	FLAPW-LDA-ST-UR	- 0.06
Cu/Fe/Cu(0 0 1)	SKKR-LDA [119]	- 0.41
	LMTO-LDA [120,121]	- 0.43
★ Fe/Au(0 0 1)	FLAPW-LDA [122]	+ 0.57
	SKKR-LDA [123]	+ 0.56

R. Wu, A.J. Freeman, J. Magn. Magn. Mater. 200, 498(1999)



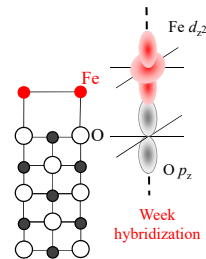
Perpendicular MCA at Fe/MgO(001) interface: Theory

$$E_{MCA} = E(\rightarrow) - E(\uparrow)$$

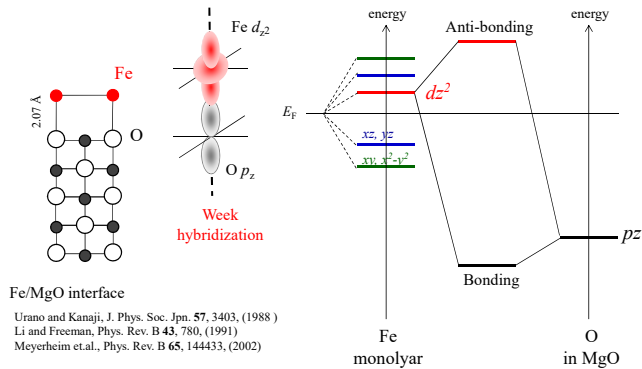
	E_{MCA} (meV/ a^2)
Free-standing Fe monolayer	0.19
Fe/MgO	1.28
Au ₃ /Fe ₃ /MgO	0.94

LDA
Nakamura et.al., PRB 81, 220409, (2010)

Origin of interfacial PMA



Fe d_{z^2} - O p_z hybridization at Fe/MgO interface



Fe/MgO interface
Urano and Kanaji, J. Phys. Soc. Jpn. 57, 3403, (1988)
Li and Freeman, Phys. Rev. B 43, 780, (1991)
Meyerheim et.al., Phys. Rev. B 65, 144433, (2002)

Shimabukuro, Physica E, 42, 1014 (2010).
Nakamura et.al., PRB 81, 220409, (2010)

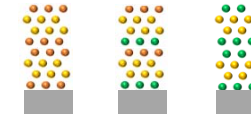
- Introduction
- Calculation method
- Electronic structures and magnetism
in bulk and at surfaces/thin films

➤ **Control of magnetism**

by tuning atomic-layer alignments and by external electric field

Control of MCA by tuning atomic-layer alignments

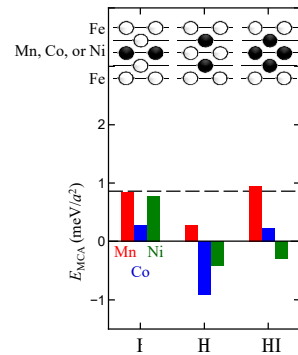
Interfaces, superlattices, multilayers



Thin films with large perpendicular MCA, consisting of only 3d metals

1s	2s	2p	3s	3p	3d	4s	4p	4d	4f	5s	5p	5d	5f	6s	6p	6d	6f	7s	7p	7d	7f	8s	8p	8d	8f	9s	9p	9d	9f
H																													He
Li	Be													B	C	N	O	F	Ne										
Na	Mg													Al	Si	P	S	Cl	Ar										
K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr												
Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe												
Cs	Ba	La	Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi	Po	At	Rn												

MCA of 3d transition metal thin films with five-atomic-layer



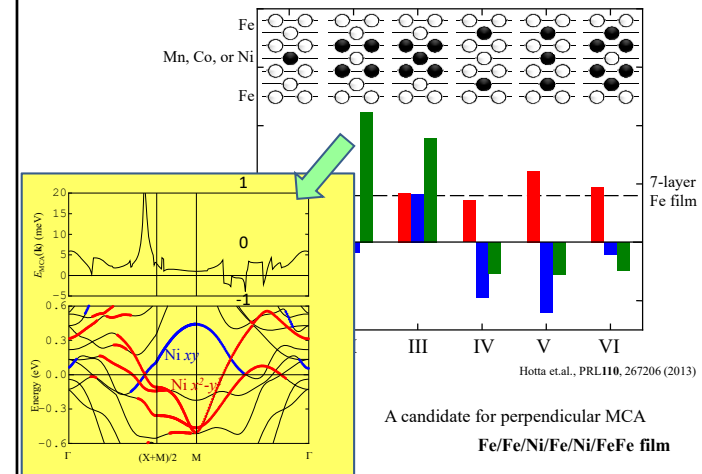
Hotta et.al., PRL110, 267206 (2013)

- Assumed
- ✓ Bcc-layer stacking
 - ✓ Surface-layer to be Fe
 - ✓ Symmetric atomic alignment along z-axis

Pur Fe five-layer film

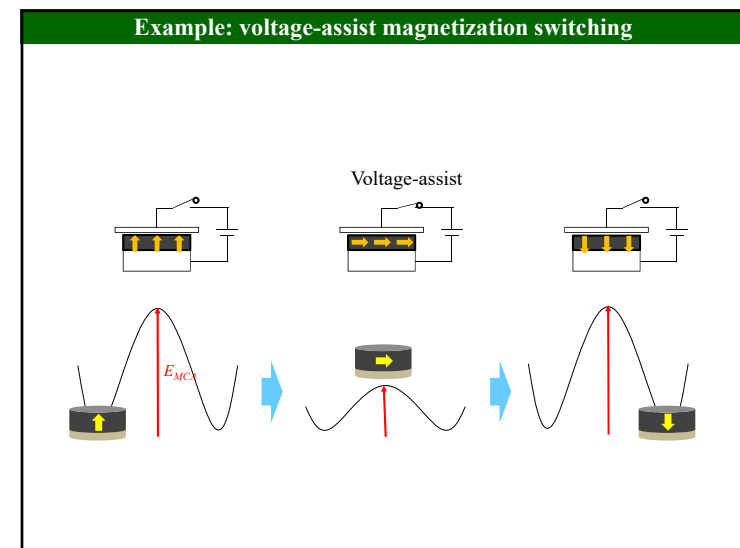
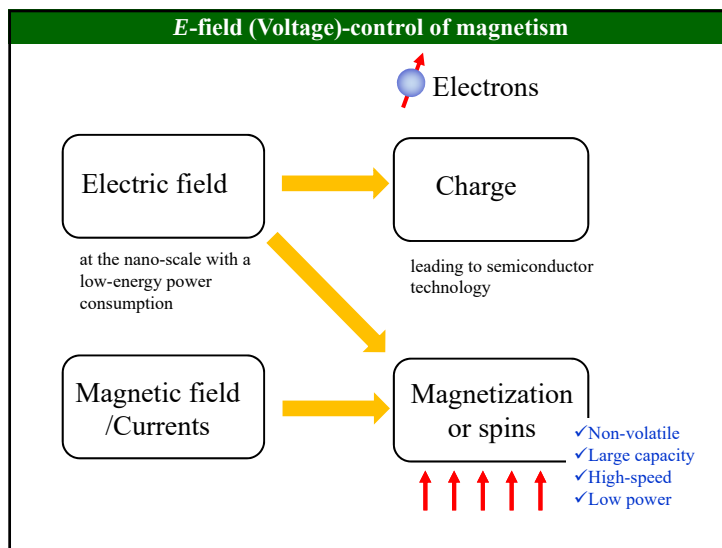
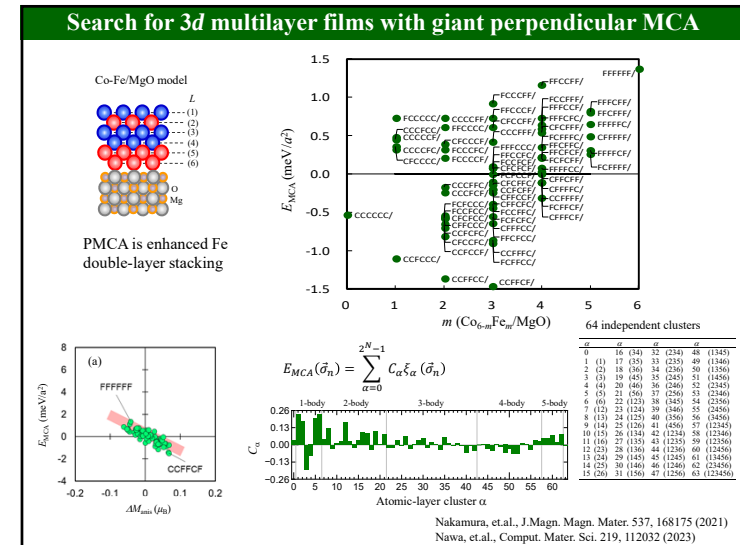
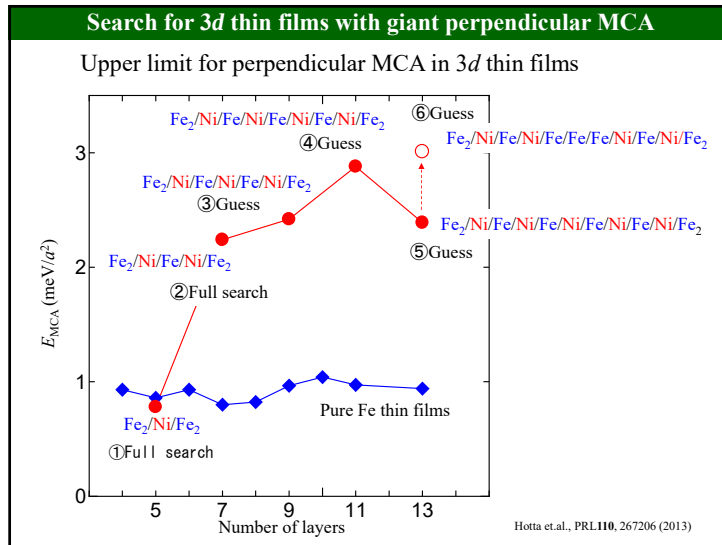
It shows no perpendicular MCA

MCA of 3d transition metal thin films with seven-atomic-layer



Hotta et.al., PRL110, 267206 (2013)

A candidate for perpendicular MCA
Fe/Fe/Ni/Fe/Ni/Fe/Fe film



Control of magnetism of metal thin films by electric field

➤ Magnetocrystalline anisotropy (MCA)

Weisheit et al., Science 315, 349 (2007)
 Maruyama et al., Nature Nanotech. 4, 349 (2009)
 Duan et al., PRL101, 137201 (2008)
 Nakamura et al., PRL102, 187201 (2009); PRB 80, 172402 (2009)
 Tsujikawa et al., PRL102, (2009)

➤ Curie temperature (T_C)

Chiba et al., Nat. Mater. 10, 853-856 (2011)
 Oba et al., PRL114, 107202 (2015)

➤ Dzyaloshinskii-Moriya interaction (DMI)

Nawaoka et al., Appl. Phys. Express 8, 063004(2015)
 Nakamura et al., (2015)

➤ Magnetic moments

Obinata et al., Sci. Rep. 5, 14303 (2015)

➤ Magneto-optical conductivity

Hibino et al., (2015)
 Nakamura et al., J. Korean Phys. Soc. 63, 612 (2013)

➤ Magnetic dumping

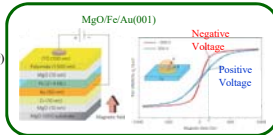
Okada et al., Appl. Phys. Lett. 105, 052415 (2014)

➤ Magnetic phase-transition (bcc Fe vs fcc Fe)

Gerhard et al., Nat. Nanotech., 5, 792 (2010)

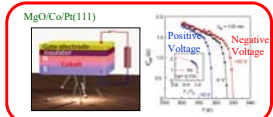
➤ etc.

Magnetocrystalline anisotropy



Maruyama et al., Nature Nanotech. 4, 349 (2009)

Curie temperatures

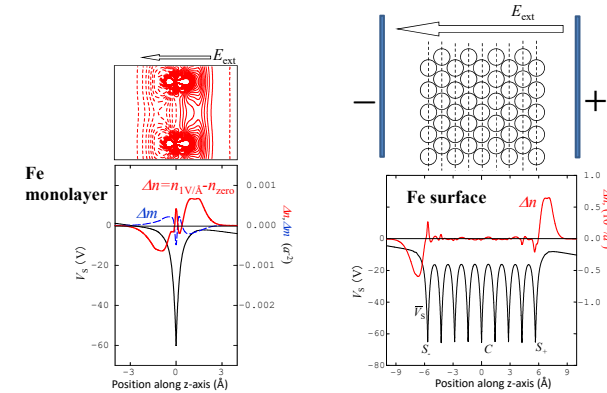


Chiba et al., Nat. Mater. 10, 853-856 (2011)

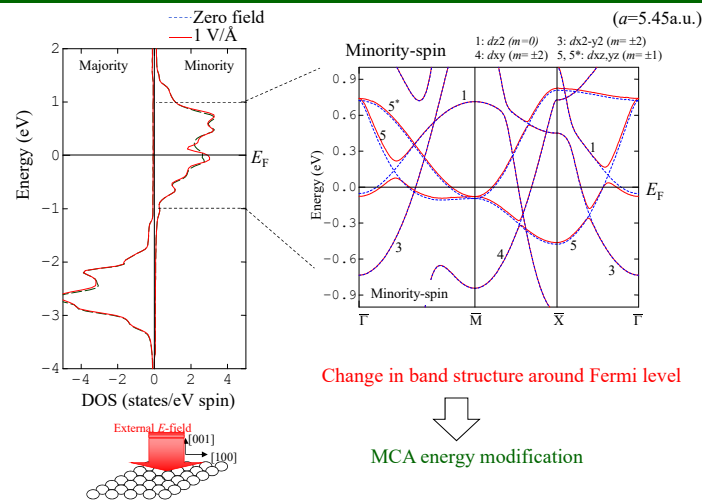
Electric-field-induced charges at surfaces/interfaces

➤ Redistribution in charge/spin densities at surfaces/interfaces by E-field.

Duan et al., PRL101, 137201 (2008),
 Nakamura et al., PRL102, 187201 (2009); PRB 80, 172402 (2009),
 Tsujikawa et al., PRL102, (2009).

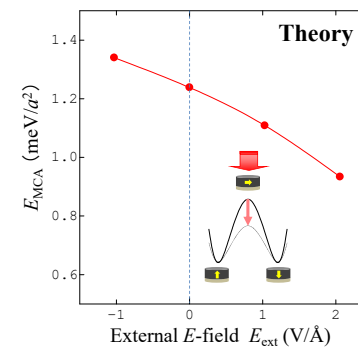


Modification of band structure of Fe monolayer by electric field

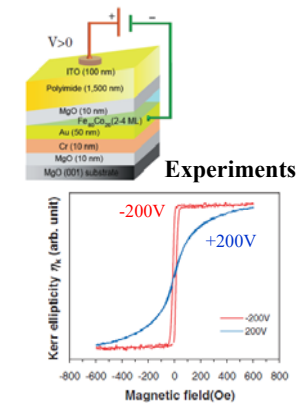


Modification of MCA energy at Fe/MgO(001)

$$E_{MCA} = E(\rightarrow) - E(\uparrow)$$



Nakamura et al., PRB 81, 220409 (2010)

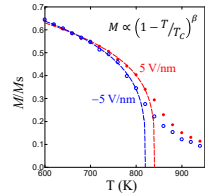
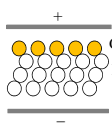


Shito et al., Appl. Phys. Express 2, 063001 (2009)

Modification of Curie temperature of Co/Pt(111)

Oba et al., PRL 114, 107202 (2015)

Co/Pt(111)



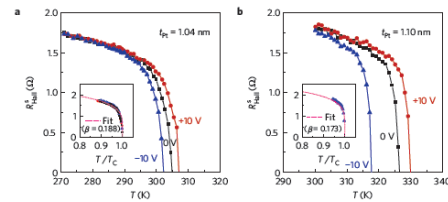
	Theory	Experiments
T_C	~830	~320
ΔT_C	~20K	~10K
β	~0.26	~0.18

Agreed with experiments qualitatively

MgO/Co/Pt(111)



Chiba et al., Nat. Mater. 10, 853-856 (2011)



Summary

Spintronic · Design · Magnetic control

Outline

- Introduction
- Calculation method
- Electronic structures and magnetism in bulk and at surfaces/thin films
- Control of magnetism by atomic-layer alignments and external electric field

