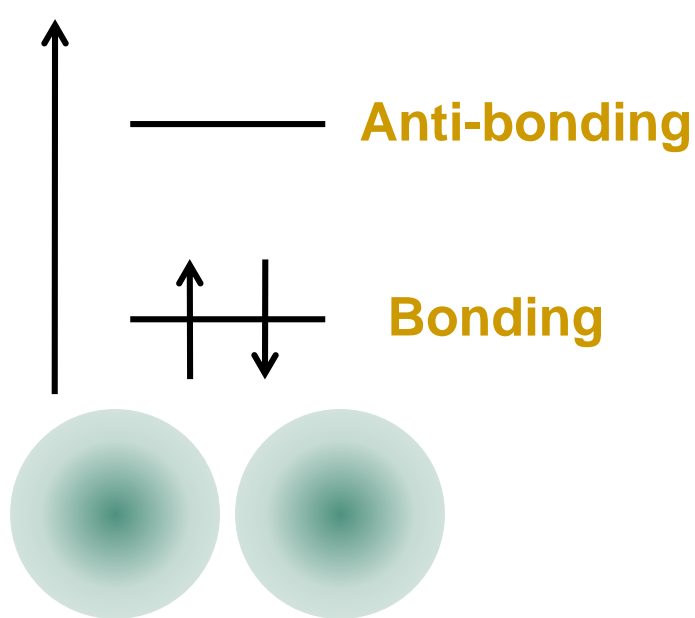# The 37th CMD workshop "Advanced study with ES-opt"
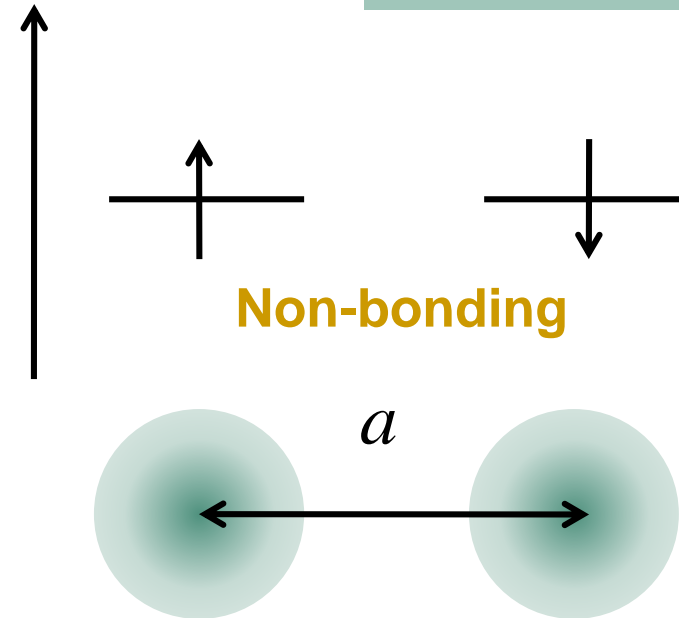
Multi-reference DFT

Koichi Kusakabe

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Electronic structure of diatomic molecules

Anti-bonding

Bonding

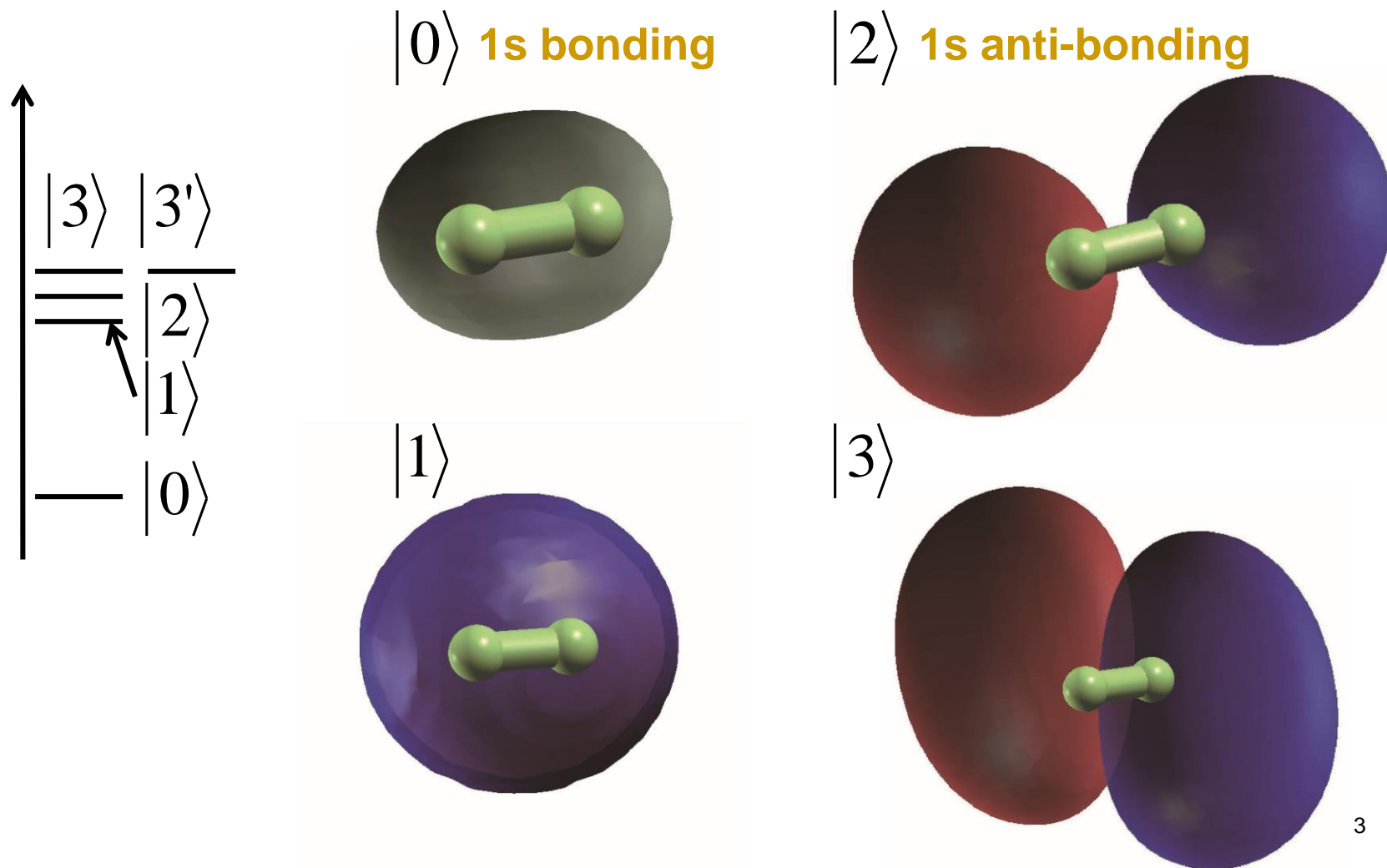**Non-bonding**

a

Molecular orbital picture
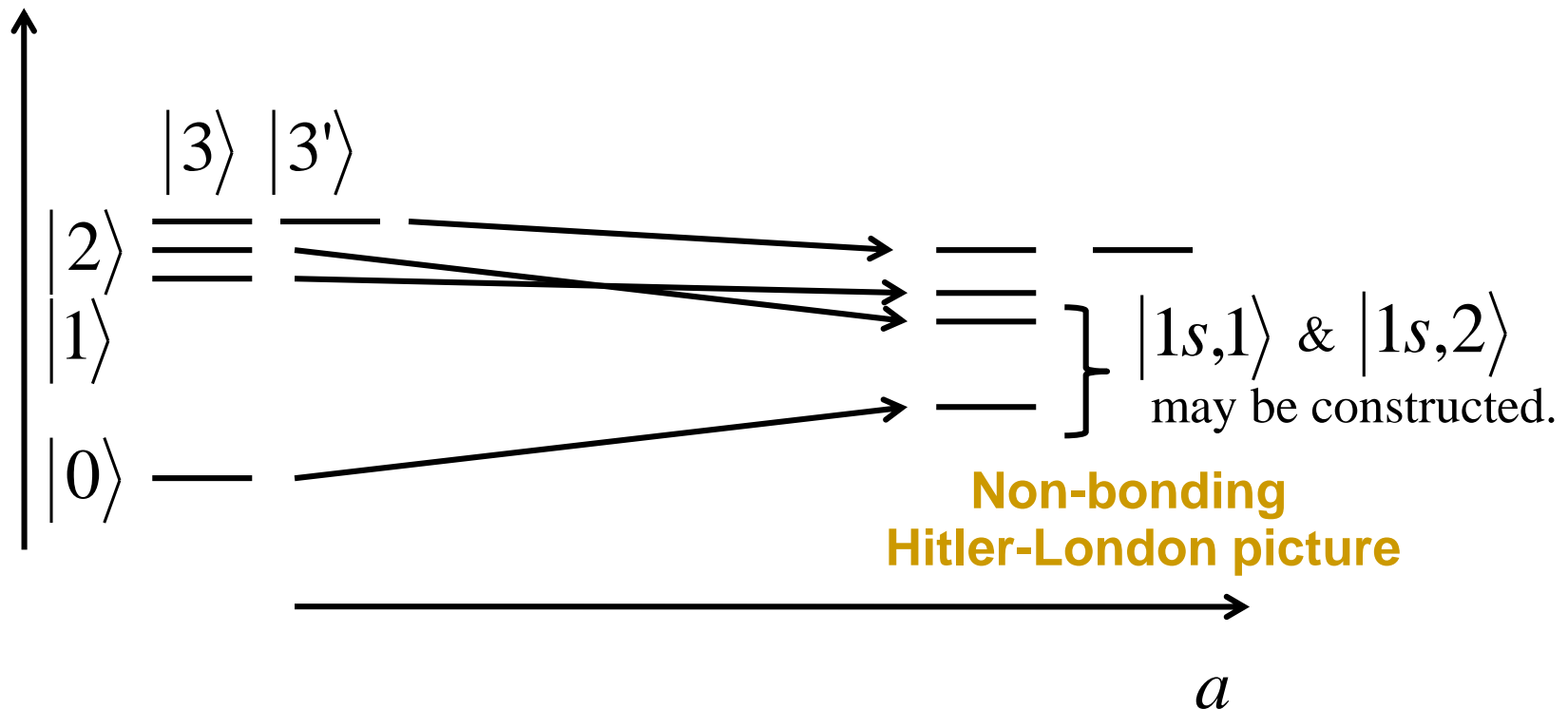
Heitler-London picture

Known results in DFT-LDA
• **Molecular orbital picture is given for any inter-atomic distance.**
• **The stable electronic state in equilibrium is well-reproduced.**

Molecular orbital given by DFT-LDA

$|0\rangle$ **1s bonding**

$|2\rangle$ **1s anti-bonding**

$|3\rangle$ $|3'\rangle$

$|2\rangle$

$|1\rangle$

$|0\rangle$

$|1\rangle$

$|3\rangle$

Dependence on inter-atomic distance

$|3\rangle \; |3'\rangle$

$|2\rangle$

$|1\rangle$

$|0\rangle$

$|1s,1\rangle$ & $|1s,2\rangle$ may be constructed.

Non-bonding Hitler-London picture

a

Is there any concise method to describe this dependence?

Correlation effects

- In DFT-LDA, the Kohn-Sham orbital is determined with correlation effects included in the effective potential.

- The correlation effect can behave differently depending on atomic configuration.

- A key is,

$$E_{xc}[n] = \frac{e^2}{2} \int d\mathbf{r}\, n(\mathbf{r}) \int d\mathbf{r}' \frac{1}{|\mathbf{r} - \mathbf{r}'|} n_{xc}(\mathbf{r}, \mathbf{r}')$$
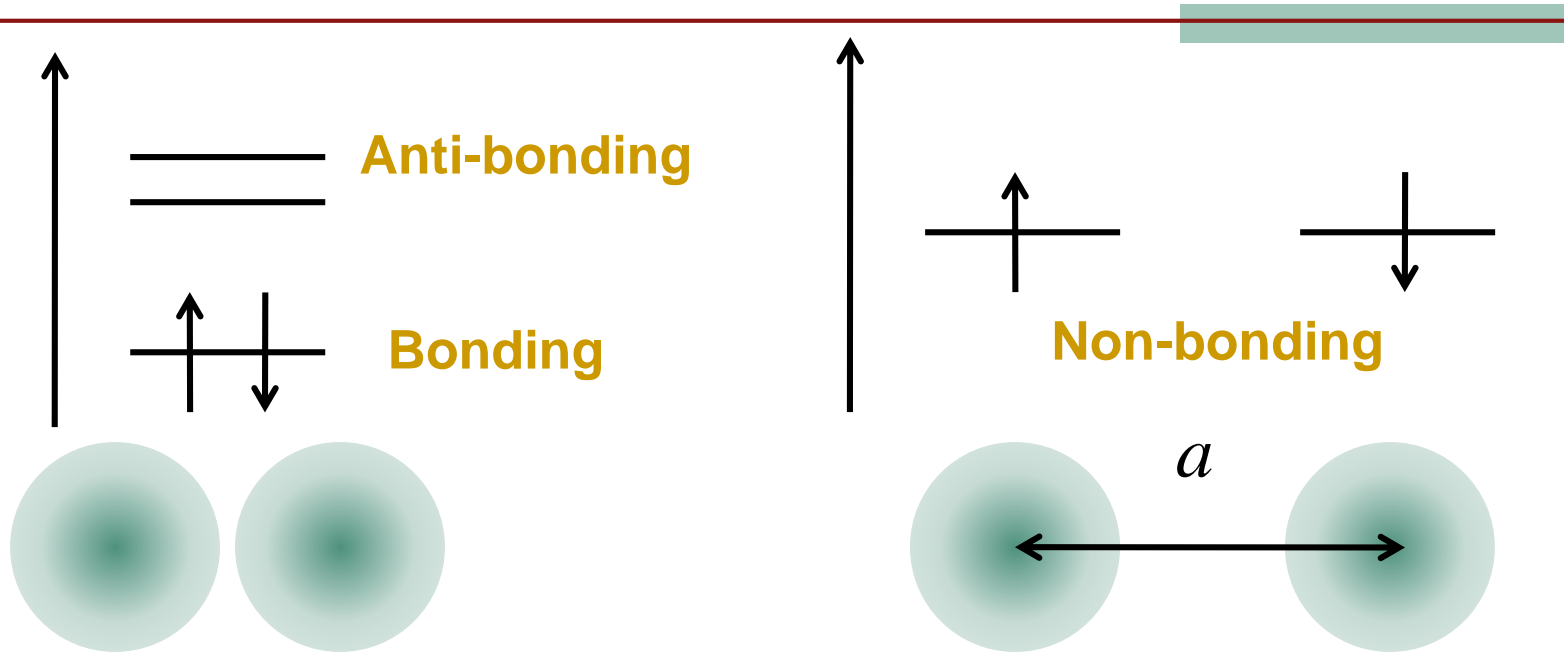
Charge fluctuation!

$$n_{xc}(\mathbf{r}, \mathbf{r}') = n(\mathbf{r}') \int_0^1 [g_n(\mathbf{r}, \mathbf{r}', \lambda) - \delta(\mathbf{r} - \mathbf{r}')] d\lambda$$

$$= \int_0^1 [\langle \Psi_\lambda | : (\hat{n}(\mathbf{r}) - n(\mathbf{r}))(\hat{n}(\mathbf{r}') - n(\mathbf{r}')) : | \Psi_\lambda \rangle / n(\mathbf{r})] d\lambda.$$
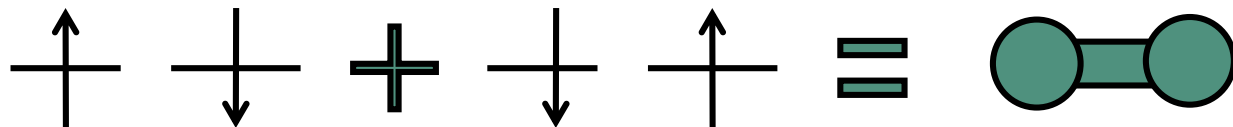
H$_2$
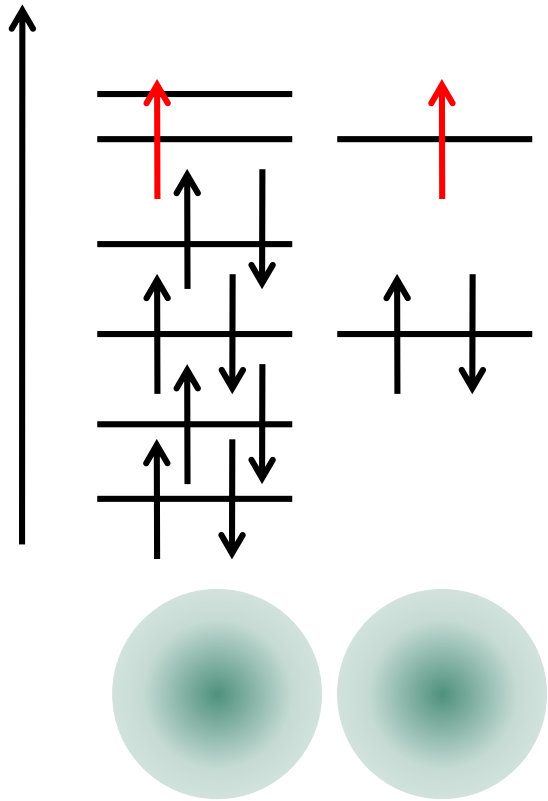
Anti-bonding

Bonding

Non-bonding

a

This is a weak correlation limit! **This is a strong correlation limit!**

In real nature & in CI, we have

Quantum mechanical superposition

O_2

This picture is not true!

2s^2 2p^4

a

S=1 ground state!
This is a strong correlation limit!
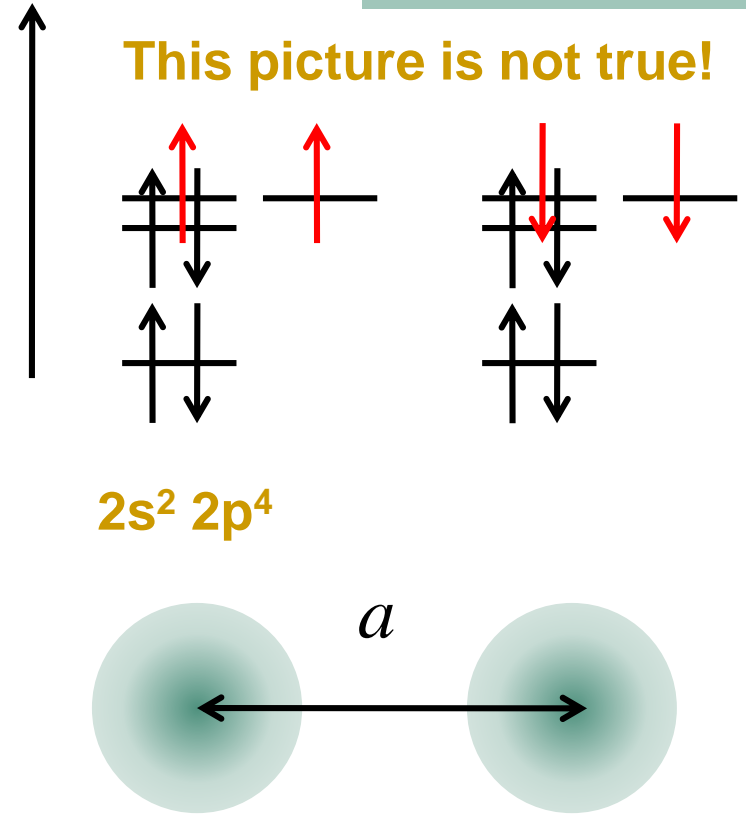
S=1 ground state!
This is a strong correlation limit!

How the multi-reference density functional theory acts?

- The Kohn-Sham single-particle description may be used, but may not always be reliable.

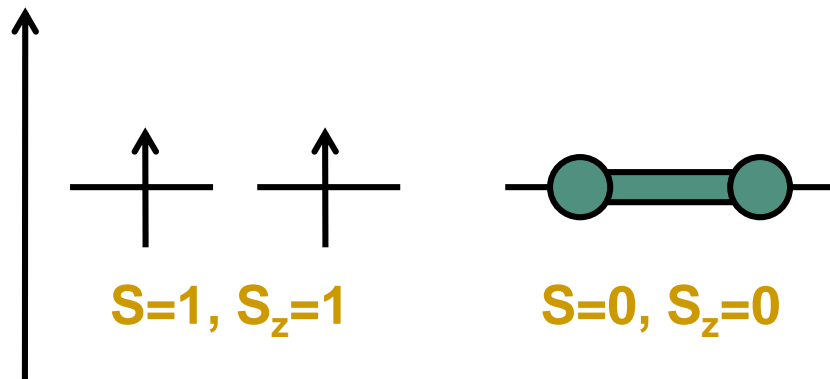- So, if a simplified description allowing a multi-Slater determinant is introduced, we may use the method as another starting point.

- A simplest case is a DFT-LDA solution for a degenerate ground state.

$S=1, S_z=1$ $S=0, S_z=0$

Determination of GS may be possible by looking at
- **charge-charge correlation or**
- **variational energy.**

Two physical quantities specifying correlation effects

- Density-density correlation : FRM

$$\left\langle \underline{n}_i^2 \right\rangle = \left\langle \left(n_{i\uparrow} + n_{i\downarrow} - \left\langle n_{i\uparrow} \right\rangle - \left\langle n_{i\downarrow} \right\rangle \right)^2 \right\rangle$$

- Variational energy : DFVT

$$E_0 \leq \min_{X_i, \varepsilon_i, g_i} \left\{ \min_{\Psi} \bar{G}_{X_i, \varepsilon_i, g_i}[\Psi] + \Delta \bar{E}_{X_i, \varepsilon_i, g_i}[\Psi] \right\},$$

$$\bar{G}_{X_i, \varepsilon_i, g_i}[\Psi] = \langle \Psi | \hat{T} + \hat{V}_{X_i} | \Psi \rangle + \frac{e^2}{2} \int \mathrm{d}^3 r \, \mathrm{d}^3 r' \frac{n_\Psi(\mathbf{r}) n_\Psi(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

$$+ E_{\varepsilon_i}[n_\Psi] + E_{g_i}[\Psi] + \int \mathrm{d}^3 r \, v_{\mathrm{ext}}(\mathbf{r}) n_\Psi(\mathbf{r}).$$

$$\Delta \bar{E}_{X_i, \varepsilon_i, g_i}[\Psi] = \frac{e^2}{2} \int \mathrm{d}^3 r \, \mathrm{d}^3 r' \frac{1}{|\mathbf{r} - \mathbf{r}'|}$$

$$\times \langle \Psi | : (\hat{n}(\mathbf{r}) - n_\Psi(\mathbf{r}))(\hat{n}(\mathbf{r}') - n_\Psi(\mathbf{r}')) : | \Psi \rangle$$

$$- E_{\varepsilon_i}[n_\Psi] - E_{g_i}[\Psi] - \langle \Psi | \hat{V}_{X_i} | \Psi \rangle.$$

Fluctuation reference method

- To have a simplified description based on DFT, we need information given by another accurate electronic structure calculation.

  - Diffusion Monte-Carlo method for electron gas ⇒ LDA, GGA

  - Complete-Active-Space Configuration-Interaction method ⇒ FRM to determine U.

  FRM : Fluctuation reference method

Two Hydrogen systems

We define ϕ_i by

$$\phi_1 = \frac{1}{\sqrt{2}}(\chi_1 + \chi_2),$$
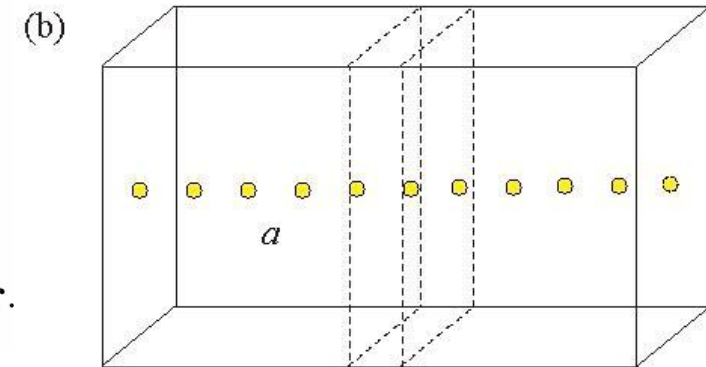
$$\phi_2 = \frac{1}{\sqrt{2}}(\chi_1 - \chi_2),$$

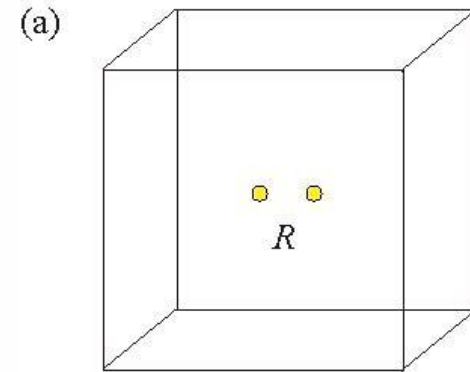for the molecule and the Wannier state

$$\phi_i = \frac{1}{\sqrt{N}}\sum_{k=1}^{N}\exp\left(i\frac{2\pi}{Na}kx_i\right)\chi_{1,k},$$

for a chain with N atoms.

$$t_{ij} = \int \phi_i^*(\mathbf{r})\left\{-\frac{\hbar^2}{2m}\Delta_{\mathbf{r}} + v_{\text{eff}}(\mathbf{r})\right\}\phi_j(\mathbf{r})d\mathbf{r}.$$
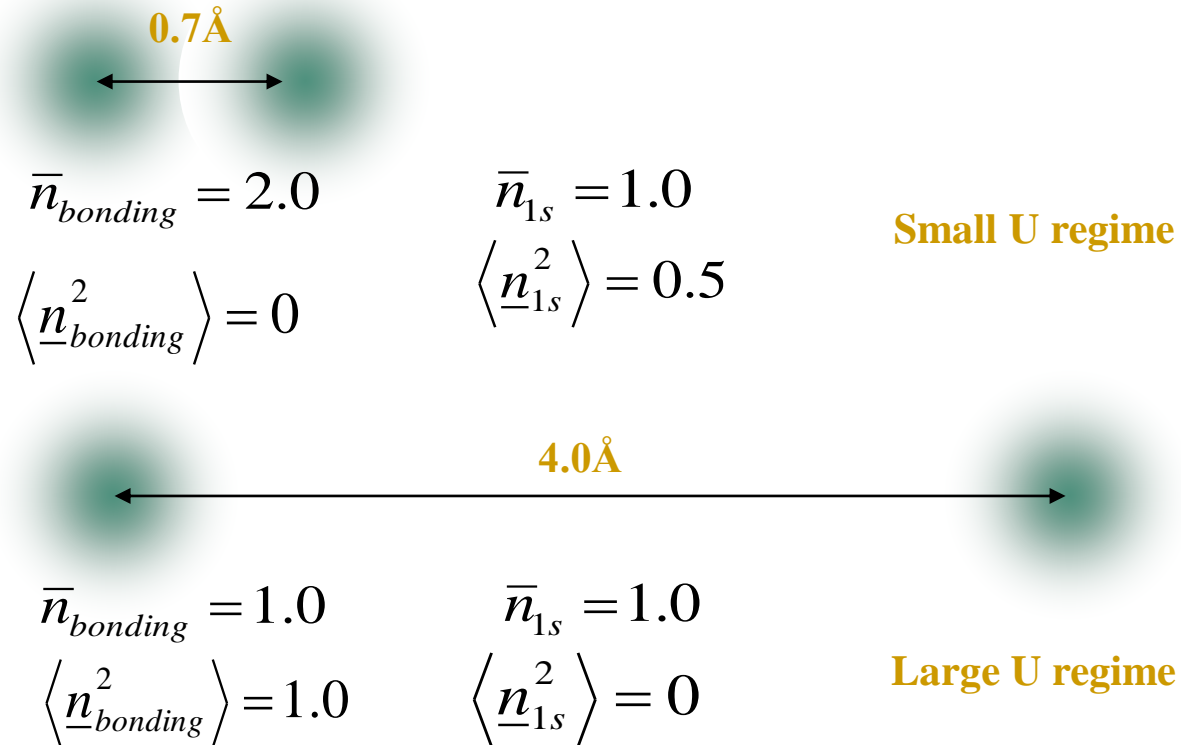
H$_2$ Molecule
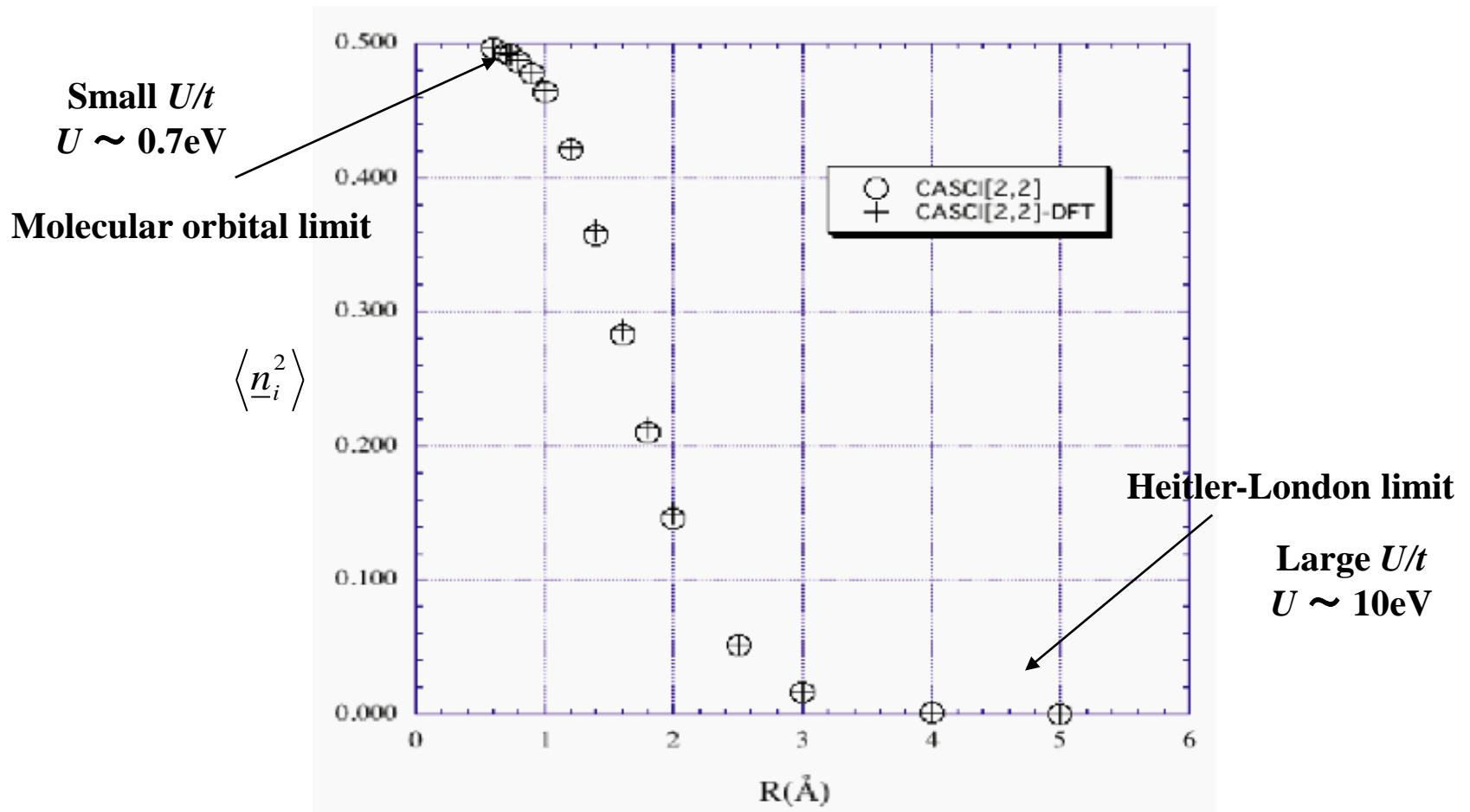
(a)

(b)

Hydrogen array

Local fluctuation on ϕ_{1s} of H_2

$\left\langle \underline{n}_i^2 \right\rangle = \left\langle \left(n_{i\uparrow} + n_{i\downarrow} - \left\langle n_{i\uparrow} \right\rangle - \left\langle n_{i\downarrow} \right\rangle \right)^2 \right\rangle$ **as a function of the inter atomic distance**

0.7Å

$\bar{n}_{bonding} = 2.0$ $\bar{n}_{1s} = 1.0$

$\left\langle \underline{n}_{bonding}^2 \right\rangle = 0$ $\left\langle \underline{n}_{1s}^2 \right\rangle = 0.5$

Small U regime

4.0Å

$\bar{n}_{bonding} = 1.0$ $\bar{n}_{1s} = 1.0$

$\left\langle \underline{n}_{bonding}^2 \right\rangle = 1.0$ $\left\langle \underline{n}_{1s}^2 \right\rangle = 0$

Large U regime

We are able to find relevant orbitals with Coulomb suppression by
1. Finding an orbital (or a set of orbitals) on which fluctuation becomes 1.
2. Using a unitary transformation to have localized orbitals.

Orbital fluctuation in CI calculation

CAS-CI calculation for H_2 by S. Yamanaka.

Small U/t
$U \sim 0.7\text{eV}$

Molecular orbital limit

$\left\langle \underline{n}_i^2 \right\rangle$

Heitler-London limit

Large U/t
$U \sim 10\text{eV}$

In the whole range, mean occupation of 1s is 1.

A test calculation of Hydrogen systems

- $E_{xc}[n]$ is given by LDA of PW91.
- Plane-wave expansion with the Troullier-Martins soft-pseudopotential is introduced for H.
- The energy cutoff for the plane-wave is 40Ry.
- The solver of single-particle part of the extended Kohn-Sham equation is given by the "opt" code.
- Effective many-body problem is solved by the exact diagonalization with e.g. the Lanczos method.

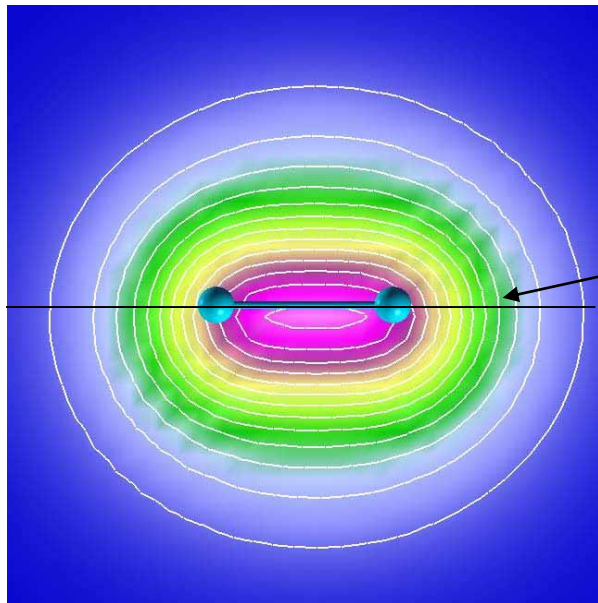# The value of U for H_2 determined by FRM

Reference calculation: CASCI method for H_2
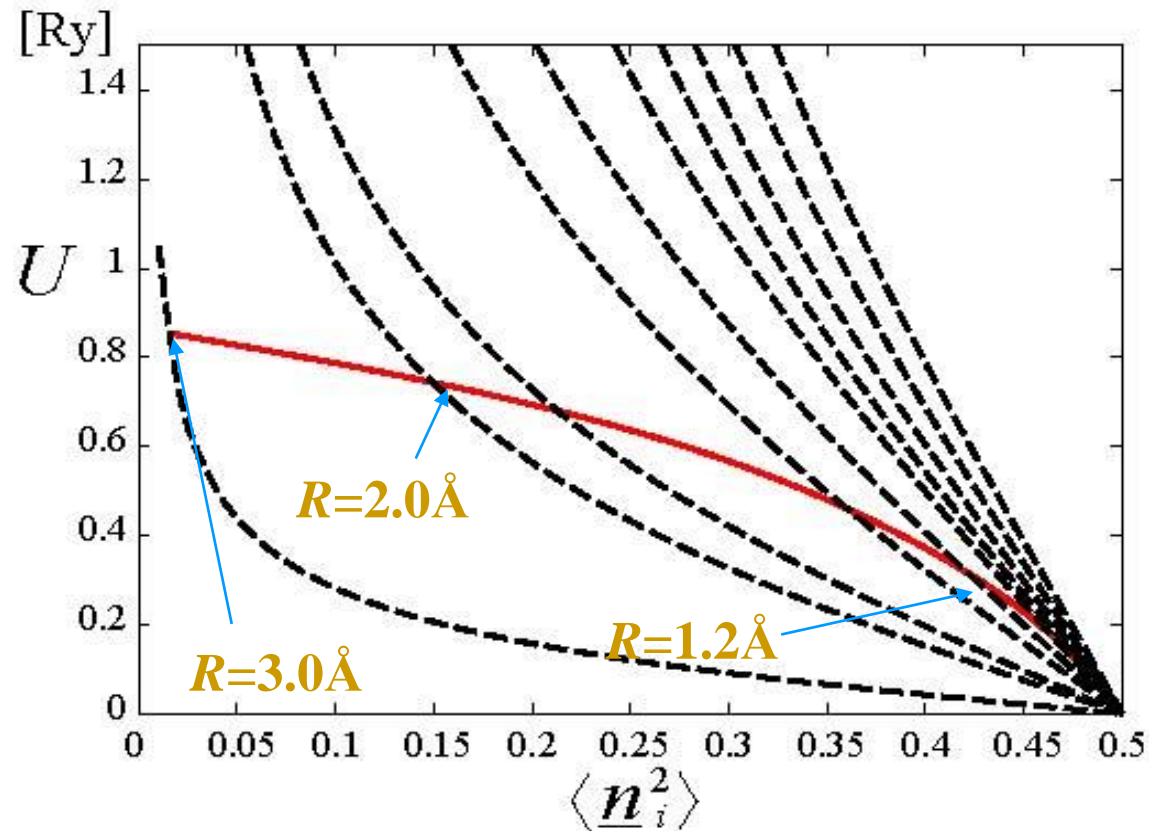
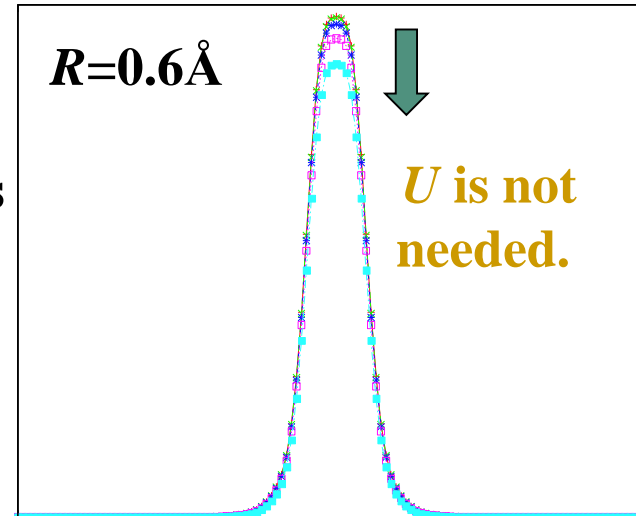
The 1st method to have U by the fluctuation reference method.

K. Kusakabe, et al. J. Phys.: Condens. Matter (2007).

Charge density of H_2 given by MR-DFT

Weak correlation regime

This section is Shown in the right panels.

R=0.6Å

U is not needed.

Strong correlation regime

R=2.0Å

U is need.
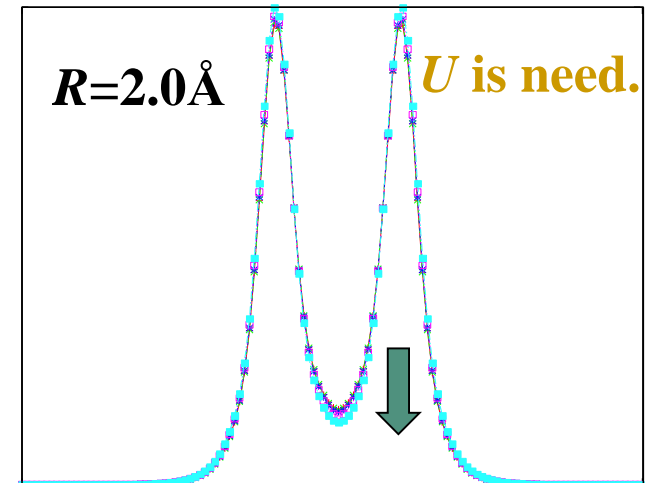
To reproduce bonding charge in correlated electron regime, we need to overcome difficulty in the single-reference description. It is achieved by using the multi-reference description as known in the Hybrid-DFT.
(Cf. S. Yamanaka, et al.)

Variational method

DFVT for MR-DFT

$$E_0 \leq \min_{X_i, \varepsilon_i, g_i} \left\{ \min_{\Psi} \bar{G}_{X_i, \varepsilon_i, g_i}[\Psi] + \Delta \bar{E}_{X_i, \varepsilon_i, g_i}[\Psi] \right\},$$

$$\bar{G}_{X_i, \varepsilon_i, g_i}[\Psi] = \langle \Psi | \hat{T} + \hat{V}_{X_i} | \Psi \rangle + \frac{e^2}{2} \int d^3 r \, d^3 r' \frac{n_\Psi(\mathbf{r}) n_\Psi(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} + E_{\varepsilon_i}[n_\Psi] + E_{g_i}[\Psi] + \int d^3 r \, v_{\text{ext}}(\mathbf{r}) n_\Psi(\mathbf{r}).$$

$$\Delta \bar{E}_{X_i, \varepsilon_i, g_i}[\Psi] = \frac{e^2}{2} \int d^3 r \, d^3 r' \frac{1}{|\mathbf{r} - \mathbf{r}'|} \times \langle \Psi | : (\hat{n}(\mathbf{r}) - n_\Psi(\mathbf{r}))(\hat{n}(\mathbf{r}') - n_\Psi(\mathbf{r}')) : | \Psi \rangle - E_{\varepsilon_i}[n_\Psi] - E_{g_i}[\Psi] - \langle \Psi | \hat{V}_{X_i} | \Psi \rangle.$$

DFT model formation

$$E_0 \leq \min_{X_i, \varepsilon_i, g_i} \left\{ \min_{\Psi} \bar{G}_{X_i, \varepsilon_i, g_i}[\Psi] + \Delta E_{X_i, \varepsilon_i, g_i}[\Psi] \right\},$$

$$\Delta E_{X_i, \varepsilon_i, g_i}[\Psi] = F[n_\Psi] - \frac{e^2}{2} \int d^3 r \, d^3 r' \frac{n_\Psi(\mathbf{r}) n_\Psi(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} - E_{\varepsilon_i}[n_\Psi] - E_{g_i}[\Psi] - \langle \Psi | \hat{T} + \hat{V}_{X_i} | \Psi \rangle.$$

PBE0+GW, PBE0+U will be refined.

It is possible to evaluate $\Delta \bar{E}[\Psi]$.

Note: If DFVT is used in a space of single Slater determinant, the Hartree-Fock approx. is given.
This is natural, but not so trivial.
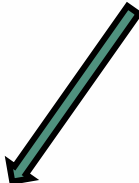Thanks to Dr. Maruyama & Dr. Friedlich.

To understand MR-DFT

$$E_0 \leq \min_{X_i, \varepsilon_i, g_i} \left\{ \min_{\Psi} \bar{G}_{X_i, \varepsilon_i, g_i}[\Psi] + \Delta E_{X_i, \varepsilon_i, g_i}[\Psi] \right\},$$

$$\Delta E_{X_i, \varepsilon_i, g_i}[\Psi] = F[n_\Psi] - \frac{e^2}{2} \int d^3 r \, d^3 r' \frac{n_\Psi(\mathbf{r}) n_\Psi(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

$$- E_{\varepsilon_i}[n_\Psi] - E_{g_i}[\Psi] - \langle \Psi | \hat{T} + \hat{V}_{X_i} | \Psi \rangle.$$

This functional contains the universal energy-density functional F[n]

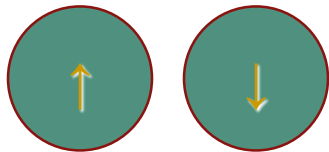**This functional should be evaluated using a multi-Slater description.**

$$\Delta \bar{E}_{X_i, \varepsilon_i, g_i}[\Psi] = \frac{e^2}{2} \int d^3 r \, d^3 r' \frac{1}{|\mathbf{r} - \mathbf{r}'|}$$

$$\times \langle \Psi | : (\hat{n}(\mathbf{r}) - n_\Psi(\mathbf{r}))(\hat{n}(\mathbf{r}') - n_\Psi(\mathbf{r}')) : | \Psi \rangle$$

$$- E_{\varepsilon_i}[n_\Psi] - E_{g_i}[\Psi] - \langle \Psi | \hat{V}_{X_i} | \Psi \rangle.$$

Heisenberg exchange:
Introduction to the localized spin model

- Heisenberg and Dirac considered quantum mechanical origin of magnetic interaction

$$-J\, S_i \cdot S_j$$

J>0 : Parallel spin configuration
Ferromagnetic
J<0 : Anti-parallel spin configuration
Antiferromagnetic

whose order of magnitude can be of the order of the Coulomb interaction.

Where does J come from?

Idea of restriction of the phase space

■ For low-energy states, we may consider only states in a restricted phase space.

E

$\vdots$

4s,4p,4d,4f
3s,3p,3d

2s,2p

1s
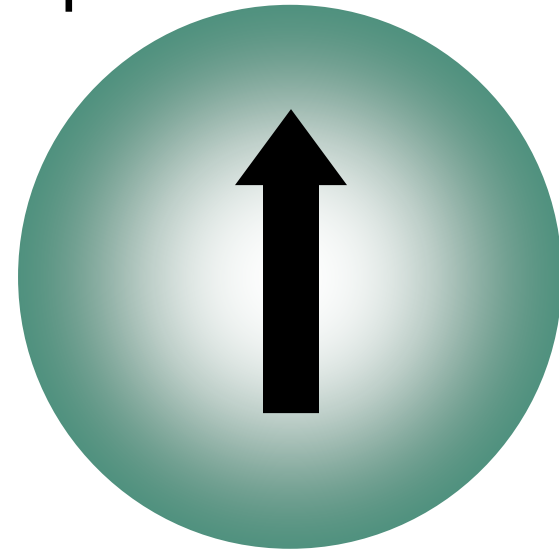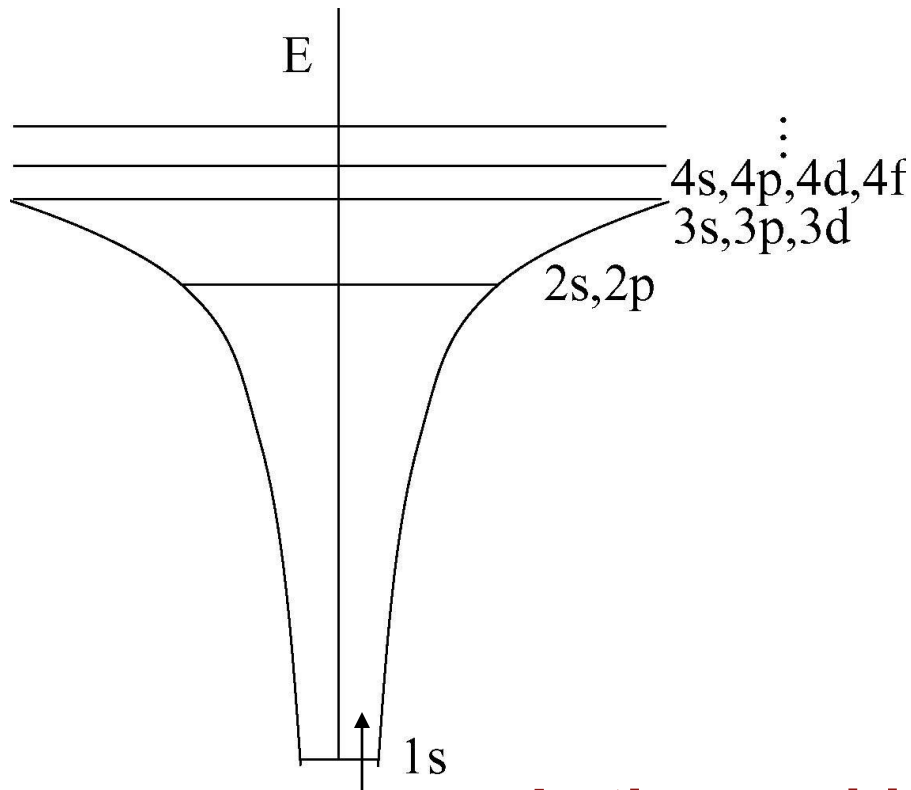
The 1s orbital of the Hydrogen atom

Let's consider the phase space spanned by spin states $|1s, \uparrow \rangle$, $|1s, \downarrow \rangle$.

Atomic orbitals and molecular orbitals

- ### Atomic orbitals
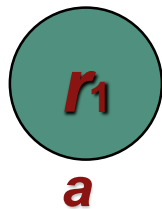
$$\phi_a(r)\xi_\uparrow(1)$$

$$c_{a\uparrow}^\dagger|0\rangle$$

- ### Molecular orbitals (in LCAO)

$$\frac{1}{\sqrt{2}}\big(\phi_a(r)+\phi_b(r)\big)\xi_\uparrow(1)$$

$$\frac{1}{\sqrt{2}}\big(c_{a\uparrow}^\dagger+c_{b\uparrow}^\dagger\big)|0\rangle$$

21

The Heitler-London theory of H_2

Two electrons repel each other and each orbital ϕ_u ($u = a$, b) is occupied by an electron **A picture in the strong correlation limit!!**

$$\Psi_1(\mathbf{r}_1, \mathbf{r}_2) = \phi_a(\mathbf{r}_1)\phi_b(\mathbf{r}_2) \qquad \Psi_2(\mathbf{r}_1, \mathbf{r}_2) = \phi_a(\mathbf{r}_2)\phi_b(\mathbf{r}_1)$$

r_1 r_2 $\Longleftrightarrow$ r_2 r_1

a   *b*   *a*   *b*
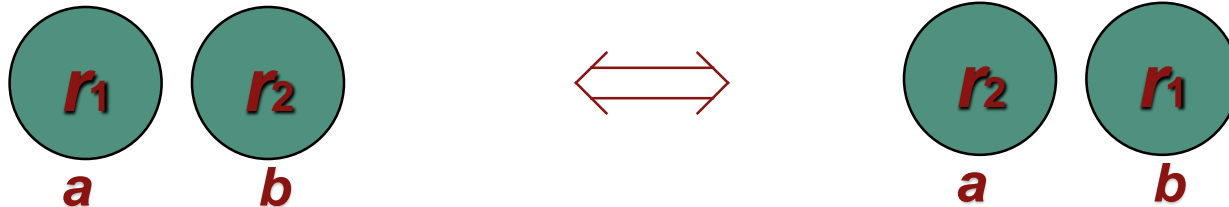
Note that there is no doubly occupied 1s orbital in this theory.

First key: Quantum states are linear combinations of Ψ_1 and Ψ_2 due to Pauli's principle.
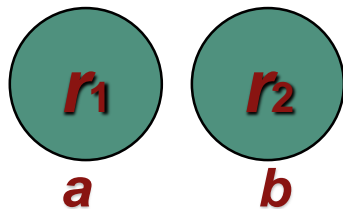
(3)   **(4)**

Singlet: $\Psi_+ = c_+(\Psi_1 + \Psi_2)$.   **Triplet:** $\Psi_- = c_-(\Psi_1 - \Psi_2)$.
$(\mathbf{S}_1 + \mathbf{S}_2)^2 = 0$. $(\mathbf{S}_1 + \mathbf{S}_2)^2 = 2$.
$E = E_+$. $E = E_-$.

22

Energy of spin states

Using following values, $E_\pm$ are obtained.

The orbital energy e_0 given by $\mathcal{H}\phi_u(\mathbf{r}) = e_0\phi_u(\mathbf{r})$. ($u = a, b$ and $\mathcal{H}$ is the Hamiltonian of an atom.)

An overlap integral $S = \int \phi_a^*(\mathbf{r})\phi_b(\mathbf{r})d\mathbf{r}$.

A Coulomb integral $U = \int \Psi_1^*(\mathbf{r}_1, \mathbf{r}_2)\left(\dfrac{e^2}{R_{ab}} + \dfrac{e^2}{r_{12}} - \dfrac{e^2}{r_{1b}} - \dfrac{e^2}{r_{2b}}\right)\Psi_1(\mathbf{r}_1, \mathbf{r}_2)d\mathbf{r}_1 d\mathbf{r}_2$.

An exchange integral $I = \int \Psi_1^*(\mathbf{r}_1, \mathbf{r}_2)\left(\dfrac{e^2}{R_{ab}} + \dfrac{e^2}{r_{12}} - \dfrac{e^2}{r_{1b}} - \dfrac{e^2}{r_{2b}}\right)\Psi_2(\mathbf{r}_1, \mathbf{r}_2)d\mathbf{r}_1 d\mathbf{r}_2$.

We have

$$E_+ = 2e_0 + \frac{U + I}{1 + S^2} , \qquad\qquad E_- = 2e_0 + \frac{U - I}{1 - S^2} .$$

Evaluation of J: Direct exchange

The triplet-singlet separation is evaluated as,

$$\Delta E = E_- - E_+ = 2\frac{US^2 - I}{1 - S^4} \ .$$

Second key: Energy separation between the magnetic state and the non-magnetic state can be as large as the Coulomb interaction.

The Heisenberg interaction is redescribed as,

$$H_{Heis} = -J_{12}\mathbf{S}_1 \cdot \mathbf{S}_2 = -J_{12}\left[\frac{(\mathbf{S}_1 + \mathbf{S}_2)^2}{2} - \frac{3}{4}\right] \ .$$

Energy of each state becomes $E_- = -\frac{1}{4}J_{12}$ for the triplets and $E_+ = \frac{3}{4}J_{12}$ for the singlet, which results $\Delta E = -J_{12}$ and thus,

$$J_{12} = -2\frac{US^2 - I}{1 - S^4} \ .$$

If $S = 0$, the interaction is ferromagnetic.

Ferromagnetic direct exchange I.

Consider atomic orbitals $\phi_u(\mathbf{r})$ which is real.

"The exchange integral I is positive, *i.e.* it is ferromagnetic, when $S = 0$."

Proof: The definition of I is,

$$I = \int \phi_a^*(\mathbf{r}_1)\phi_b^*(\mathbf{r}_2)\left(\frac{e^2}{R_{ab}} + \frac{e^2}{r_{12}} - \frac{e^2}{r_{1b}} - \frac{e^2}{r_{2b}}\right)\phi_a(\mathbf{r}_2)\phi_b(\mathbf{r}_1)d\mathbf{r}_1 d\mathbf{r}_2 .$$

Let S=0, then we have,

$$I = \int \frac{e^2}{r_{12}}\phi_a^*(\mathbf{r}_1)\phi_b(\mathbf{r}_1)\phi_a(\mathbf{r}_2)\phi_b^*(\mathbf{r}_2)d\mathbf{r}_1 d\mathbf{r}_2 .$$

Ferromagnetic direct exchange II.

- Exchange integral of two orthogonal orbitals with $S=0$ becomes ferromagnetic.

A proof: Consider the exchange integral when $S=0$.

$$I = \int d\mathbf{r}_1 d\mathbf{r}_2 \frac{e^2}{|\mathbf{r}_1 - \mathbf{r}_2|} \phi_a^*(\mathbf{r}_1)\phi_b(\mathbf{r}_1)\phi_b^*(\mathbf{r}_2)\phi_a(\mathbf{r}_2)$$

Note that the expression below.

$$\phi_a^*(\mathbf{r}_1)\phi_b(\mathbf{r}_1) = \mathrm{Re}\left(\phi_a^*(\mathbf{r}_1)\phi_b(\mathbf{r}_1)\right) + i\,\mathrm{Im}\left(\phi_a^*(\mathbf{r}_1)\phi_b(\mathbf{r}_1)\right)$$

$$I = \int d\mathbf{r}_1 d\mathbf{r}_2 \frac{e^2}{|\mathbf{r}_1 - \mathbf{r}_2|} \phi_a^*(\mathbf{r}_1)\phi_b(\mathbf{r}_1)\left(\phi_a^*(\mathbf{r}_2)\phi_b(\mathbf{r}_2)\right)^*$$

$$= \int d\mathbf{r}_1 d\mathbf{r}_2 \frac{e^2}{|\mathbf{r}_1 - \mathbf{r}_2|} \mathrm{Re}\left(\phi_a^*(\mathbf{r}_1)\phi_b(\mathbf{r}_1)\right) \times \mathrm{Re}\left(\phi_a^*(\mathbf{r}_2)\phi_b(\mathbf{r}_2)\right)$$

$$+ \int d\mathbf{r}_1 d\mathbf{r}_2 \frac{e^2}{|\mathbf{r}_1 - \mathbf{r}_2|} \mathrm{Im}\left(\phi_a^*(\mathbf{r}_1)\phi_b(\mathbf{r}_1)\right) \times \mathrm{Im}\left(\phi_a^*(\mathbf{r}_2)\phi_b(\mathbf{r}_2)\right)$$

To show the positivity of I, it is enough to consider an integral of real functions,

$$I = \int d\mathbf{r}_1 d\mathbf{r}_2 \frac{e^2}{|\mathbf{r}_1 - \mathbf{r}_2|} \Phi(\mathbf{r}_1)\Phi(\mathbf{r}_2)$$

Ferromagnetic direct exchange III.

Substitute next expressions into I.

$$\frac{e^2}{r} = \int d\mathbf{q} \frac{4\pi e^2}{q^2} \exp(i\mathbf{q} \cdot \mathbf{r}) \ ,$$

$$\mathbf{\Phi}(\mathbf{r}) = \int d\mathbf{q} v(\mathbf{q}) \exp(i\mathbf{q} \cdot \mathbf{r}) \ .$$

Since $\mathbf{\Phi}$ are real, $v^*(\mathbf{q}) = v(\text{-q})$. Then, we can show,

$$
\begin{aligned}
I &= \int d\mathbf{r}_1 d\mathbf{r}_2 \int d\mathbf{q} \frac{4\pi e^2}{q^2} \exp(i\mathbf{q} \cdot (\mathbf{r}_1 - \mathbf{r}_2)) \\
&\quad \int d\mathbf{q}_1 v(\mathbf{q}_1) \exp(i\mathbf{q}_1 \cdot \mathbf{r}_1) \int d\mathbf{q}_2 v(\mathbf{q}_2) \exp(i\mathbf{q}_2 \cdot \mathbf{r}_2) \\
&= \int d\mathbf{q} \frac{4\pi e^2}{q^2} v(-\mathbf{q}) v(\mathbf{q}) \\
&= \int d\mathbf{q} \frac{4\pi e^2}{q^2} v^*(\mathbf{q}) v(\mathbf{q}) > 0 \ .
\end{aligned}
$$

Slater determinants for a system with two electrons I.

1. Triplet states

$$\langle \mathbf{r}_1, 1; \mathbf{r}_2, 2 | S_{tot} = 1, S_z = 1 \rangle = \langle \mathbf{r}_1, 1; \mathbf{r}_2, 2 | \phi_a \uparrow, \phi_b \uparrow \rangle$$

$$= \frac{1}{\sqrt{2}} \begin{vmatrix} \phi_a(\mathbf{r}_1)\xi_\uparrow(1) & \phi_a(\mathbf{r}_2)\xi_\uparrow(2) \\ \phi_b(\mathbf{r}_1)\xi_\uparrow(1) & \phi_b(\mathbf{r}_2)\xi_\uparrow(2) \end{vmatrix} = \frac{1}{\sqrt{2}} \begin{vmatrix} \phi_a(\mathbf{r}_1) & \phi_a(\mathbf{r}_2) \\ \phi_b(\mathbf{r}_1) & \phi_b(\mathbf{r}_2) \end{vmatrix} \xi_\uparrow(1)\xi_\uparrow(2) \quad (1)$$

$$\langle \mathbf{r}_1, 1; \mathbf{r}_2, 2 | S_{tot} = 1, S_z = 0 \rangle = \langle \mathbf{r}_1, 1; \mathbf{r}_2, 2 | S_{tot}^- | \phi_a \uparrow, \phi_b \uparrow \rangle$$

$$= \frac{1}{\sqrt{2}} \begin{vmatrix} \phi_a(\mathbf{r}_1) & \phi_a(\mathbf{r}_2) \\ \phi_b(\mathbf{r}_1) & \phi_b(\mathbf{r}_2) \end{vmatrix} \frac{1}{\sqrt{2}} \left(\xi_\uparrow(1)\xi_\downarrow(2) + \xi_\downarrow(1)\xi_\uparrow(2) \right)$$

$$= \frac{1}{2} \{ (\phi_a(\mathbf{r}_1)\xi_\uparrow(1)\phi_b(\mathbf{r}_2)\xi_\downarrow(2) - \phi_b(\mathbf{r}_1)\xi_\uparrow(1)\phi_a(\mathbf{r}_2)\xi_\downarrow(2))$$

$$+ (\phi_a(\mathbf{r}_1)\xi_\downarrow(1)\phi_b(\mathbf{r}_2)\xi_\uparrow(2) - \phi_b(\mathbf{r}_1)\xi_\downarrow(1)\phi_a(\mathbf{r}_2)\xi_\uparrow(2)) \}$$

$$= \frac{1}{\sqrt{2}} \left\{ \frac{1}{\sqrt{2}} \begin{vmatrix} \phi_a(\mathbf{r}_1)\xi_\uparrow(1) & \phi_a(\mathbf{r}_2)\xi_\uparrow(2) \\ \phi_b(\mathbf{r}_1)\xi_\downarrow(1) & \phi_b(\mathbf{r}_2)\xi_\downarrow(2) \end{vmatrix} + \frac{1}{\sqrt{2}} \begin{vmatrix} \phi_a(\mathbf{r}_1)\xi_\downarrow(1) & \phi_a(\mathbf{r}_2)\xi_\downarrow(2) \\ \phi_b(\mathbf{r}_1)\xi_\uparrow(1) & \phi_b(\mathbf{r}_2)\xi_\uparrow(2) \end{vmatrix} \right\} \quad (2)$$

Slater determinants for a system with two electrons II.

Thus,

$$\langle \mathbf{r}_1, 1; \mathbf{r}_2, 2 | S_{tot} = 1, S_z = 1, 0, -1 \rangle$$

$$= \frac{1}{\sqrt{2}} \begin{vmatrix} \phi_a(\mathbf{r}_1) & \phi_a(\mathbf{r}_2) \\ \phi_b(\mathbf{r}_1) & \phi_b(\mathbf{r}_2) \end{vmatrix} \begin{cases} \xi_\uparrow(1)\xi_\uparrow(2) \\ \frac{1}{\sqrt{2}}(\xi_\uparrow(1)\xi_\downarrow(2) + \xi_\downarrow(1)\xi_\uparrow(2)) \\ \xi_\downarrow(1)\xi_\downarrow(2) \end{cases} \quad (3)$$

2. Singlet states

We have another state which is orthogonal to $|S_{tot} = 1, S_z = 0\rangle$.

Multi Slater determinants

$$\langle \mathbf{r}_1, 1; \mathbf{r}_2, 2 | S_{tot} = 0, S_z = 0 \rangle$$

$$= \frac{1}{\sqrt{2}} \left\{ \frac{1}{\sqrt{2}} \begin{vmatrix} \phi_a(\mathbf{r}_1)\xi_\uparrow(1) & \phi_a(\mathbf{r}_2)\xi_\uparrow(2) \\ \phi_b(\mathbf{r}_1)\xi_\downarrow(1) & \phi_b(\mathbf{r}_2)\xi_\downarrow(2) \end{vmatrix} - \frac{1}{\sqrt{2}} \begin{vmatrix} \phi_a(\mathbf{r}_1)\xi_\downarrow(1) & \phi_a(\mathbf{r}_2)\xi_\downarrow(2) \\ \phi_b(\mathbf{r}_1)\xi_\uparrow(1) & \phi_b(\mathbf{r}_2)\xi_\uparrow(2) \end{vmatrix} \right\}$$

$$= \frac{1}{2} \left\{ (\phi_a(\mathbf{r}_1)\xi_\uparrow(1)\phi_b(\mathbf{r}_2)\xi_\downarrow(2) - \phi_b(\mathbf{r}_1)\xi_\uparrow(1)\phi_a(\mathbf{r}_2)\xi_\downarrow(2)) \right.$$

$$\left. - (\phi_a(\mathbf{r}_1)\xi_\downarrow(1)\phi_b(\mathbf{r}_2)\xi_\uparrow(2) - \phi_b(\mathbf{r}_1)\xi_\downarrow(1)\phi_a(\mathbf{r}_2)\xi_\uparrow(2)) \right\}$$

$$= \frac{1}{\sqrt{2}} (\phi_a(\mathbf{r}_1)\phi_b(\mathbf{r}_2) + \phi_b(\mathbf{r}_1)\phi_a(\mathbf{r}_2)) \frac{1}{\sqrt{2}} \begin{vmatrix} \xi_\uparrow(1) & \xi_\uparrow(2) \\ \xi_\downarrow(1) & \xi_\downarrow(2) \end{vmatrix} \quad (4)$$

Spin states in the second quantization

In the second-quantization description,

$$\langle \mathbf{r}_1, 1; \mathbf{r}_2, 2 | S_{tot} = 1, S_z = 1 \rangle = \frac{1}{\sqrt{2}} \begin{vmatrix} \phi_a(\mathbf{r}_1)\xi_\uparrow(1) & \phi_a(\mathbf{r}_2)\xi_\uparrow(2) \\ \phi_b(\mathbf{r}_1)\xi_\uparrow(1) & \phi_b(\mathbf{r}_2)\xi_\uparrow(2) \end{vmatrix}$$

$$\longrightarrow \quad |S_{tot} = 1, S_z = 1 \rangle = c_{a,\uparrow}^\dagger c_{b,\uparrow}^\dagger |0\rangle$$

Similarly,

$$|S_{tot} = 1, S_z = 0 \rangle = \frac{1}{\sqrt{2}}(c_{a,\uparrow}^\dagger c_{b,\downarrow}^\dagger + c_{a,\downarrow}^\dagger c_{b,\uparrow}^\dagger)|0\rangle$$

$$|S_{tot} = 1, S_z = -1 \rangle = c_{a,\downarrow}^\dagger c_{b,\downarrow}^\dagger |0\rangle$$

$$|S_{tot} = 0, S_z = 0 \rangle = \frac{1}{\sqrt{2}}(c_{a,\uparrow}^\dagger c_{b,\downarrow}^\dagger - c_{a,\downarrow}^\dagger c_{b,\uparrow}^\dagger)|0\rangle$$

2-electrons in molecular orbitals

Let us consider two electrons in the bonding orbital.

$$\langle \mathbf{r}_1, 1; \mathbf{r}_2, 2 | \frac{1}{\sqrt{2}}(\phi_a + \phi_b) \uparrow, \frac{1}{\sqrt{2}}(\phi_a + \phi_b) \downarrow \rangle$$

$$= \frac{1}{2}(\phi_a(\mathbf{r}_1) + \phi_b(\mathbf{r}_1))(\phi_a(\mathbf{r}_2) + \phi_b(\mathbf{r}_2))\frac{1}{\sqrt{2}}\begin{vmatrix} \xi_\uparrow(1) & \xi_\uparrow(2) \\ \xi_\downarrow(1) & \xi_\downarrow(2) \end{vmatrix}$$

$$\longrightarrow | \frac{1}{\sqrt{2}}(\phi_a + \phi_b) \uparrow, \frac{1}{\sqrt{2}}(\phi_a + \phi_b) \downarrow \rangle$$

$$= \frac{1}{2}(c_{a,\uparrow}^\dagger + c_{b,\uparrow}^\dagger)(c_{a,\downarrow}^\dagger + c_{b,\downarrow}^\dagger)|0\rangle \qquad (5)$$

If we expand eq. (5), we obtain,

$$| \frac{1}{\sqrt{2}}(\phi_a + \phi_b) \uparrow, \frac{1}{\sqrt{2}}(\phi_a + \phi_b) \downarrow \rangle$$

$$= \frac{1}{\sqrt{2}}\left\{ \frac{1}{\sqrt{2}}(c_{a,\uparrow}^\dagger c_{b,\downarrow}^\dagger - c_{a,\downarrow}^\dagger c_{b,\uparrow}^\dagger)|0\rangle + \frac{1}{\sqrt{2}}(c_{a,\uparrow}^\dagger c_{a,\downarrow}^\dagger + c_{b,\uparrow}^\dagger c_{b,\downarrow}^\dagger)|0\rangle \right\} \qquad (6)$$

The Heitler-London wf.   **Ionic states!**

The Heitler-London state v.s. the molecular orbitals

We can consider a trial (variational) state in which ionic states are mixed with the Heitler-London wavefunction.

$$|S_{tot} = 0; \alpha\rangle$$

$$= \frac{1}{\sqrt{1+\alpha^2}} \left\{ \frac{1}{\sqrt{2}} (c_{a,\uparrow}^\dagger c_{b,\downarrow}^\dagger - c_{a,\downarrow}^\dagger c_{b,\uparrow}^\dagger)|0\rangle + \frac{\alpha}{\sqrt{2}} (c_{a,\uparrow}^\dagger c_{a,\downarrow}^\dagger + c_{b,\uparrow}^\dagger c_{b,\downarrow}^\dagger)|0\rangle \right\} \quad (7)$$

1. $\alpha = 0$: the Heiter-London wavefunction (Localized electron picture)

2. $\alpha = 1$: the doubly-occupied bonding state (Delocalized electron picture)

The true w.f. is close to a variational state with 0<α<1.

α is dependent on the inter-electron interaction or choice of the model (interaction parameters etc.).

The Hubbard model : I.

1. Consider a system with atomic sites. Each site is assumed to have an orbital for the conduction electrons
2. Electrons can hop between neighboring sites by a transfer integral *t*.
3. Two electrons coming across at a site feels a repulsive interaction *U*.
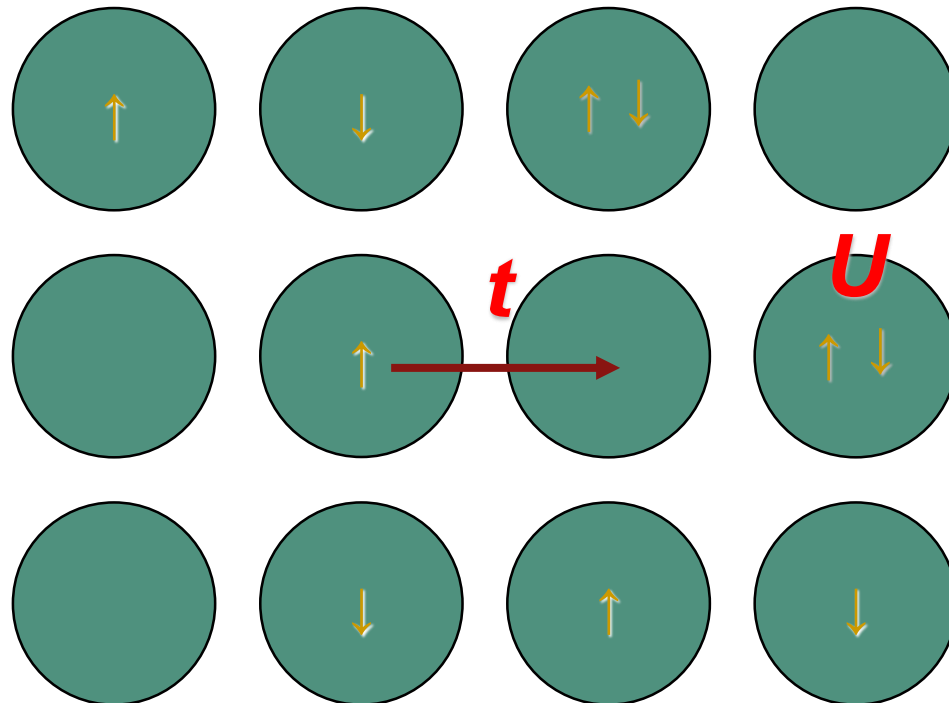
The Hubbard model : II.

1) The Bloch function

For crystals, the single particle wave function $\phi_{i,\sigma}(\mathbf{r})$ is given in a Bloch form as,

$$\phi_{n,\mathbf{k},\sigma}(\mathbf{r}) = e^{i\mathbf{k}\mathbf{r}} u_{n,\sigma}(\mathbf{r}) \,,$$

where n and $\mathbf{k}$ are the band index and the wave vector, respectively. $u_{n,\sigma}(\mathbf{r})$ is a function with periodicity of the lattice. $\phi_{n,\mathbf{k},\sigma}$ satisfies,

$$\left\{ -\frac{\hbar^2}{2m}\Delta_{\mathbf{r}} + v_{\text{eff}}(\mathbf{r}) \right\} \phi_{n,\mathbf{k},\sigma}(\mathbf{r}) = \varepsilon_{n,\mathbf{k}}\phi_{n,\mathbf{k},\sigma}(\mathbf{r}) \,.$$

2) The Wannier function

A localized function at a lattice point $\mathbf{r}_i$ can be made as,

$$\phi_{n,\sigma}(\mathbf{r} - \mathbf{r}_i) = \frac{1}{\sqrt{N}} \sum_{\mathbf{k}} e^{-i\mathbf{k}\mathbf{r}_i}\phi_{n,\mathbf{k},\sigma}(\mathbf{r})$$

Notes:

1) $\phi_{n,\mathbf{k},\sigma}(\mathbf{r})$ (and $\phi_{n,\sigma}(\mathbf{r} - \mathbf{r}_i)$) give an orthonormal complete set.
2) The definition of the Wannier func. has a degree of freedom.
   Cf. Marzari and Vanderbilt, PRB **56** (1997) 12847.

The Hubbard model : III.

Introduce creation (annihilation) operators $c_{n\mathbf{k}\sigma}^{\dagger}$ ($c_{n\mathbf{k}\sigma}$) satisfying

$$\psi_{\sigma}^{\dagger}(\mathbf{r}) = \sum_{n\mathbf{k}} \phi_{n,\mathbf{k},\sigma}^{*}(\mathbf{r}) c_{n\mathbf{k}\sigma}^{\dagger} \ ,$$

the effective Hamiltonian is written as,

$$\hat{H}_{\text{eff}} = \sum_{n\mathbf{k}\sigma} \varepsilon_{n,\mathbf{k}} c_{n\mathbf{k}\sigma}^{\dagger} c_{n\mathbf{k}\sigma} \ .$$

Introduce a unitary transformation,

$$c_{ni\sigma} = \frac{1}{\sqrt{N}} \sum_{\mathbf{k}} \sum_{m} e^{i\mathbf{k}\mathbf{r}_i} U_{mn}^{-1}(\mathbf{k}) c_{m\mathbf{k}\sigma}$$

$$c_{ni\sigma}^{\dagger} = \frac{1}{\sqrt{N}} \sum_{\mathbf{k}} \sum_{m} e^{-i\mathbf{k}\mathbf{r}_i} U_{nm}(\mathbf{k}) c_{m\mathbf{k}\sigma}^{\dagger} \ .$$

The effective Hamiltonian is rewritten as,

$$\hat{H}_{\text{eff}} = \sum_{nmij} \sum_{\sigma} t_{ij}^{nm} (c_{ni\sigma}^{\dagger} c_{mj\sigma} + \text{H.c.})$$

This model is called the tight-binding model, where $t_{ij}^{nm} = \frac{1}{N} \sum_{\mathbf{k}} e^{i\mathbf{k}(\mathbf{r}_i - \mathbf{r}_j)} U_{nl}^{-1}(\mathbf{k}) \varepsilon_{l,\mathbf{k}} U_{lm}(\mathbf{k})$.

The Hubbard model : IV.

Consider the Hartree approximation to make an effective potential. (The method below is applicable for the Hartree-Fock approximation, but we need another formulation for DFT.)

$$\hat{H}_0 = \sum_{n,\mathbf{k},\sigma} \varepsilon_{n,\mathbf{k}} c^\dagger_{n\mathbf{k}\sigma} c_{n\mathbf{k}\sigma}$$

$$\hat{H}_1 = \frac{1}{2} \int d^3r d^3r' \frac{e^2}{|\mathbf{r}-\mathbf{r'}|} \sum_{\sigma\sigma'} \psi^\dagger_\sigma(\mathbf{r}) \psi^\dagger_{\sigma'}(\mathbf{r'}) \psi_{\sigma'}(\mathbf{r'}) \psi_\sigma(\mathbf{r})$$

$$- \left[\int d^3r d^3r' \frac{e^2}{|\mathbf{r}-\mathbf{r'}|} \rho(\mathbf{r}) \sum_\sigma \psi^\dagger_\sigma(\mathbf{r}) \psi_\sigma(\mathbf{r}) - \frac{1}{2} \int d^3r d^3r' \frac{e^2}{|\mathbf{r}-\mathbf{r'}|} \rho(\mathbf{r})\rho(\mathbf{r'}) \right]$$

Expand $\psi^\dagger_\sigma(\mathbf{r})$ $(\psi_\sigma(\mathbf{r}))$ into $c^\dagger_{n\mathbf{k}\sigma}$ $(c_{n\mathbf{k}\sigma})$, utilize the unitary transformation into $c^\dagger_{ni\sigma}$ $(c_{ni\sigma})$, then H_1 is rewritten as,

$$\hat{H}_1 = \frac{1}{2} \sum_{ijkl} \sum_{m_1 \cdots m_4} \sum_{\sigma\sigma'} (m_1 i, m_2 j | \frac{1}{r} | m_3 k, m_4 l) c^\dagger_{m_1 i\sigma} c^\dagger_{m_2 j\sigma'} c_{m_3 k\sigma'} c_{m_4 l\sigma}$$

$$- \sum_{ijkl} \sum_{m_1 \cdots m_4} \sum_\sigma (m_1 i, m_2 j | \frac{1}{r} | m_3 k, m_4 l) \nu_{m_2 j m_3 k} c^\dagger_{m_1 i\sigma} c_{m_4 l\sigma}$$

$$+ \frac{1}{2} \int d^3r d^3r' \frac{e^2}{|\mathbf{r}-\mathbf{r'}|} \rho(\mathbf{r})\rho(\mathbf{r'})$$

where $\nu_{m_1 i m_2 j} = (1/N) \Sigma_{n\mathbf{k}} e^{-i\mathbf{k}(\mathbf{r}_i - \mathbf{r}_j)} U_{nm_1}(\mathbf{k}) \langle n_{n\mathbf{k}} \rangle U^{-1}_{nm_2}(\mathbf{k})$.

The Hubbard model : V.

Consider a single band model with well-localized Wannier orbitals. The local interactions are classified into groups. Hubbard estimated the values as follows. (Hubbard 1963)

1. Intra-orbital repulsion: $U = (ii|\frac{1}{r}|ii) \sim 20$ eV

2. Inter-orbital repulsion: $V = (ij|\frac{1}{r}|ji) \sim 6$ eV

3. Correlated hopping: $X = (ii|\frac{1}{r}|ji) \sim 1/2$ eV

4. Off-site corr. hop.: $X' = (ij|\frac{1}{r}|ki) \sim 1/20$ eV

5. Inter-orbital exchange: $J = (ij|\frac{1}{r}|ij) \sim 1/40$ eV

6. Pair hopping: $J' = (ii|\frac{1}{r}|jj) \sim 1/40$ eV

Note that we may utilize screened values, if we consider screening effects due to electrons in other bands. Usually the counter term with ν_{ij} is assumed to be irrelavant for discussion of low-energy phenomena.

Thus, we may utilize the Hubbard model

$$H = - \sum_{\langle i,j \rangle} \sum_{\sigma} t_{i,j}(c_{i,\sigma}^{\dagger} c_{j,\sigma} + \text{H.c.}) + U \sum_{i} n_{i,\uparrow} n_{i,\downarrow}$$

for the discussion of magnetism.

The Hubbard model : VI.

Introduce a creation operator (an annihilation operator) $c_{i,\sigma}^{\dagger}$ $(c_{i,\sigma})$ for each local orbital. They satisfy commutation relation.

$$\{c_{i,\sigma}, c_{j,\sigma'}^{\dagger}\} = \delta_{i,j}\delta_{\sigma,\sigma'} \ .$$

The Hubbard model is given by,

$$H = -\sum_{\langle i,j\rangle}\sum_{\sigma}t_{i,j}(c_{i,\sigma}^{\dagger}c_{j,\sigma} + \text{H.c.}) + U\sum_{i}n_{i,\uparrow}n_{i,\downarrow}$$

- $t_{i,j}$: Transfer integral

- U : The Hubbard interaction parameter

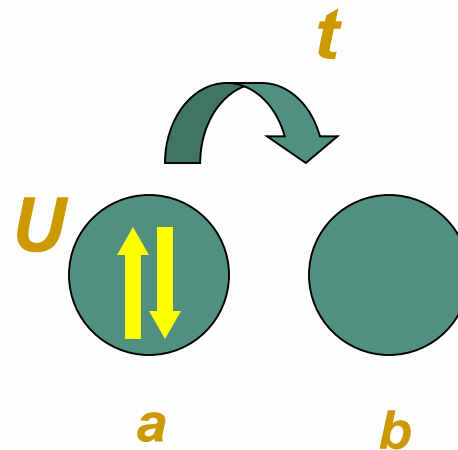# 2-site Hubbard model : I.

Consider a 2-site Hubbard model.

Each site allows 4 states, $|0\rangle$, $c_{i,\uparrow}^{\dagger}|0\rangle$, $c_{i,\downarrow}^{\dagger}|0\rangle$, and $c_{i,\uparrow}^{\dagger}c_{i,\downarrow}^{\dagger}|0\rangle$. $(i = a, b)$

- States available : $4^2 = 16$

- States with 2 electrons : 6.

- States with one $\uparrow$ spin and one $\downarrow$ spin : 4.

We consider the next 4 states.

- $|a,\uparrow;b,\downarrow\rangle \equiv c_{a,\uparrow}^{\dagger}c_{b,\downarrow}^{\dagger}|0\rangle$.

- $|a,\downarrow;b,\uparrow\rangle \equiv c_{a,\downarrow}^{\dagger}c_{b,\uparrow}^{\dagger}|0\rangle$.

- $|a,\uparrow,\downarrow;b,0\rangle \equiv c_{a,\uparrow}^{\dagger}c_{a,\downarrow}^{\dagger}|0\rangle$.

- $|a,0;b,\uparrow,\downarrow\rangle \equiv c_{b,\uparrow}^{\dagger}c_{b,\downarrow}^{\dagger}|0\rangle$.

2-site Hubbard model : II.

The Hamiltonian is given by,

$$H = -t \sum_\sigma (c_{b,\sigma}^\dagger c_{a,\sigma} + c_{a,\sigma}^\dagger c_{b,\sigma}) + U \sum_{i=a,b} n_{i,\uparrow} n_{i,\downarrow} \tag{8}$$

A matrix representation of H in terms of the above 4 states are,

$$H \doteq \begin{bmatrix} 0 & 0 & t & t \\ 0 & 0 & -t & -t \\ t & -t & U & 0 \\ t & -t & 0 & U \end{bmatrix} \tag{9}$$

Using a unitary transformation U which gives spin eigen states,

- $|S_{tot} = 0, \alpha\rangle \equiv \frac{1}{\sqrt{2}} (c_{a,\uparrow}^\dagger c_{b,\downarrow}^\dagger - c_{a,\downarrow}^\dagger c_{b,\uparrow}^\dagger)|0\rangle.$

- $|S_{tot} = 0, \beta\rangle \equiv \frac{1}{\sqrt{2}} (c_{a,\uparrow}^\dagger c_{a,\downarrow}^\dagger + c_{b,\uparrow}^\dagger c_{b,\downarrow}^\dagger)|0\rangle.$

- $|S_{tot} = 0, \gamma\rangle \equiv \frac{1}{\sqrt{2}} (c_{a,\uparrow}^\dagger c_{a,\downarrow}^\dagger - c_{b,\uparrow}^\dagger c_{b,\downarrow}^\dagger)|0\rangle.$

- $|S_{tot} = 1, S_z = 0\rangle \equiv \frac{1}{\sqrt{2}} (c_{a,\uparrow}^\dagger c_{b,\downarrow}^\dagger + c_{a,\downarrow}^\dagger c_{b,\uparrow}^\dagger)|0\rangle.$

2-site Hubbard model : III.

U is defined by,

$$U \doteq \begin{bmatrix} \frac{1}{\sqrt{2}} & 0 & 0 & \frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{2}} & 0 & 0 & \frac{1}{\sqrt{2}} \\ 0 & \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 \\ 0 & \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} & 0 \end{bmatrix} \tag{10}$$
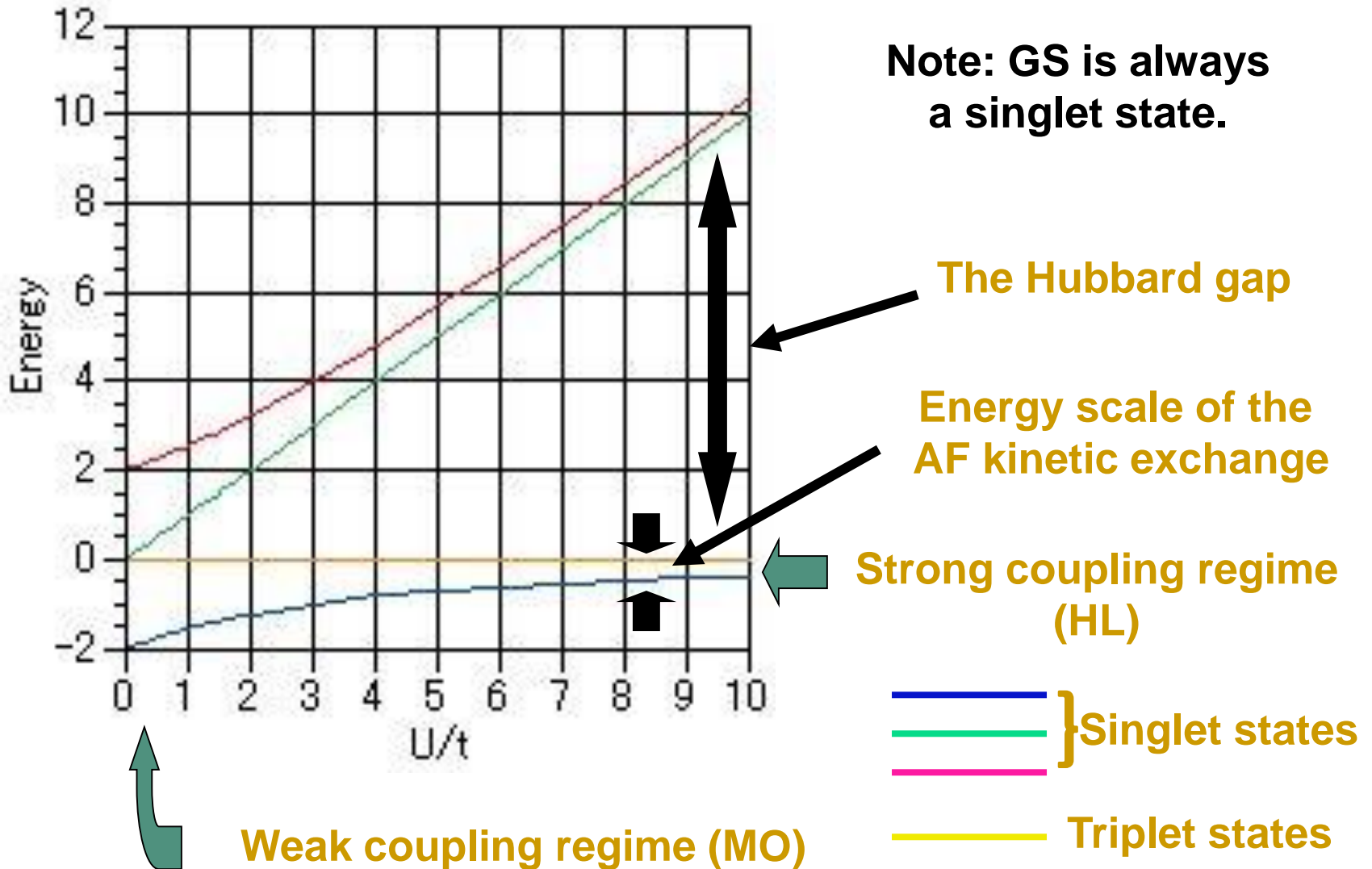
Then, we have a block-diagonalized H,

$$U^{-1}HU \doteq \begin{bmatrix} 0 & 2t & 0 & 0 \\ 2t & U & 0 & 0 \\ 0 & 0 & U & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \tag{11}$$

and eigen values are

- $E_1 = \frac{1}{2}\left(U - \sqrt{U^2 + 16t^2}\right) \sim -\frac{4t^2}{U}$.
- $E_2 = \frac{1}{2}\left(U + \sqrt{U^2 + 16t^2}\right) \sim U + \frac{4t^2}{U}$.
- $E_3 = U$.   • $E_4 = 0$.

Crossover from weak coupling regime to strong coupling regime

2-site Hubbard model : IV.

Note: GS is always a singlet state.

The Hubbard gap

Energy scale of the AF kinetic exchange

Strong coupling regime (HL)

}**Singlet states**

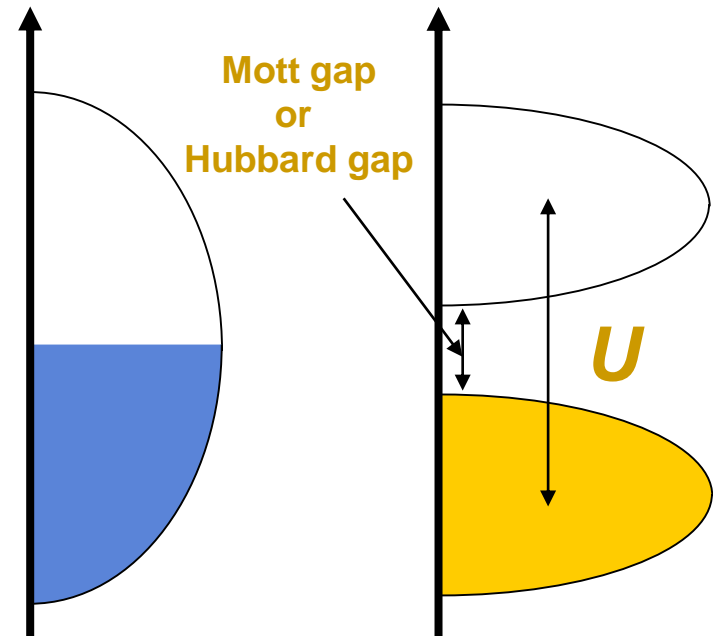**Triplet states**

Weak coupling regime (MO)

Concept of the Mott insulator

■ Electron transfer from an orbital to an occupied orbital is prohibited by the Coulomb repulsion. This leads to an insulating state.
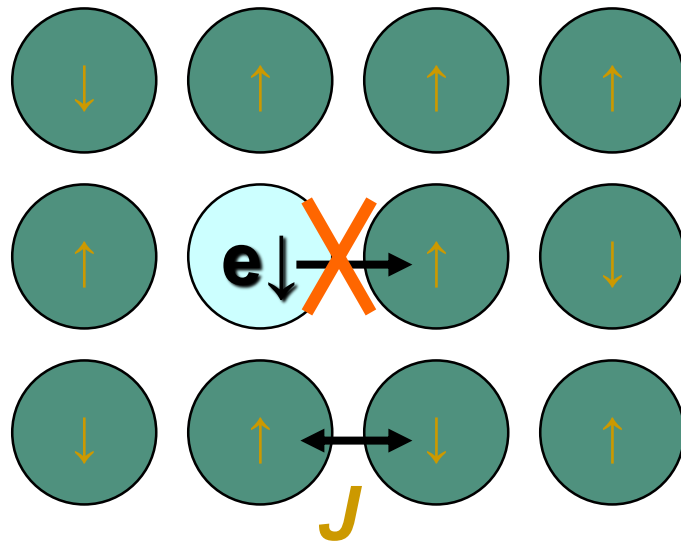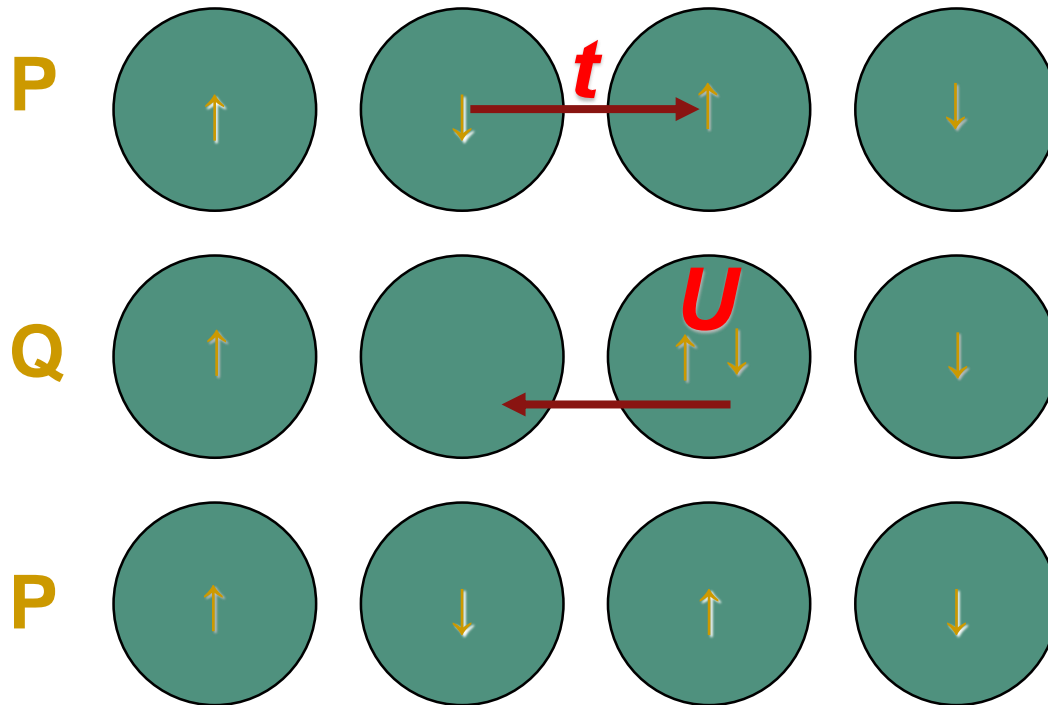
The half-filled Hubbard model shows the Mott insulating ground state.

Formation of the Mott gap in the single particle excitation.

Mott gap
or
Hubbard gap

J

U

e↓

The kinetic exchange

- Effective interaction at around $U=\infty$


P   **Q**   **P**   with labels *t*, *U*

We have another process starting from transfer of the up spin moving to the neighboring site.

$$H_{eff} = \frac{2t^2}{U} \sum_{\langle i,j \rangle} (\mathbf{S}_i \cdot \mathbf{S}_j - \frac{1}{4})$$

Antiferromagnetic Heisenberg model.

Superexchange interaction

■ In the transition-metal oxides, there is a structure of M-O-M (M: transition metal, O: oxygen). Along this structure, two localized spins couples via the superexchane through non-magnetic oxygen atom.

Often, Mott insulator is formed.
→ M has an open d-shell.

d_{x2-y2} @ M

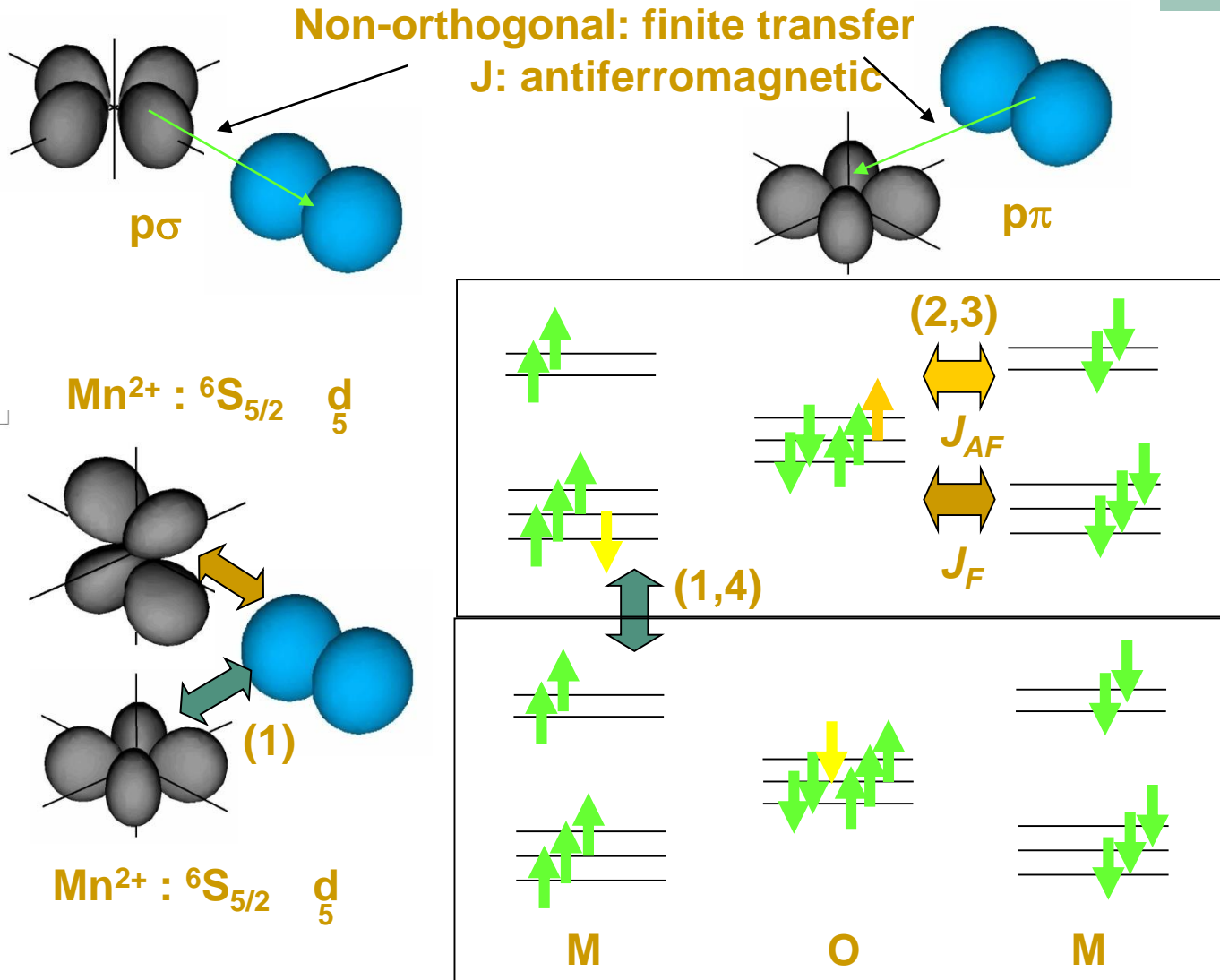p_x @ O

d_{x2-y2} @ M

J_{AF}
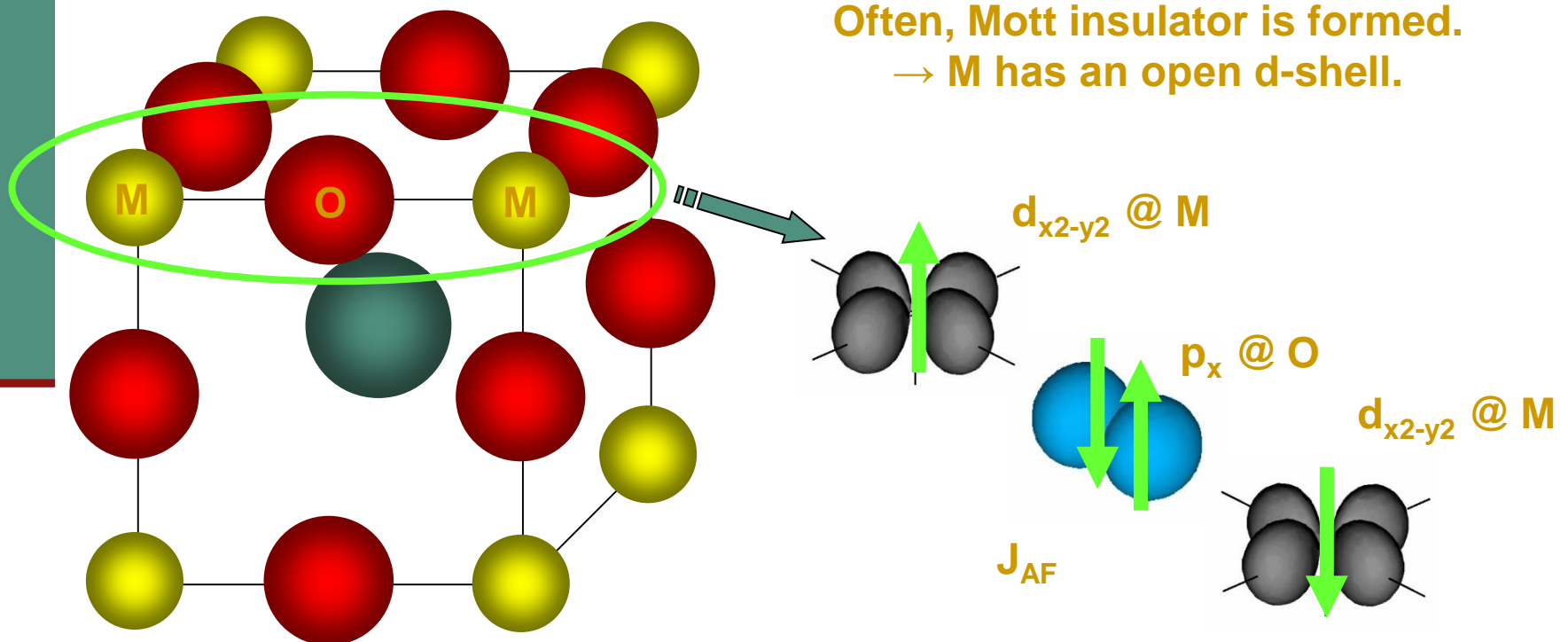
The Kanamori-Goodenough rule

- Consider description by non-orthogonalized atomic orbitals (LCAO picture)

- The electron transfer is possible from O to M when two neiboring orbitals are not orthogonal.

- The sign of exchange interaction between d and p is,
    - Ferromagnetic if two orbitals are orthogonal with each other,
    - Antiferromagnetic if two orbitals are not orthogonal

- Occupation of d-orbitals are determined by the crystal field splitting and the Hund rule.

Following the above rule, we can determine the sign of superexchange in a qualitative manner.

P.W. Anderson reformulated the rule using the orthogonalized Wannier basis allowing the second quantization scheme.

Exanple of antiferromagnetic exchange

Non-orthogonal: finite transfer
J: antiferromagnetic

pσ

pπ

Mn^{2+} : $^6S_{5/2}$ d$_5$

Mn^{2+} : $^6S_{5/2}$ d$_5$

(1)

(2,3)

J_{AF}

J_F

(1,4)

Non-perturbative States in exci-tation

Non-perturbative ground state

U

M   O   M

Various types of exchange interaction

- Direct exchange : (Cf. Heitler-London theory)
- Kinetic exchange : (Cf. the Hubbard model)
- Super exchange
- Double exchange
- RKKY interaction
- Anisotropic exchange interaction
- Dzyaloshinskii-Moriya interaction

1. **Exchange interactions except for the 1st one are effective interactions.**
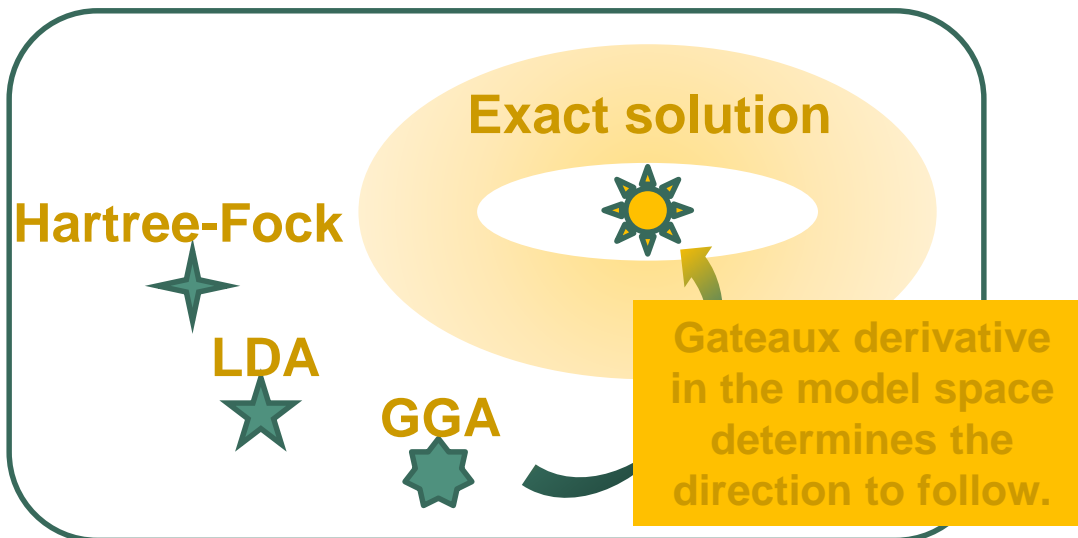2. **Effective interactions can be derived from a model with local interactions.**

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An almighty method

- Space of DFT models
  - Distance :

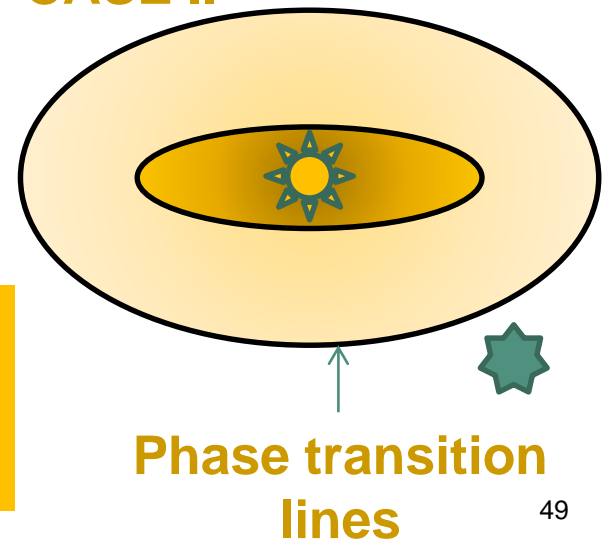$$\left\| n_i - n_{i+1} \right\|.$$

  - Variational principle: DFVT by K.K. (2009).

CASE I
ε-vicinity around exact solution

CASE II

Exact solution

Hartree-Fock

LDA

GGA

Gateaux derivative in the model space determines the direction to follow.

Phase transition lines

Super processes

To have an initial state of correlated electron system, we may utilize

The upconversion Hamiltonian

$$\left(\hat{H}_{1,i} + \hat{H}_2^{A_i \to A_i} - \hat{H}_2^{A_i^c \to A_i} P_B \frac{1}{\hat{H}_{1,i} + \hat{H}_2^{A_i \to A_i} + \hat{H}_2^{A_i \leftrightarrow A_i^c} - E} P_B \hat{H}_2^{A_i \to A_i^c} \right) \left| \Psi_i \right\rangle = E \left| \Psi_i \right\rangle.$$

which leads us to the convergent model series in the DFT model space.
Cf. K. Kusakabe, I. Maruyama, "Electronic state calculation method, electronic state calculation device, computer program", PCT filed No. PCT/JP2011/068589, date 2011.8.18.

Problems in static state

Magnetism   **Superconductivity**

Semi-classical & quantum dynamics

Quantum processes in QED!

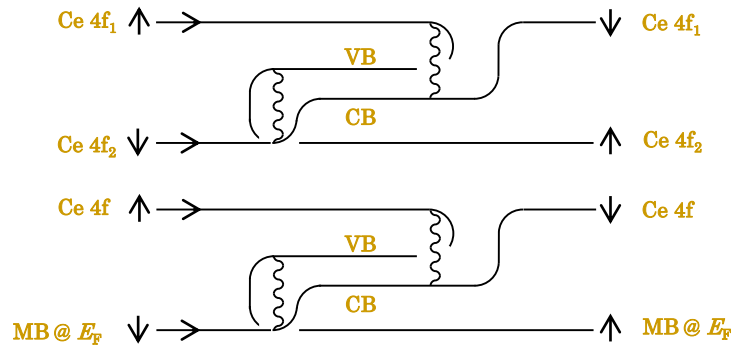| **The super-exchange process** | **The super-pair hopping process** | **The van-der-Waals interaction** | **Two & multi-photon processes** |

Exchange scattering

Ce 4f₁   VB   CB   **Ce 4f₁**

Evaluation of c-f exchange

$$\left(J_{eff-cf} \right)_{c_j f_i}$$

AF effective exchange
⇒ The Kondo Screening

Upconversion Hamiltonian with super processes

$$\left(\hat{H}_{1,i} + \hat{H}_2^{A_i \to A_i} - \hat{H}_2^{A_i^c \to A_i} P_B \frac{1}{\hat{H}_{1,i} + \hat{H}_2^{A_i \to A_i} + \hat{H}_2^{A_i \leftrightarrow A_i^c} - E} P_B \hat{H}_2^{A_i \to A_i^c} \right) |\Psi_i\rangle = E|\Psi_i\rangle,$$

$$H_{1,i} = \hat{T} + \hat{V} = -\frac{\hbar^2}{2m} \int d^3r \lim_{r' \to r} \hat{\psi}^\dagger(\mathbf{r}')\Delta_\mathbf{r}\hat{\psi}(\mathbf{r}) + \int d^3r \bar{v}_{eff}(\mathbf{r})\hat{n}(\mathbf{r}) + \frac{\delta}{\delta\langle\Psi|}E_{g_i}[\Psi],$$
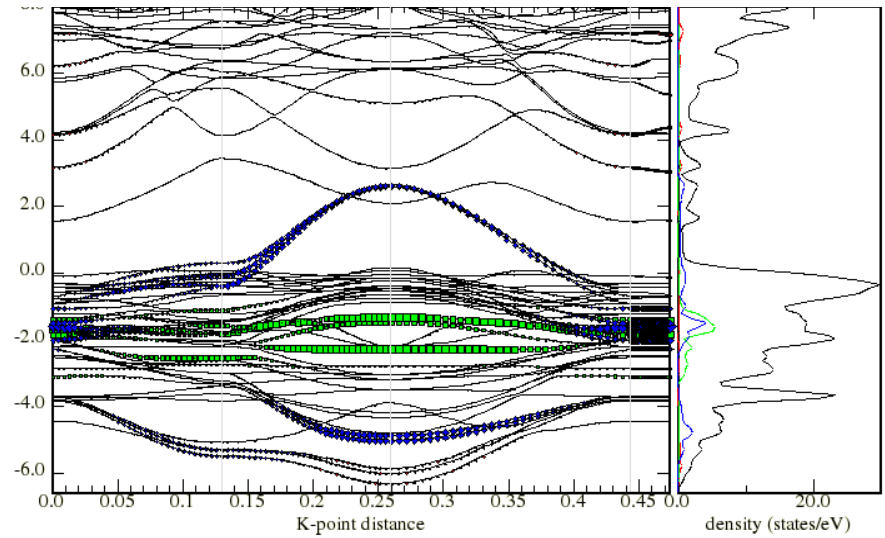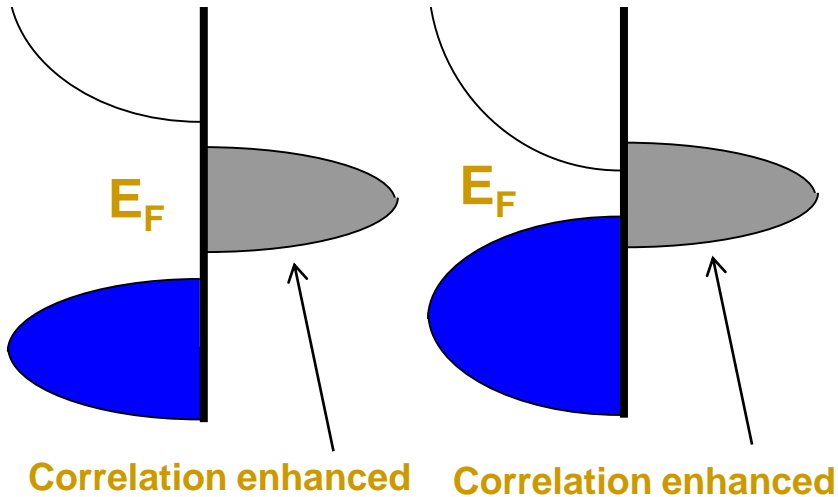
$$\bar{v}_{eff}(\mathbf{r}) = \frac{\delta}{\delta n(\mathbf{r})}\left\{ \frac{e^2}{2} \int d^3r d^3r' \frac{n(\mathbf{r})n(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} + E_{\varepsilon_i}[n] + \int d^3r v_{ext}(\mathbf{r})n(\mathbf{r}) \right\}.$$

A unified picture of high-T_c materials

Isolated center band   **Embedded center band**   **Band structure of Tl-cuprates**

E_F   **E_F**

Correlation enhanced   **Correlation enhanced**

Empirical rule :
1. **Embedded type preferable**
2. **Larger band width of center band preferable**
3. **Correlation induced by localized Wannier orbital preferable**

Evaluation of matrix elements

$(V_{\text{ee}})_{(n,\mathbf{K}),\mathbf{k}':\mathbf{k}'(m,\mathbf{K})}$

$$= \int d^3 r\, d^3 r' \frac{e^2 \phi^*_{\mathbf{k}'}(\mathbf{r}) \varphi^*_{n,\mathbf{K}}(\mathbf{r}') \phi_{\mathbf{k}'}(\mathbf{r}') \varphi_{m,\mathbf{K}}(\mathbf{r})}{2|\mathbf{r} - \mathbf{r}'|}.$$

4x4x1 mesh

| Material	Bloch diag.	BI off-diag.	On-site	n.n. off-site
Hg-comp.	18.6 eV at Γ	12.0 eV etc.	10.2 eV	6.4 eV
Tl-comp.	18.6 eV at Γ	5.4 eV etc.	8.3 eV	5.1 eV

Process A, C : A correlator as energy increment by fluctuation in the number density of particles in the correlated band allows an approximated form via introduction of approximated Green's functions (resolvents).

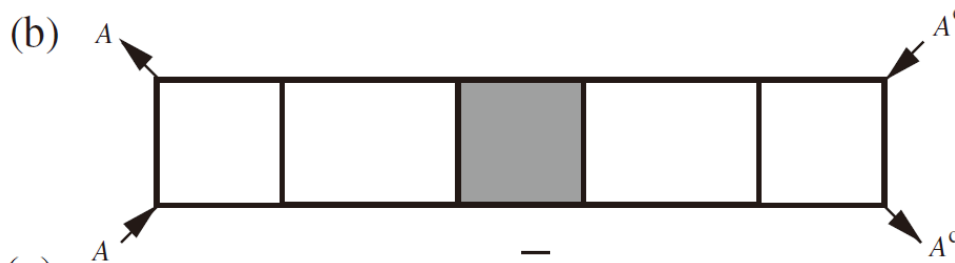Process B : A gap in un-correlated bands allows an approximated form as,

$$J^{\text{eff}}_{\mathbf{k},\mathbf{k}'} = - \sum_{\mathbf{K}, n \in A_1^c, m \in A_2^c} \frac{(V_{\text{ee}})_{\mathbf{k},(m,\mathbf{K}):(n,\mathbf{K}),\mathbf{k}}(V_{\text{ee}})_{(n,\mathbf{K}),\mathbf{k}':\mathbf{k}'(m,\mathbf{K})}}{\epsilon_{\mathbf{K},n} + \tilde{\epsilon}_{\mathbf{K},m}}.$$

$$H_{\text{scat}} = \sum_{\mathbf{k},\mathbf{k}',\sigma,\sigma'} J^{\text{eff}}_{\mathbf{k},\mathbf{k}'} c^\dagger_{\mathbf{k}',\sigma} c^\dagger_{\mathbf{k},\sigma'} c_{\mathbf{k}',\sigma'} c_{\mathbf{k},\sigma}.$$
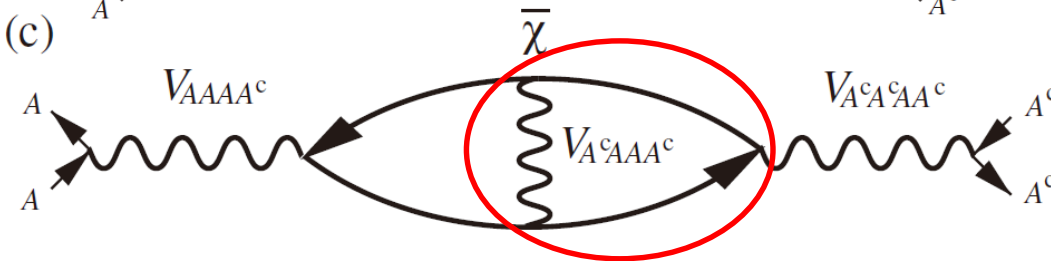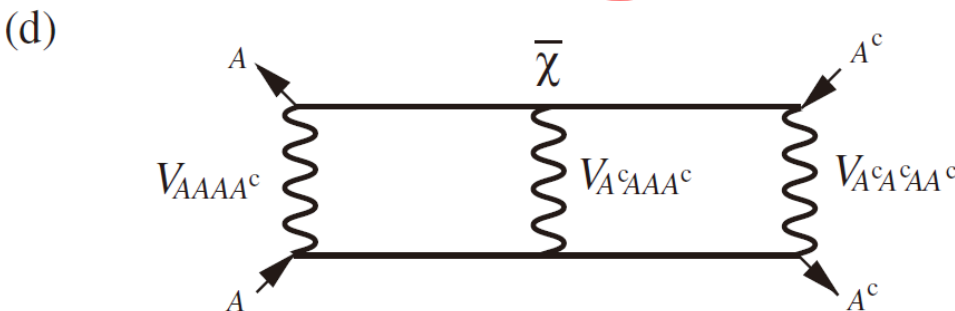
Super process as two-particle G.F.
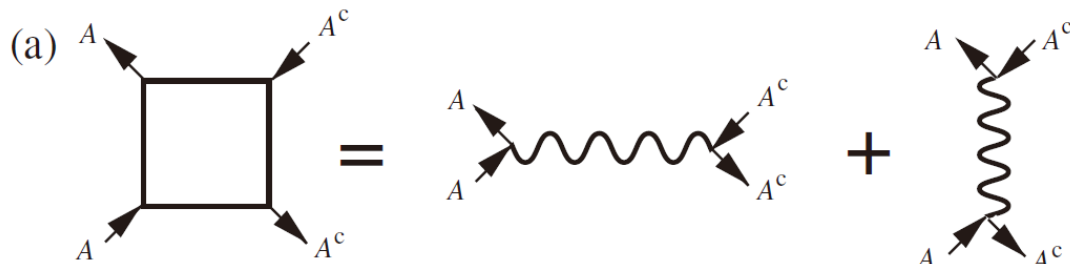
(a) Symmetric form of the bare Coulomb interaction

(b) Effective interaction diagram leading the Bethe-Salpeter equation (BSE)

(c) Bubble diagram by the polarization function with vertex correction (red part)

(d) Ladder diagram

Derivation of screened interaction

$$\hat{V}_{X_1 X_2 X_3 X_4} = \sum_{l_1 \in X_1, l_2 \in X_2, l_3 \in X_3, l_4 \in X_4} \sum_{\sigma, \sigma'} V_{l_1 l_2 : l_3 l_4} c^{\dagger}_{l_1,\sigma} c^{\dagger}_{l_2,\sigma'} c_{l_3,\sigma'} c_{l_4,\sigma}.$$

Bare interaction term

$$V^{\mathrm{scr}}_{A^c A A A^c}(\omega)$$

$$= V_{A^c A A A^c} + V_{A A A A^c} \chi V_{A^c A^c A A^c} + V_{A A A A^c} \chi V_{A A^c A A^c} \chi V_{A^c A^c A A^c}$$

$$\quad + V_{A A A A^c} \chi V_{A A^c A A^c} \chi V_{A A^c A A^c} \chi V_{A^c A^c A A^c} + \cdots$$

Polarization function

$$= V_{A^c A A A^c} + V_{A A A A^c} \chi \frac{1}{1 - V_{A A^c A A^c} \chi} V_{A^c A^c A A^c}$$

$$= V_{A^c A A A^c} + V_{A A A A^c} \chi^{1/2} \frac{1}{1 - \chi^{1/2} V_{A A^c A A^c} \chi^{1/2}} \chi^{1/2} V_{A^c A^c A A^c}.$$

Screening in c-RPA level : Vertex correction can be considered.

$$V^{\mathrm{scr}}_{A^c A A A^c}(\omega)$$

$$= V^{(l)}_{A A} \frac{1}{1 - V^{(r)}_{A A^c} \chi(\omega) V^{(l)}_{A^c A}} V^{(r)}_{A^c A^c} = V^{(l)}_{A A} \epsilon^{-1}(\omega) V^{(r)}_{A^c A^c}.$$

Screened interaction term

Super process in general forms

General operator expression

$$\hat{H}_{\text{super}} = P_A \hat{H}_{AB} \left(E - \hat{H}_A - \hat{h}_B - P_B \hat{H}_{AB} - \hat{H}_{BA} \right)^{-1} \hat{H}_{BA} P_A.$$

$$\hat{G} = \left(E - \hat{H}_A - \hat{h}_B - P_B \hat{H}_{AB} - \hat{H}_{BA} \right)^{-1}.$$
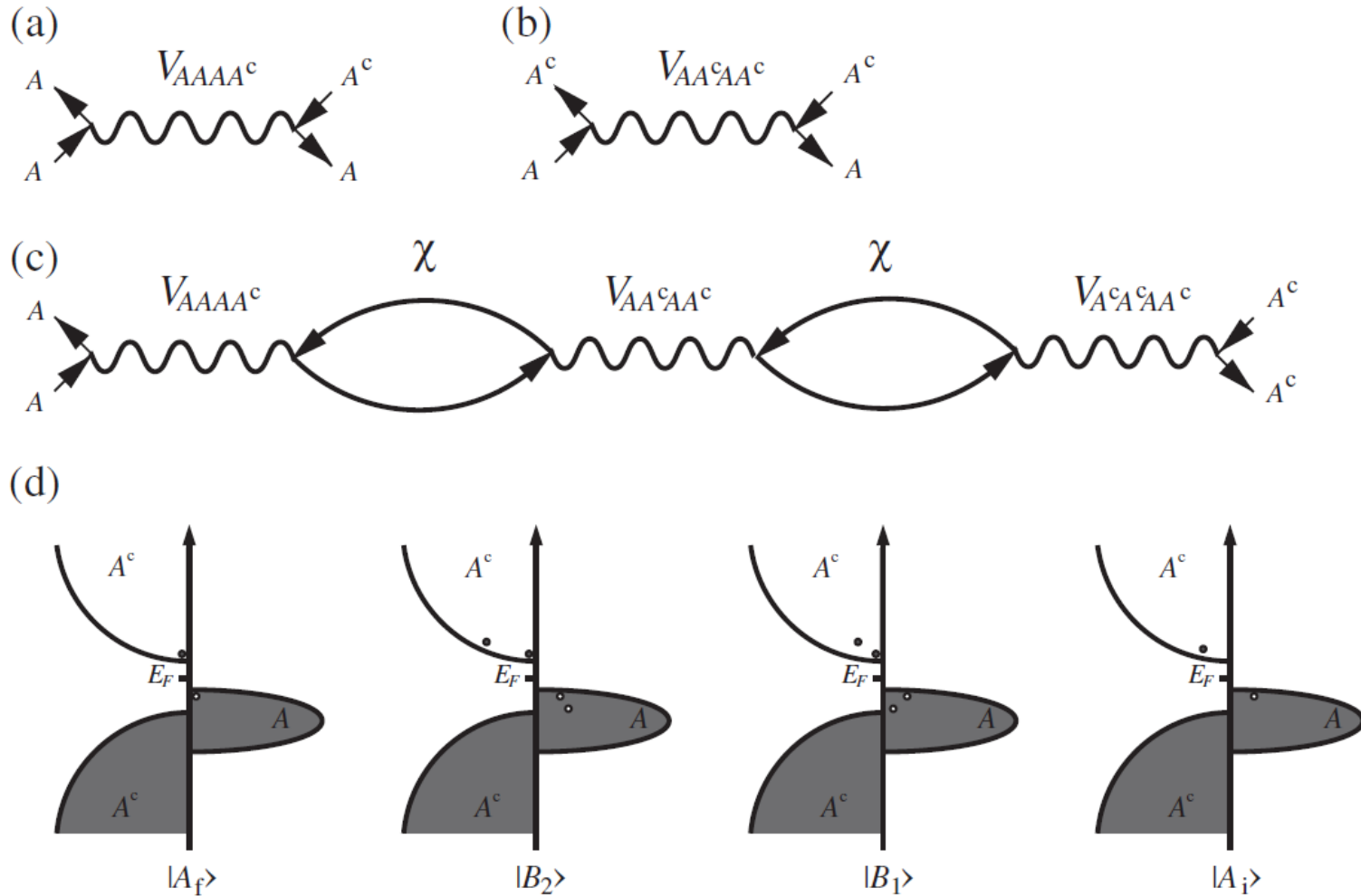
Phase-space dependent expression

$$\hat{H}_{\text{super}}$$

$$= P_A \sum_{i=1}^{\mathcal{N}} |A_i\rangle\langle A_i| \hat{H}_{AB} \sum_{k=1}^{\infty} |B_k\rangle\langle B_k| \hat{G}^{(B)} \sum_{l=1}^{\infty} |B_l\rangle\langle B_l| \hat{H}_{BA} \sum_{j=1}^{\mathcal{N}} |A_j\rangle\langle A_j| P_A$$

$$= P_A \sum_{i,j} |A_i\rangle \sum_{k,l} \langle A_i|\hat{H}_{AB}|B_k\rangle\langle B_k|\hat{G}^{(B)}|B_l\rangle\langle B_l|\hat{H}_{BA}|A_j\rangle\langle A_j| P_A$$

$$= \sum_{i,j} P_A |A_i\rangle \left(H_{\text{super}} \right)_{ij} \langle A_j| P_A.$$

A matrix element in an approximated expression corresponding to a 3rd order diagram

$$\left(H_{\text{super}} \right)^{(d)}$$

$$= (V_{AAAA^c})_{2\bar{2}_f\bar{1}_f2} \chi(\omega = E - (e_2 + \bar{e}_1))$$

$$\times (V_{AA^cAA^c})_{2112} \chi(\omega = E - (e_2 + \bar{e}_1))(V_{A^cA^cAA^c})_{2_i111_i}.$$

Electron-hole interaction in semiconductors

The GW scheme

$$\left[\hat{T} + \hat{V}_{ext} + \hat{V}_{Hartree} + \hat{V}_{xc}\right]\left|\psi_{i,\sigma}^{KS}\right\rangle = \varepsilon_{i,\sigma}^{KS}\left|\psi_{i,\sigma}^{KS}\right\rangle.$$

Kohn-Sham equation

$$\left[\hat{T} + \hat{V}_{ext} + \hat{V}_{Hartree} + \hat{\Sigma}\left(\varepsilon_{i,\xi}^{QP}\right)\right]\left|\psi_{i,\xi}^{QP}\right\rangle = \varepsilon_{i,\xi}^{QP}\left|\psi_{i,\xi}^{QP}\right\rangle.$$

Quasi-particle equation

$$\left\langle\mathbf{r},\sigma\left|\hat{\Sigma}(\omega)\right|\psi_{i,\xi}^{QP}\right\rangle = \int d\mathbf{r}' \sum_{\sigma'}\left\langle\mathbf{r},\sigma\left|\hat{\Sigma}(\omega)\right|\mathbf{r}',\sigma\right\rangle\delta_{\sigma,\sigma'}\left\langle\mathbf{r}',\sigma'\left|\psi_{i,\xi}^{QP}\right\rangle\right.$$

Self-energy expression

$$= \int d\mathbf{r}'\Sigma_{\sigma=\xi}(\mathbf{r},\mathbf{r}';\omega)\psi_{i,\xi}^{QP}(\mathbf{r}').$$

$$\Sigma_{\sigma}(\mathbf{r},\mathbf{r}';\omega) = \frac{i}{2\pi}\int d\omega' e^{i\omega'\delta}G_{\sigma}(\mathbf{r},\mathbf{r}';\omega+\omega')W(\mathbf{r},\mathbf{r}';\omega')d\omega'.$$

The GW self-energy

$$G_{\sigma}^{0}(\mathbf{r},\mathbf{r}';\omega) = \sum_{i}\frac{\psi_{i,\sigma}^{KS}(\mathbf{r})\left(\psi_{i,\sigma}^{KS}(\mathbf{r}')\right)^{*}}{\omega - \varepsilon_{i,\sigma}^{KS}}$$

The 0-th order Green function

$$W(\mathbf{r},\mathbf{r}';\omega) = \sum_{\mathbf{GG'}}e^{i(\mathbf{G'}+\mathbf{q})\cdot\mathbf{r}'}W_{\mathbf{GG'}}^{KS}(\mathbf{q};\omega)e^{-i(\mathbf{G}+\mathbf{q})\cdot\mathbf{r}'},$$

Screened interaction

$$W_{\mathbf{GG'}}^{KS}(\mathbf{q};\omega) = -4\pi e^{2}\frac{1}{\left|\mathbf{G}+\mathbf{q}\right|}\varepsilon_{\mathbf{GG'}}^{KS,-1}(\mathbf{q};\omega)\frac{1}{\left|\mathbf{G'}+\mathbf{q}\right|},$$
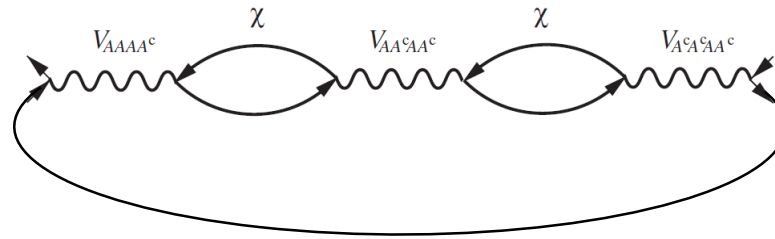
Screened interaction matrix

$$\varepsilon_{\mathbf{GG'}}^{KS}(\mathbf{q};\omega) = \delta_{\mathbf{G,G}} - 4\pi e^{2}\frac{1}{\left|\mathbf{G}+\mathbf{q}\right|}\chi_{\mathbf{GG'}}^{KS-0}(\mathbf{q};\omega)\frac{1}{\left|\mathbf{G'}+\mathbf{q}\right|},$$

Dielectric function

GW as an approximation

$$\left(\hat{H}_{1,i} + \hat{H}_2^{A_i \to A_i} - \hat{H}_2^{A_i^c \to A_i} P_B \frac{1}{\hat{H}_{1,i} + \hat{H}_2^{A_i \to A_i} + \hat{H}_2^{A_i \leftrightarrow A_i^c} - E} P_B \hat{H}_2^{A_i \to A_i^c} \right) |\Psi_i\rangle = E |\Psi_i\rangle.$$

$$|\Psi_i\rangle \Box |\psi_{i,\xi}^{QP}\rangle = \tilde{c}_{i,\xi}^\dagger |\Psi_0\rangle$$

$$-\hat{H}_2^{A_j^c \to A_j} P_B \frac{1}{\hat{H}_{1,\bar{i}} \delta_{\bar{i},\bar{j}} + \hat{\bar{H}}_2^{A_{\bar{i}}^c \to A_{\bar{j}}^c} - \omega \delta_{\omega,E}} P_B \hat{H}_2^{A_i \to A_i^c} |\psi_{i,\xi}^{QP}\rangle$$

$$\Box \, W_{j\xi,\bar{j},\bar{\xi};i\xi,\bar{i},\bar{\xi}}(\omega) G(\bar{j}, \bar{\bar{\xi}}; \bar{i}, \bar{\bar{\xi}}; \omega) \tilde{c}_{j,\xi}^\dagger \tilde{c}_{i,\xi} |\psi_{i,\xi}^{QP}\rangle$$

$$\Rightarrow \quad \left[\hat{T} + \hat{V}_{ext} + \hat{V}_{Hartree} + \hat{\Sigma}(\varepsilon_{i,\xi}^{QP}) \right] |\psi_{i,\xi}^{QP}\rangle = \varepsilon_{i,\xi}^{QP} |\psi_{i,\xi}^{QP}\rangle.$$

Thank you for your attention.

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